

Logic and Quantum Information

Lecture II: The Topology of Paradox

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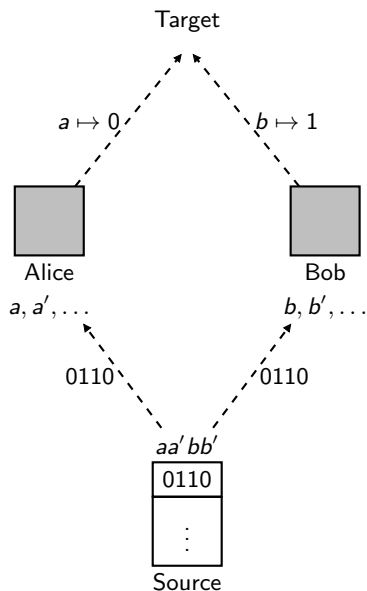
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However, this view is **impossible to sustain** in the light of our **actual observations of (micro)-physical reality**.

Hidden Variables: The Mermin instruction set picture



The 'Hardy Paradox': Bell's theorem without inequalities

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Hardy models: those whose **support** satisfies

	(0, 0)	(0, 1)	(1, 0)	(1, 1)
(a_1, b_1)	1			
(a_1, b_2)	0			
(a_2, b_1)	0			
(a_2, b_2)				0

Which 'instruction set' λ could the outcomes (0, 0) for measurements (a_1, b_1) could come? Clearly, we must have

$$\lambda : a_1 \mapsto 0, \quad b_1 \mapsto 0.$$

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However, this would require the outcome (0,0) for measurements (a_2, b_1) to be possible, and this is **precluded**.

Thus Hardy models are **contextual**. They cannot be explained by a classical source.

Bundle Pictures

Logical Contextuality

- Ignore precise probabilities
- Events are possible or not
- E.g. the Hardy model:

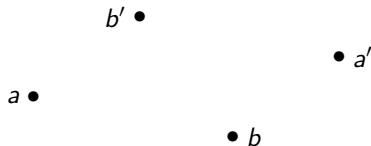
	00	01	10	11
ab	✓	✓	✓	✓
ab'	×	✓	✓	✓
$a'b$	×	✓	✓	✓
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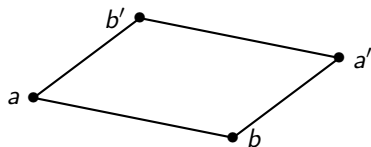


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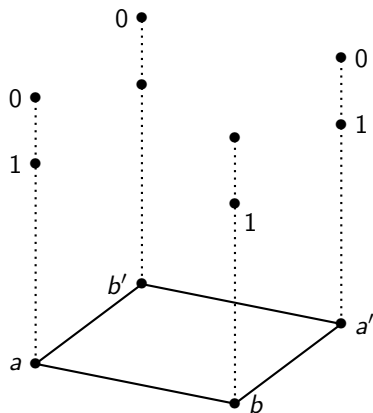


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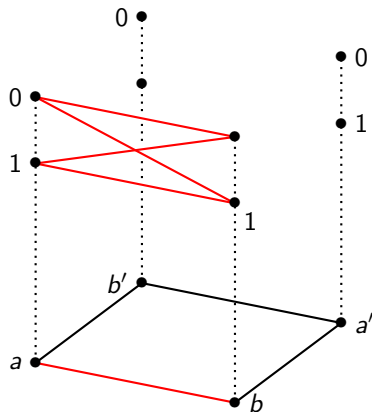


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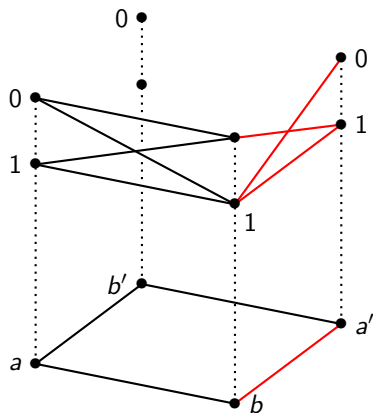


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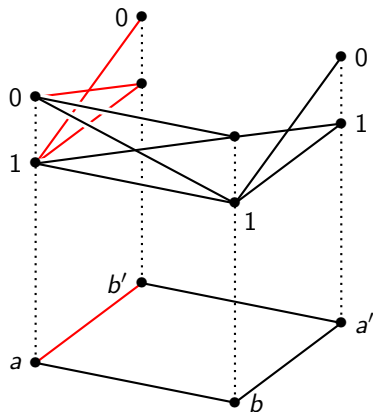


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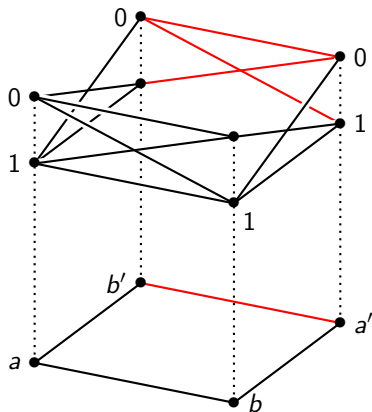


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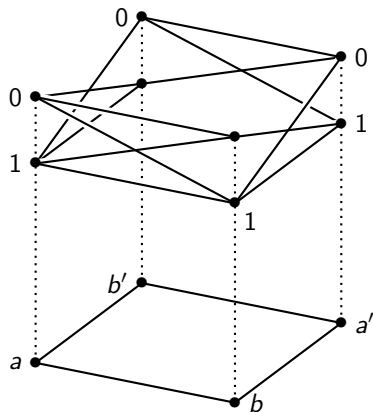


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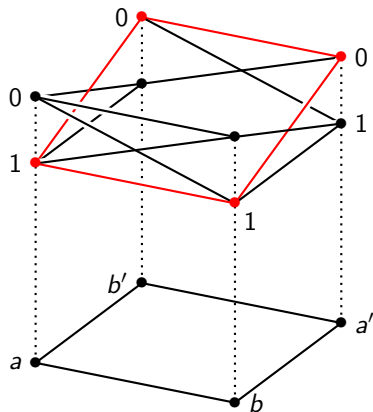


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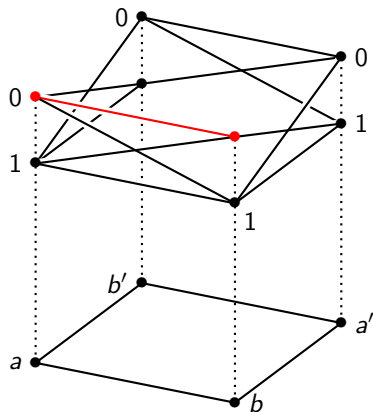


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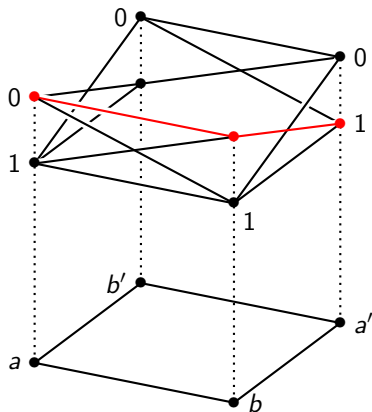


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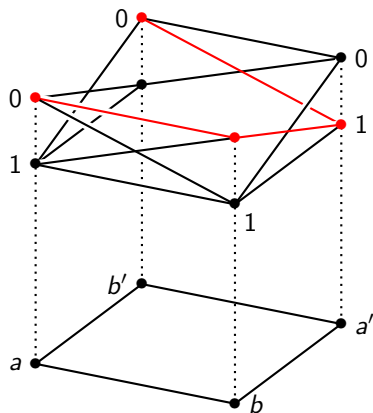


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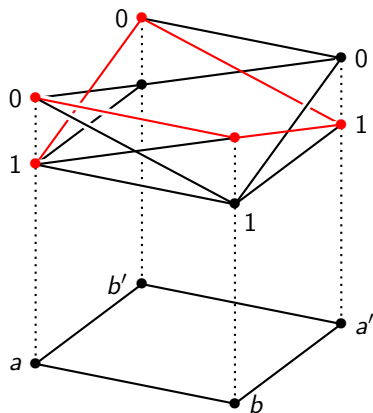


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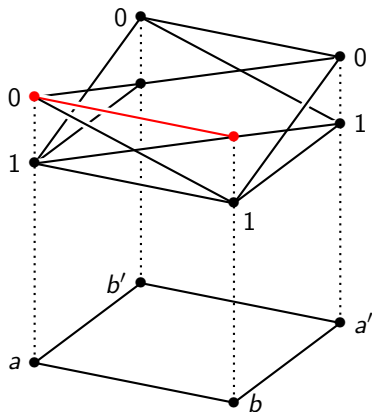


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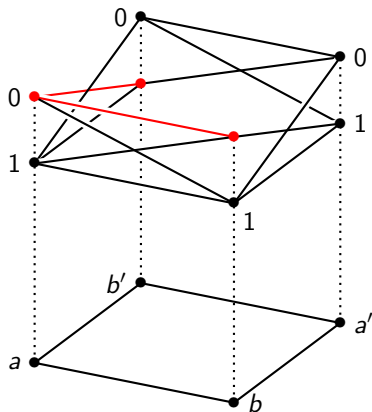


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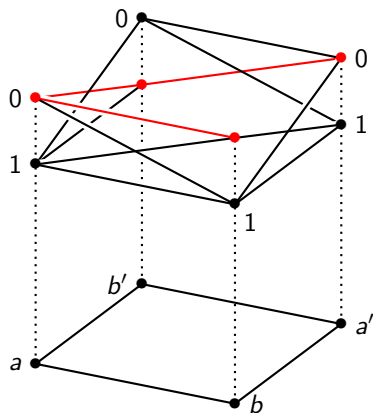


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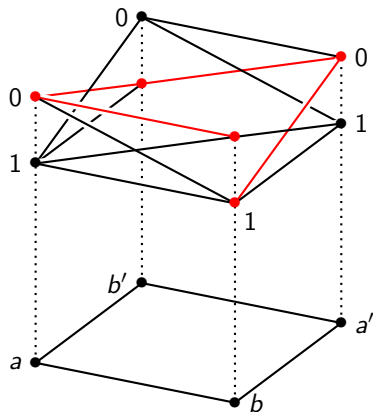


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Strong Contextuality

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a_1	b_2	1	0	0	1
a_2	b_1	1	0	0	1
a_2	b_2	0	1	1	0

The PR Box

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The PR Box

Note this achieves the algebraic maximum of 4 for our logical Bell inequality.

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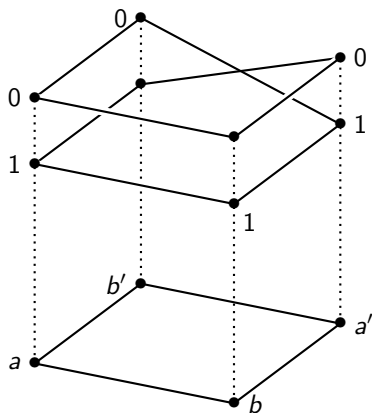
In terms of the XOR game, it is a **winning strategy**.

Bundle Pictures

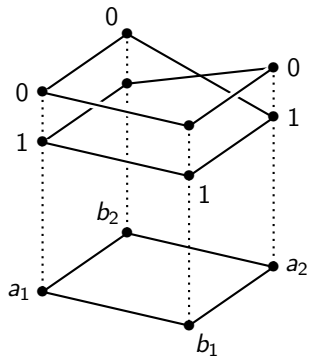
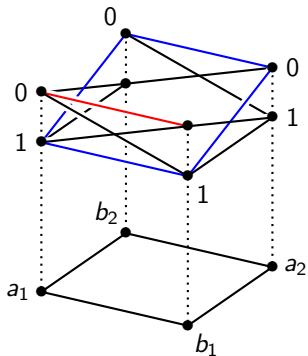
Strong Contextuality

- E.g. the PR box:

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ab'	✓	×	×	✓
$a'b$	✓	×	×	✓
$a'b'$	×	✓	✓	×

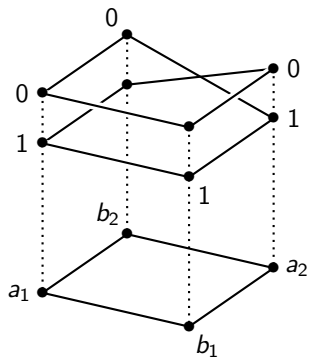
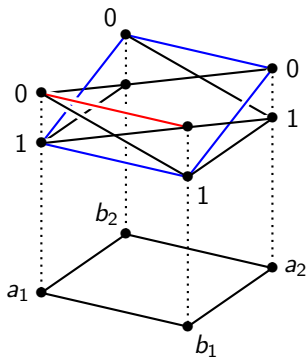


Visualizing Contextuality



The Hardy table and the PR box as bundles

Visualizing Contextuality



The Hardy table and the PR box as bundles

A hierarchy of degrees of contextuality:

$$\text{Bell} < \text{Hardy} < \text{GHZ}$$

Contextuality, Logic and Paradoxes

Contextuality, Logic and Paradoxes

Liar cycles. A Liar cycle of length N is a sequence of statements

$S_1 : S_2$ is true,

$S_2 : S_3$ is true,

\vdots

$S_{N-1} : S_N$ is true,

$S_N : S_1$ is false.

For $N = 1$, this is the classic Liar sentence

$S : S$ is false.

Contextuality, Logic and Paradoxes

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Following Cook, Walicki et al. we can model the situation by boolean equations:

$$x_1 = x_2, \dots, x_{n-1} = x_n, x_n = \neg x_1$$

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$$x_1 = x_2, \dots, x_{n-1} = x_n, x_n = \neg x_1$$

The “paradoxical” nature of the original statements is now captured by the inconsistency of these equations.

Contextuality in the Liar; Liar cycles in the PR Box

Contextuality in the Liar; Liar cycles in the PR Box

We can regard each of these equations as fibered over the set of variables which occur in it:

$$\{x_1, x_2\} : x_1 = x_2$$

$$\{x_2, x_3\} : x_2 = x_3$$

$$\vdots$$

$$\{x_{n-1}, x_n\} : x_{n-1} = x_n$$

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Any subset of up to $n - 1$ of these equations is consistent; while the whole set is inconsistent.

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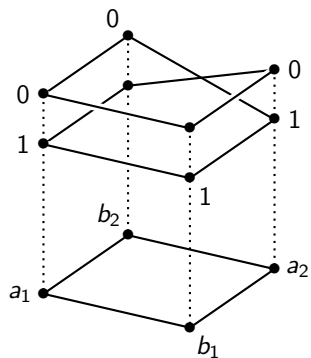
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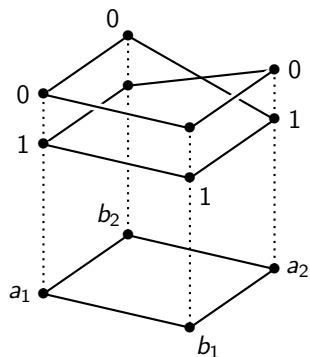
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The usual reasoning to derive a contradiction from the Liar cycle corresponds precisely to the attempt to find a univocal path in the bundle diagram.

Paths to contradiction



Paths to contradiction

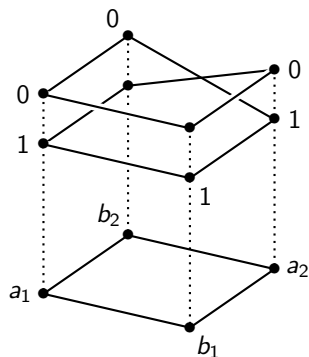


Suppose that we try to set a_2 to 1. Following the path on the right leads to the following local propagation of values:

$$a_2 = 1 \rightsquigarrow b_1 = 1 \rightsquigarrow a_1 = 1 \rightsquigarrow b_2 = 1 \rightsquigarrow a_2 = 0$$

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We have discussed a specific case here, but the analysis can be generalised to a large class of examples.

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This example is well-known in the quantum contextuality literature as the **Specker triangle**.

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A joint outcome or **event** in a context C is $s \in O^C$, e.g. $s = \{a \mapsto 0, b \mapsto 1\}$.

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C	F	C	G	M	N	D	F	M
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The original K-S construction used 117 variables!

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These distributions are the rows of our probability tables.

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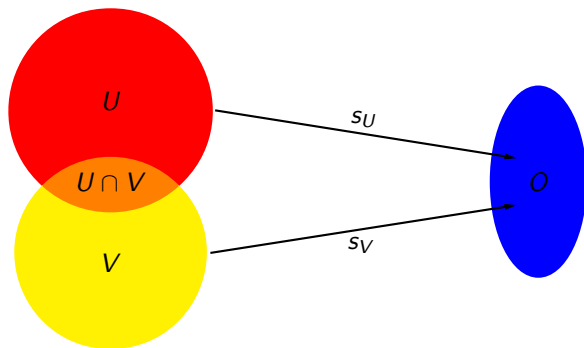
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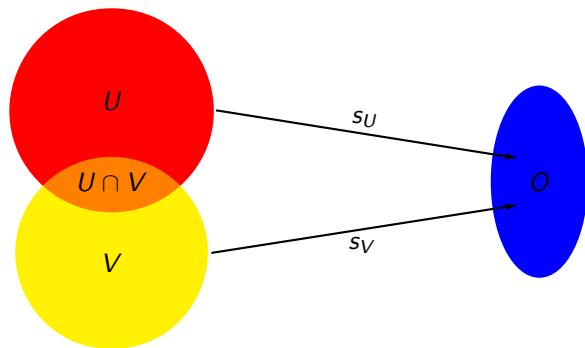
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Each row of the table specifies a **probability distribution** on events O^C for a given choice of measurements C .

Gluing functional sections



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If $s_U|_{U \cap V} = s_V|_{U \cap V}$, they can be glued to form

$$s : U \cup V \longrightarrow O$$

such that $s|_U = s_U$ and $s|_V = s_V$.

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So compatibility says that the distributions on different contexts have consistent marginals.

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- This is captured by saying that the distribution on $\{a\} = \{a, b\} \cap \{a, b'\}$ is the same whether we marginalize from the distribution e_C , or the distribution $e_{C'}$.
- This condition is generalized by compatibility – and this general form is satisfied by quantum systems.

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Consider the following schematic representation of an Alice-Bob table:

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The no-signalling conditions for the non-empty intersections of contexts are given by the following equations:

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Moreover, the PR box has a **unique family of distributions** which satisfy these conditions.

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An empirical model $\{d_C\}_{C \in \mathcal{M}}$ on a measurement scenario (X, \mathcal{M}, O) is **non-contextual** if there is a distribution $d \in \text{Prob}(O^X)$ such that, for all $C \in \mathcal{M}$:

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The import of Bell's theorem and similar results is that there are empirical models arising from quantum mechanics which are contextual.

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However, there are compatible (*i.e.* No-Signalling) empirical models which are **not** quantum realizable.

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It turns out that **all quantum realizable models are compatible**.

Compatibility is in fact the general form of an important physical principle known as No-Signalling, which ensures the consistency of quantum mechanics with Special Relativity.

However, there are compatible (*i.e.* No-Signalling) empirical models which are **not** quantum realizable.

We thus get a strict hierarchy of empirical models:

$$\text{NC} \subset \text{QM} \subset \text{NS}$$

The PR Box

A	B	(0,0)	(1,0)	(0,1)	(1,1)
a_1	b_1	1	0	0	1
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This satisfies No-Signalling, so is consistent with SR, but it is **not** quantum realisable.

Empirical Models as Vectors

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We can regard an empirical model $\{d_C\}_{C \in \mathcal{M}}$ as a vector

$$\mathbf{v} = (\mathbf{v}_{C,s})_{C \in \mathcal{M}, s \in O^C}, \quad \mathbf{v}_{C,s} := d_C(s)$$

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Moreover, convex combinations of compatible models are compatible.

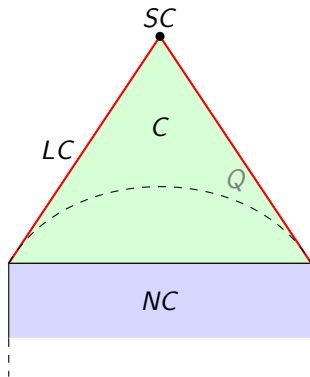
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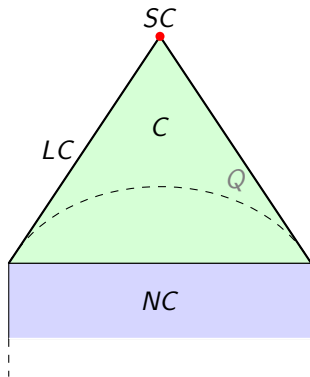
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Key question: find compelling principles to explain why Nature picks out the quantum set.

The Support of a Model

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This is equivalent to saying that, for all $C \subseteq C'$, the restriction map

$$\rho_C^{C'} : \mathcal{S}(C') \longrightarrow \mathcal{S}(C) \quad :: s \mapsto s|_C$$

is surjective.

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Obviously, strong contextuality implies logical contextuality.

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Thus in terms of well-known quantum examples, we have

Bell $<$ Hardy $<$ GHZ

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- We can similarly define logical contextuality for states; we say that a state is logically contextual if for some choice of local observables, the resulting empirical model is logically contextual; while the state is **not** strongly contextual.
- Finally, a state is weakly contextual if it is contextual, but neither of the previous two cases apply.

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We believe that an answer to this problem will shed considerable light on the structure of multipartite states, not least because it will necessitate solving the following task:

Given a multipartite state, find local observables which witness its highest degree of contextuality.

A Recent Result and a Question

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This question remains open, and appears difficult!