

# Phase transitions in the complexity of simulating random shallow quantum circuits



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Simons workshop

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# Complexity from entanglement

Original motivation for quantum computing [Feynman '82]



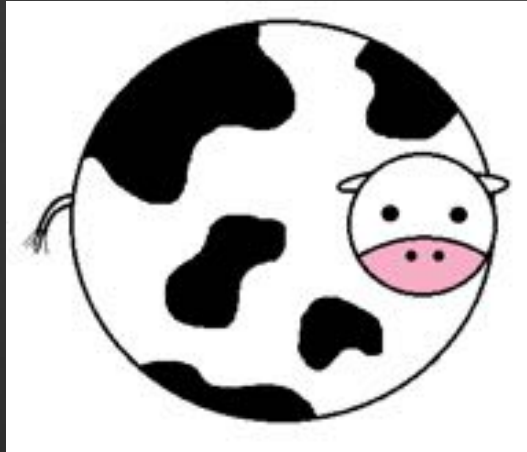
Nature isn't classical, dammit, and if you want to make a simulation of Nature, you'd better make it quantum mechanical, and by golly it's a wonderful problem, because it doesn't look so easy.

$N$  systems in **product state**  $\rightarrow O(N)$  degrees of freedom  
 $N$  **entangled** systems  $\rightarrow \exp(N)$  degrees of freedom

Describes cost of simulating dynamics or even describing a state.

**This talk:** do typical quantum dynamics achieve this?

# easier quantum simulations



1. solve trivial special case (e.g. non-interacting theory)
2. treat corrections to theory as perturbations



# easier quantum simulation

## Lightly entangling dynamics

product states + non-interacting gates are easy.  
Cost grows exponentially with # of entangling gates.

## Stabilizer circuits

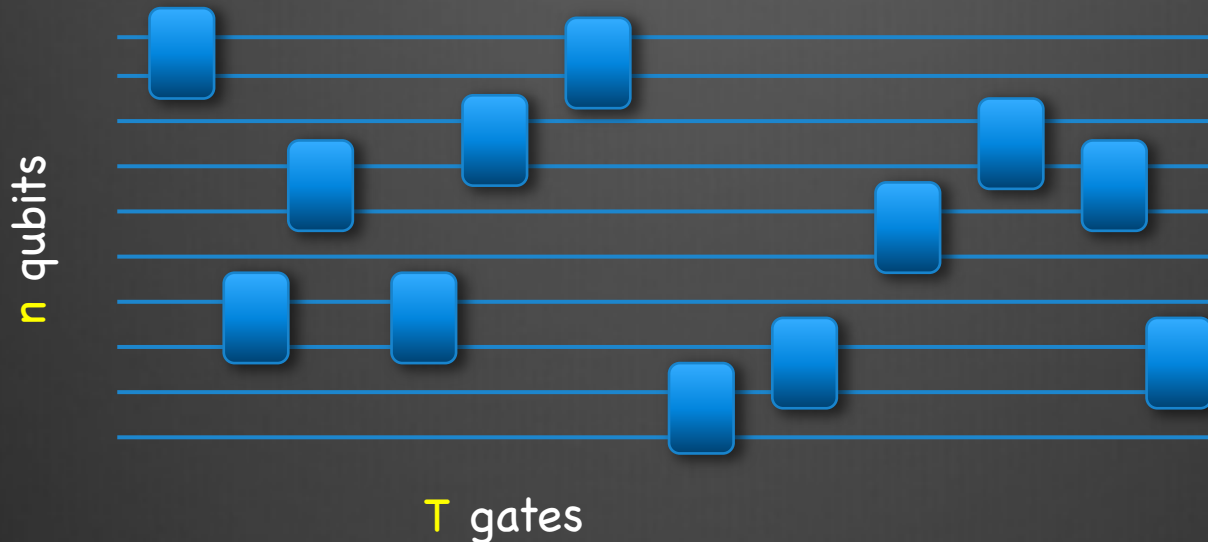
Poly-time simulation of stabilizer circuits,  
growing exponentially with # of non-stabilizer gates.

Likewise for matchgates / non-interacting fermions.

## Ground states of 1-D systems

Effort grows exponentially with correlation length.

# quantum circuits



Classical simulation possible in time  $O(T) \cdot \exp(k)$ , where

- $k$  = treewidth [Markov-Shi '05]
- $k$  = max # of gates crossing any single qubit [Yoran-Short '06, Jozsa '06]

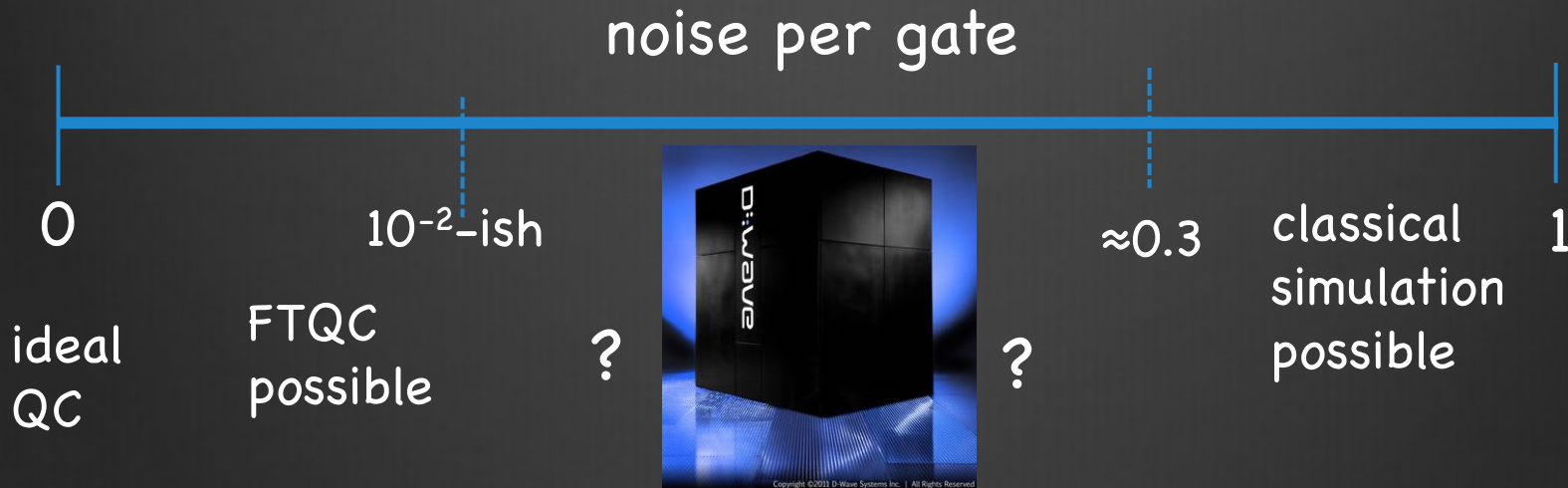
- + Complexity interpolates between linear and exponential.
- Treating all gates as "potentially entangling" is too pessimistic.



# noisy dynamics?

Time evolution of quantum systems

$$\frac{d\rho}{dt} = -i(H\rho - \rho H) + \text{noise terms that are linear in } \rho$$



conjectured to exhibit phase transition  
(possibly with intermediate phases)

# phase transitions?

Complexity **smoothly increases** with

- entanglement
- correlation length
- # of non-stabilizer gates

Complexity **jumps discontinuously** with

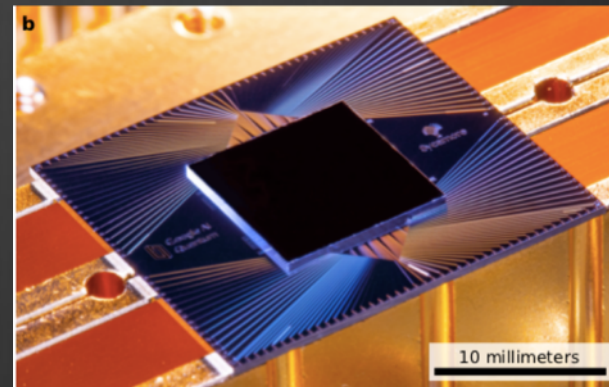
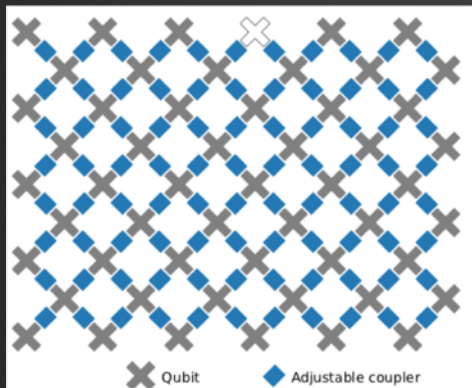
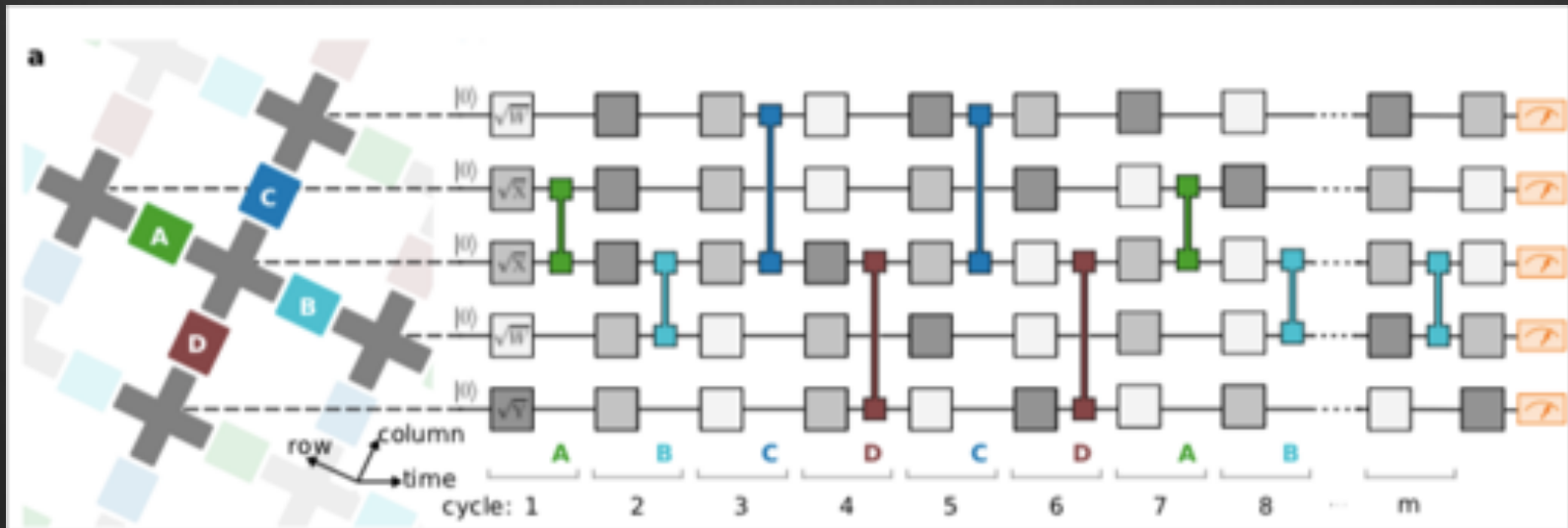
- noise rate

Today: what about **circuit depth**?

# Quantum supremacy using a programmable superconducting processor

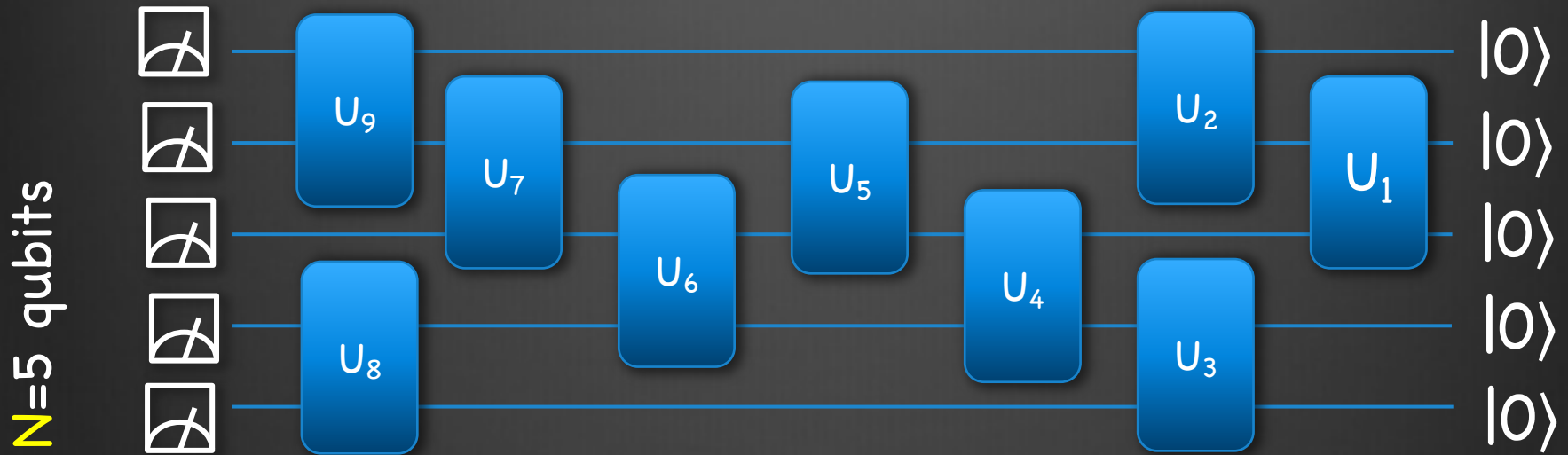
Google AI Quantum and collaborators<sup>†</sup>

task chosen to favor quantum computers and clear comparison





# quantum circuits

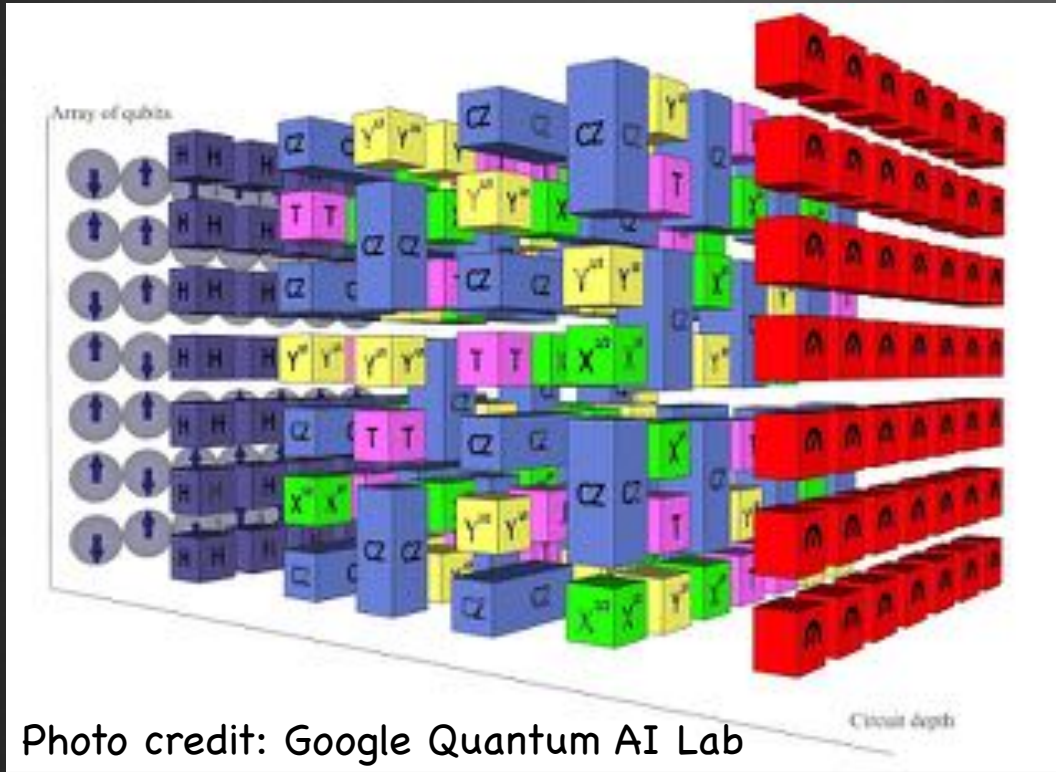


depth  $T=7$

$$p(z) = p(z_1, z_2, z_3, z_4, z_5) = |\langle z_1 z_2 z_3 z_4 z_5 | U_9 U_8 U_7 U_6 U_5 U_4 U_3 U_2 U_1 | 00000 \rangle|^2$$

Other parameters: connectivity, # of gates, fidelity.

# random circuit sampling



Conjecture:

Output distribution  $p(z)$  is hard to sample from on classical computer.

Google used  $N=53$  qubits in 2D geometry with  $T=20$ .

Conjecture:  $T \geq \sqrt{N} \rightarrow$  classical simulation time  $\exp(N)$ .

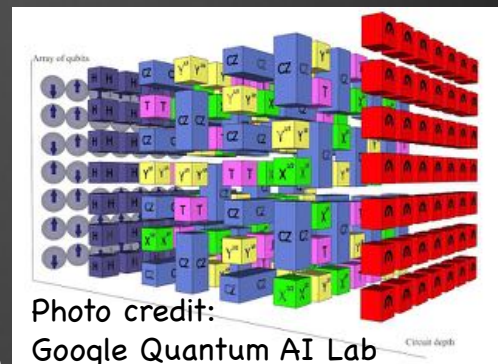
[Aaronson, Bremner, Jozsa, Montanaro, Shepherd, ...]

# low-depth circuits

Google proposal is  $\sqrt{N} \times \sqrt{N}$  grid for depth  $T \sim \sqrt{N}$ .

How low can we make depth?

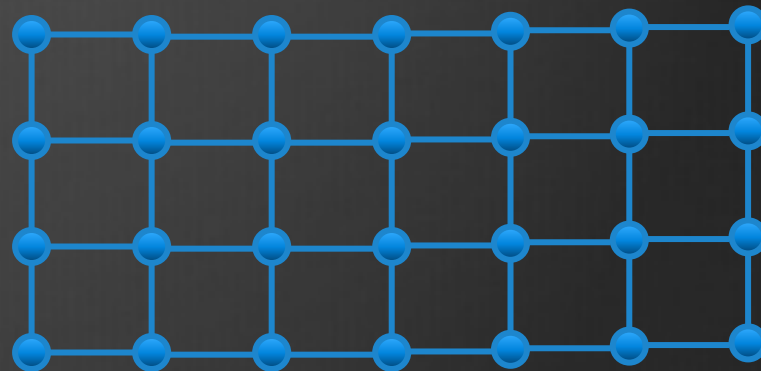
[Terhal-DiVincenzo '04] showed worst-case hardness of simulation as soon as  $T \geq 3$ . ( $T=2$  is easy.)



measurement-based  
quantum computing (MBQC)

W

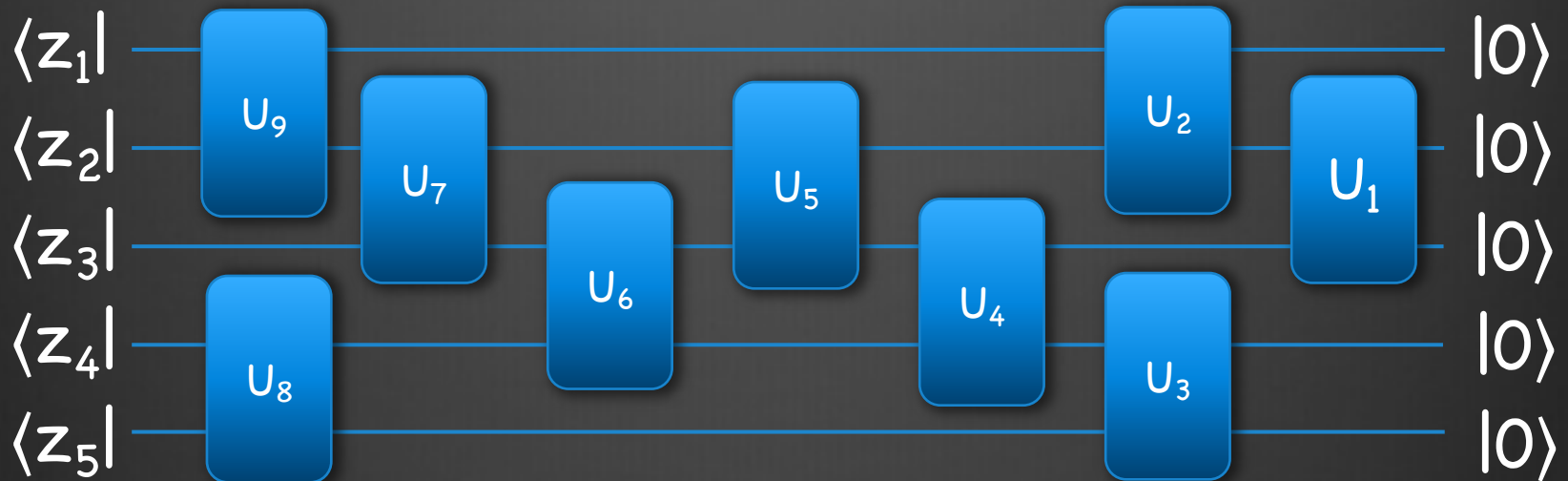
L



- prepare  $L \times W$  cluster state in  $O(1)$  depth
- single-qubit measurements simulate depth- $W$  circuit on line of  $L$  qubits
- implies classical hardness is  $\geq \exp(\min(L,W))$ .
- tensor contraction achieves this

# tensor contraction

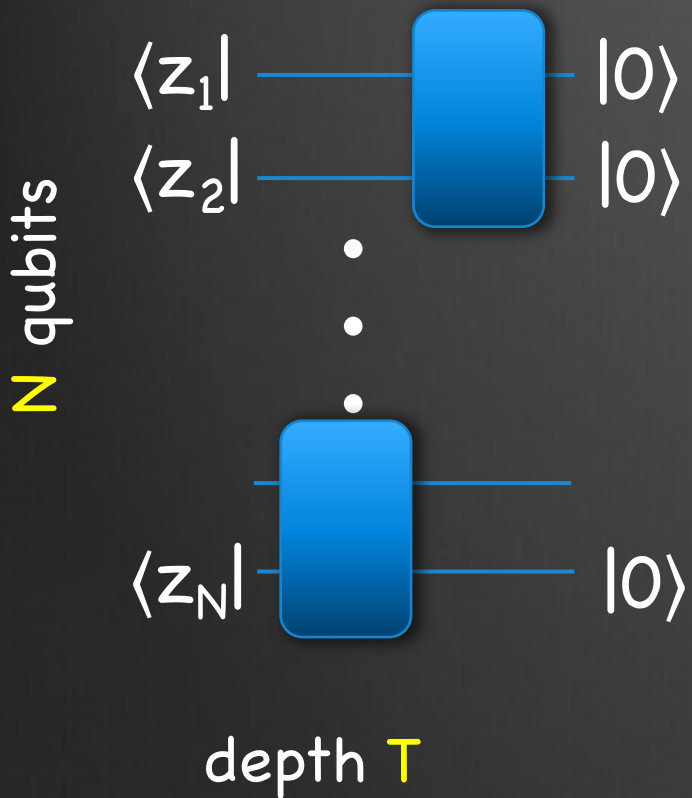
$$\langle z_1 z_2 z_3 z_4 z_5 | U_9 U_8 U_7 U_6 U_5 U_4 U_3 U_2 U_1 | 00000 \rangle =$$



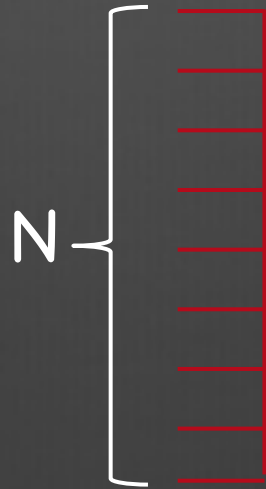
= tensor with 4 indices, each dim 2



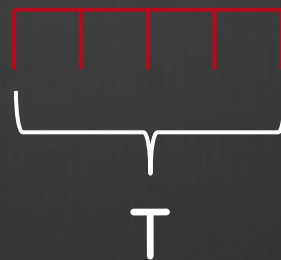
# tensor contraction in 1-D



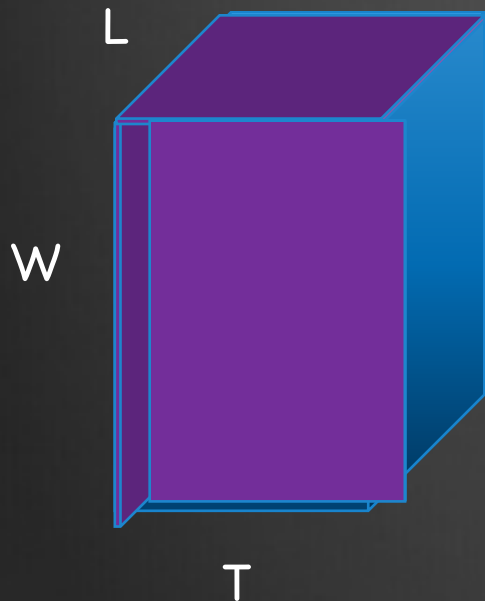
intermediate  
tensors can be



run time:  
 $T \exp(N)$  or  
 $N \exp(T)$



# simulating 2-D circuits



Depth  $T=O(1)$  circuit  
on  $\sqrt{N} \times \sqrt{N}$  grid



Naively takes time  $2^{O(\sqrt{N})}$   
 $\approx \sqrt{N}$  qubits on line for time  $\sqrt{N}$ .

But 1-D effective evolution is  
not unitary.

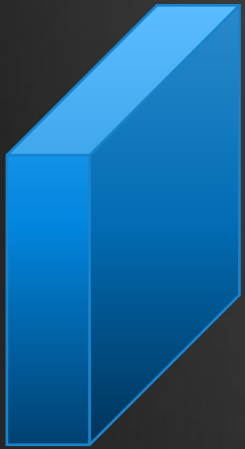
Entanglement has phase transition  
from area law  $\rightarrow$  volume law.

can be simulated in  
time  $2^{LW}$  or  $2^{LT}$  or  $2^{WT}$

$T=3$  in area law phase  $\rightarrow$   
 $N^{O(1)}$ -time classical simulation for  
**approximate** sampling of **random** circuits.  
Exact or worst-case is #P-hard.

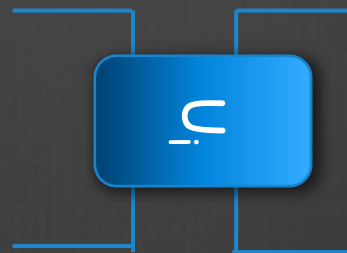
# cheaper tensor contraction

$T=O(1)$   
 $\sqrt{N} \times \sqrt{N}$  grid



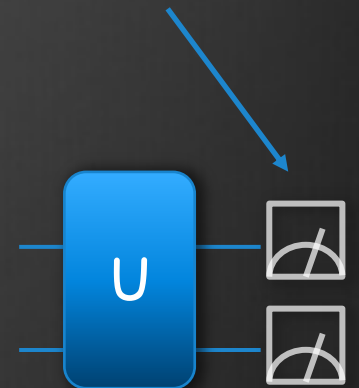
$\sqrt{N}$  qubits  
evolving for  
 $\sqrt{N}$  time

effective  
time  
evolution

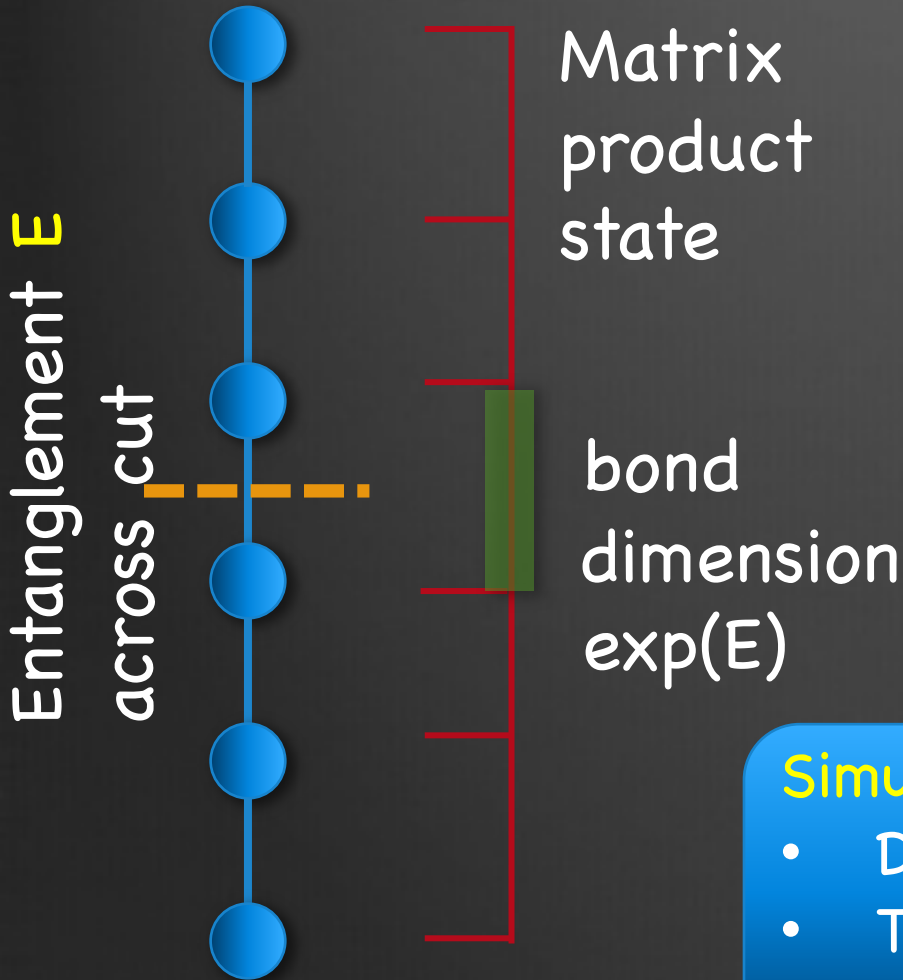


sideways gates  
generically not  
unitary

weak  
measurement



# Approximate simulation



Example:



$$\begin{aligned}
 |\psi\rangle &= \sum_{i,j} c_{i,j} |i\rangle \otimes |j\rangle \\
 &= \sum_k \lambda_k |\alpha_k\rangle |\beta_k\rangle
 \end{aligned}$$

$$\begin{aligned}
 &\sum_k |k\rangle |\alpha_k\rangle \\
 &\left[ \sum_k \lambda_k \langle k | \langle k | \right. \\
 &\left. \sum_k |k\rangle |\beta_k\rangle \right]
 \end{aligned}$$

- Simulation algorithm:**
- Do tensor contraction
  - Truncate bonds to dim  $\exp(O(E))$ .

Run-time is  $N2^{O(E)}$ .



# Does the algorithm work?

"Beware of bugs in the above code; I have only proved it correct, not tried it."  
--Donald Knuth

## 1. Yes.

We tested it and simulated 400x400 grids on a laptop.

## 2. Probably.

We proved a phase transition in something like the effective entanglement.

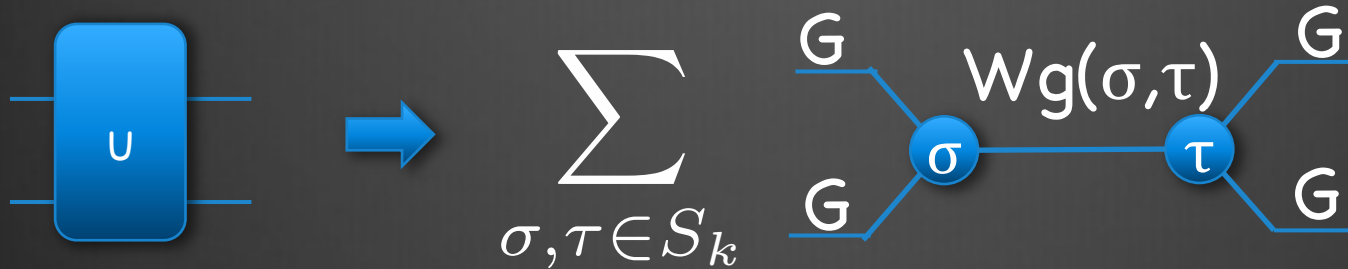
## 3. Sometimes.

The extended brickwork architecture is #P-hard to simulate exactly but our algorithm is proven to work on it.

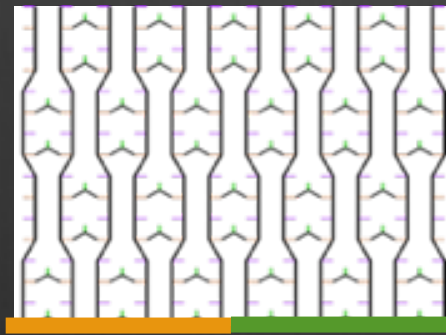
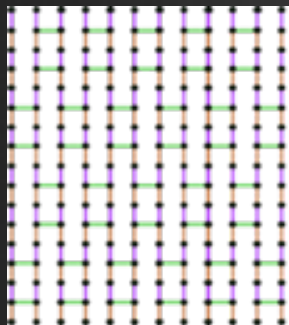
# stat mech model

Qubits  $\rightarrow$  dim- $q$  particles.  
 $E[\text{tr}[\rho_A^k]] =$  partition function

$$\begin{aligned} Wg^{-1}(\sigma, \tau) &= G(\sigma, \tau) \\ &= q^{-\text{dist}(\sigma, \tau)} \end{aligned}$$

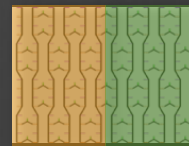


$k=2 \leftrightarrow$  anisotropic Ising

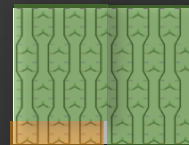


A  $A^c$

ordered



disordered



volume law

$$E = O(\sqrt{N})$$

$\gg 1 \sim$  holography

$$q_c = 3.249\dots$$

$1+\epsilon \sim$  high-T Potts

area law

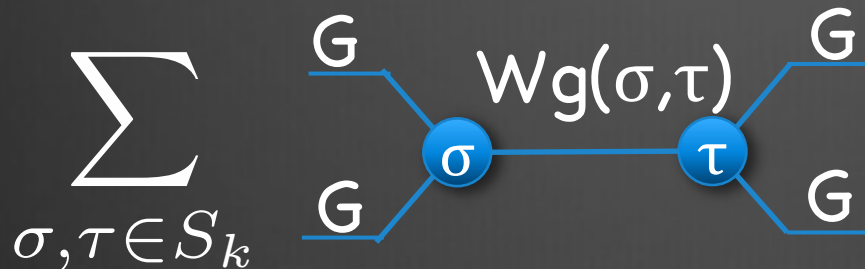
$$E = O(1)$$

$q$

conjecture: decimating yields nonnegative weights for all  $k, q$

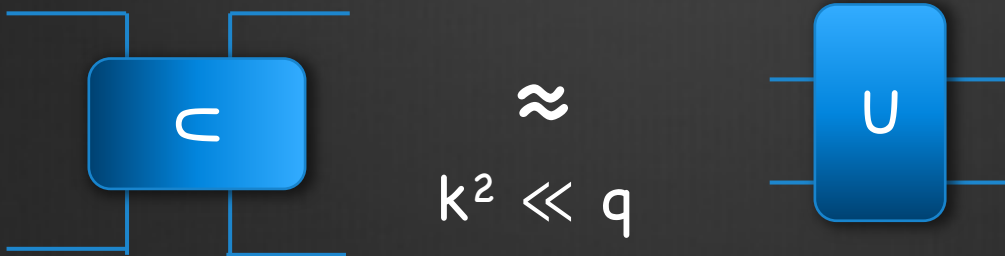
# random tensor networks

$q = \text{local dim, } E[\text{tr}[\rho^k]]$



$$\begin{aligned} Wg^{-1}(\sigma, \tau) &= G(\sigma, \tau) \\ &= q^{-\text{dist}(\sigma, \tau)} \end{aligned}$$

- $E_{U \sim \text{Haar}}[U^{\otimes k} \otimes (U^*)^{\otimes k}] = \sum_{\sigma, \tau} Wg(\sigma, \tau) |\sigma\rangle\langle\tau|$
- $E_{G \sim \text{GUE}}[G^{\otimes k} \otimes (G^*)^{\otimes k}] = \sum_{\sigma} |\sigma\rangle\langle\sigma|$
- $\langle\sigma|\tau\rangle = G(\sigma, \tau)$
- $\|G - I\|_{\text{op}}, \|Wg - I\|_{\text{op}} \leq k^2 / q$



# Open questions

- Rigorously prove the location of the phase transition and the correctness of the algorithm.
- Random tensor networks with low bond dimension
- Universality classes in random circuits?
- (time-independent) Hamiltonian versions?
- Where exactly is the boundary between easy and hard?



# Thanks!



Fernando  
Brandão



Alex  
Dalzell



Rolando  
La Placa



John Napp

arXiv:2001.00021