

the talk

Daniel Nagaj



universität
wien

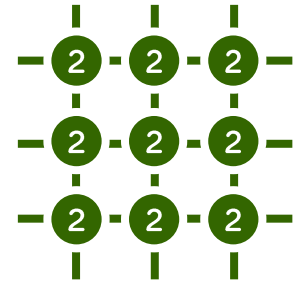


2014 | 2 | 24
Simons Institute

1

strong promises

and eigenvalue gaps



2

the history

of the history state



3

running the clock

precise/faulty, qubit/qudit, sequential/parallel



4

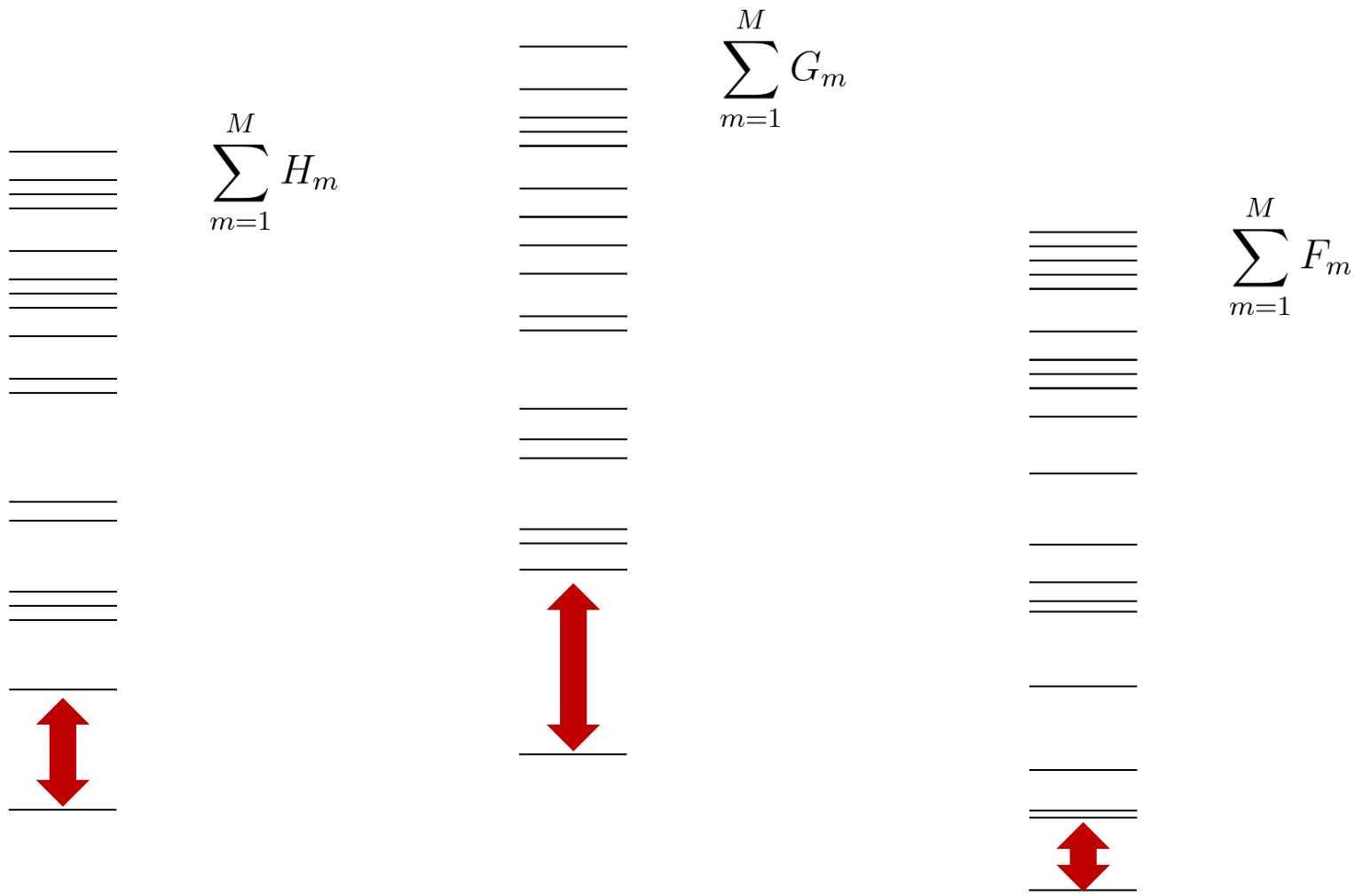
on the qPCP road

questions & warnings



1

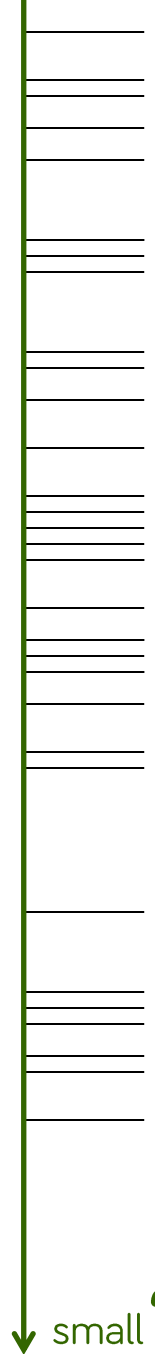
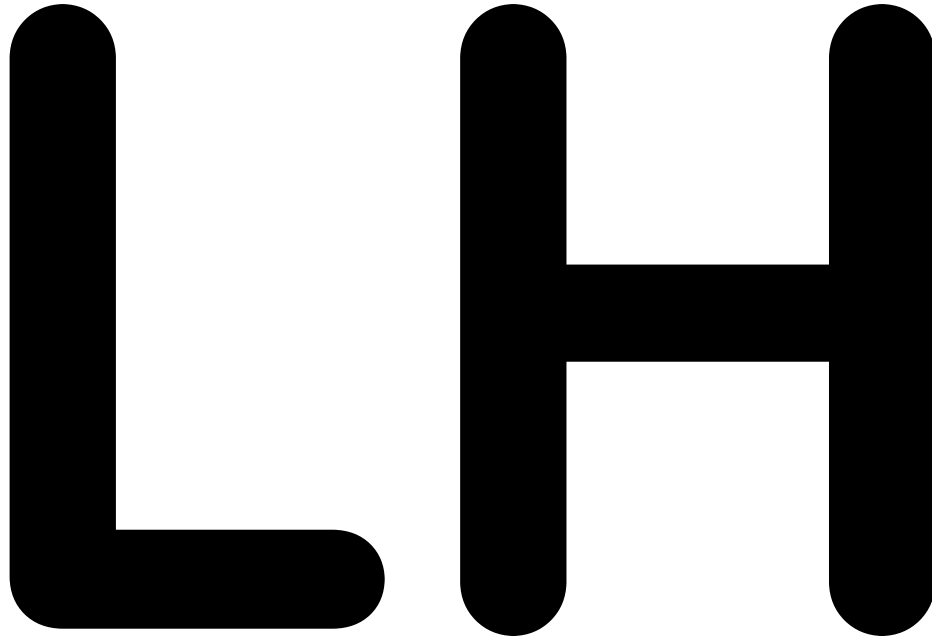
Hamiltonians and their eigenvalue gaps



1

Hamiltonians and their ground states

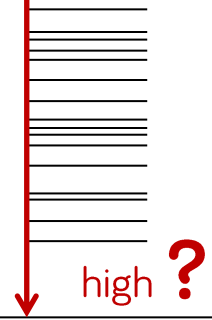
Is
the
ground
state
energy
of a



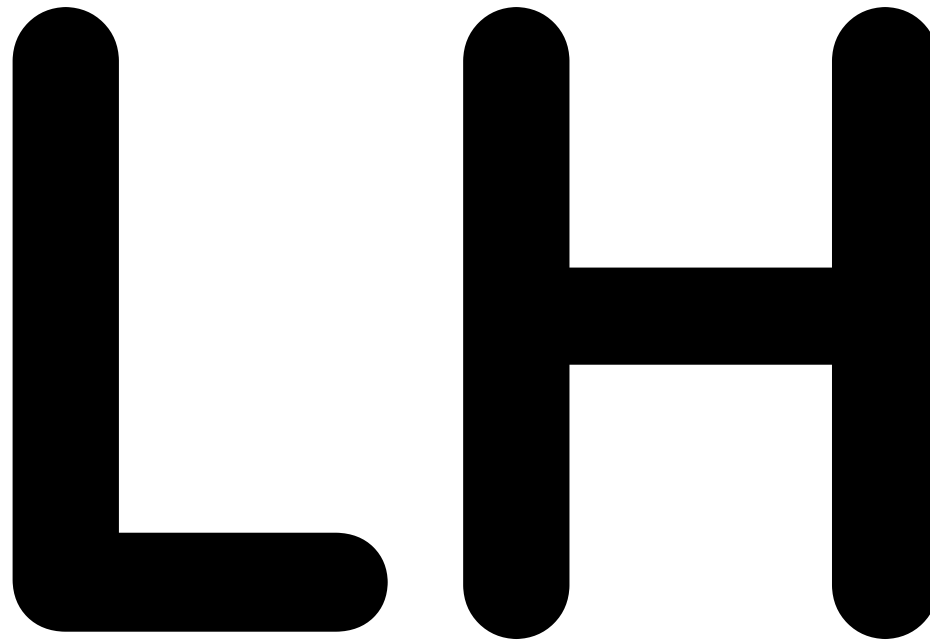
small ?

1

Hamiltonians and their ground states

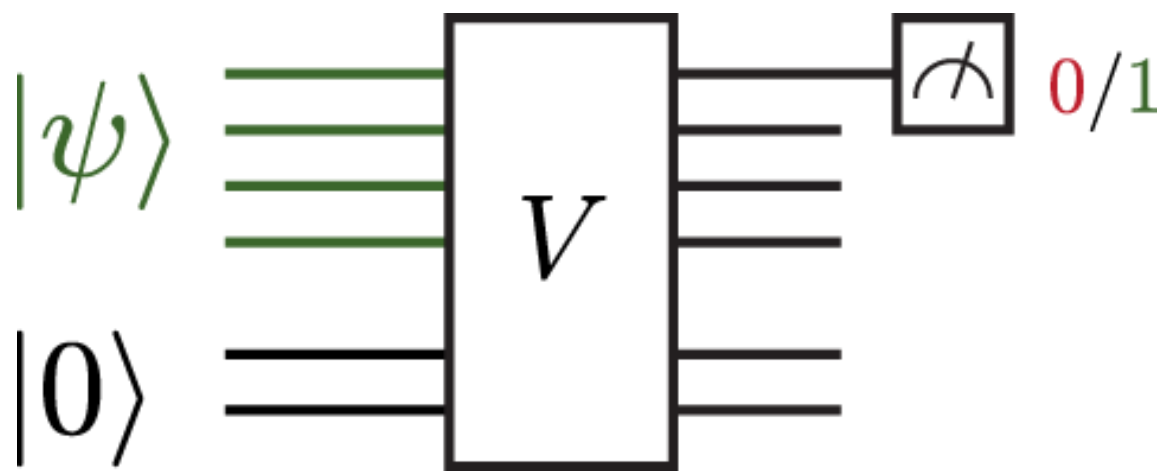
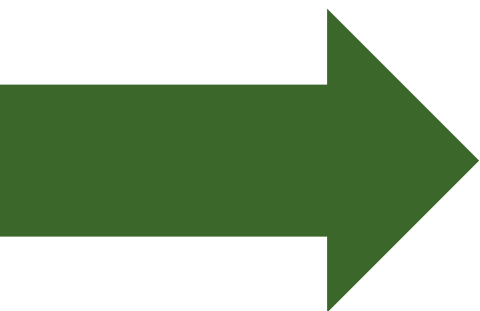
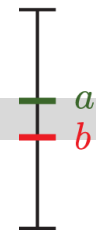


Is
the
ground
state
energy
of a



1 The QMA protocol

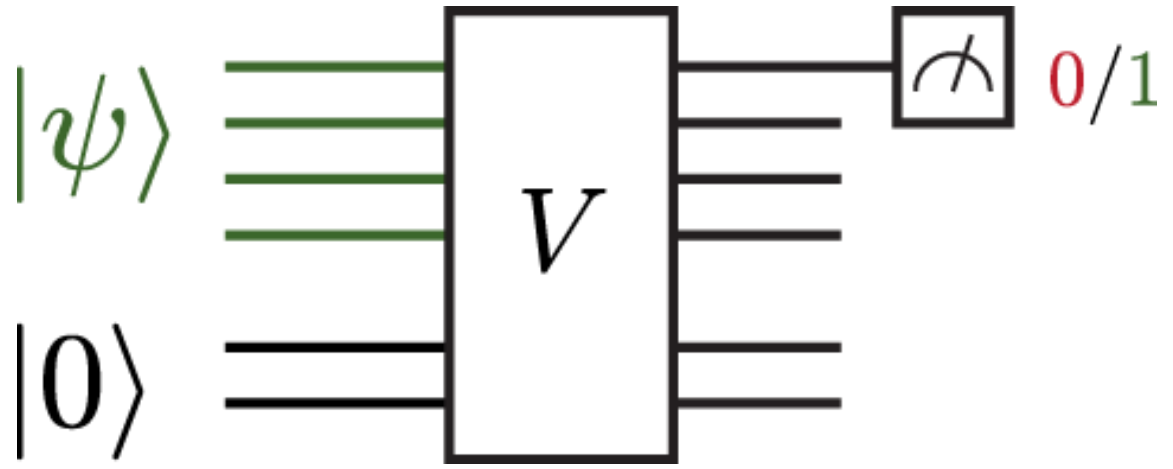
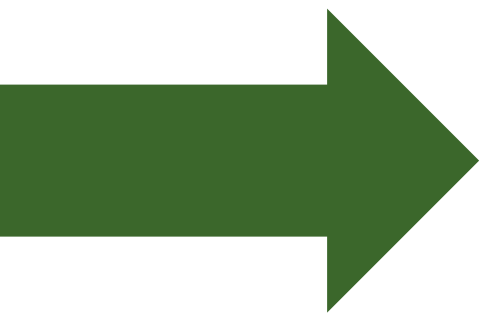
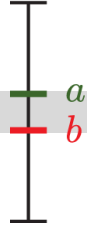
YES? Accept a good proof with $p > a$.
the promise
NO? Probability of accepting $p < b$.



- Is there an acceptable witness for this circuit?
- Is some local Hamiltonian (nearly) frustration-free?

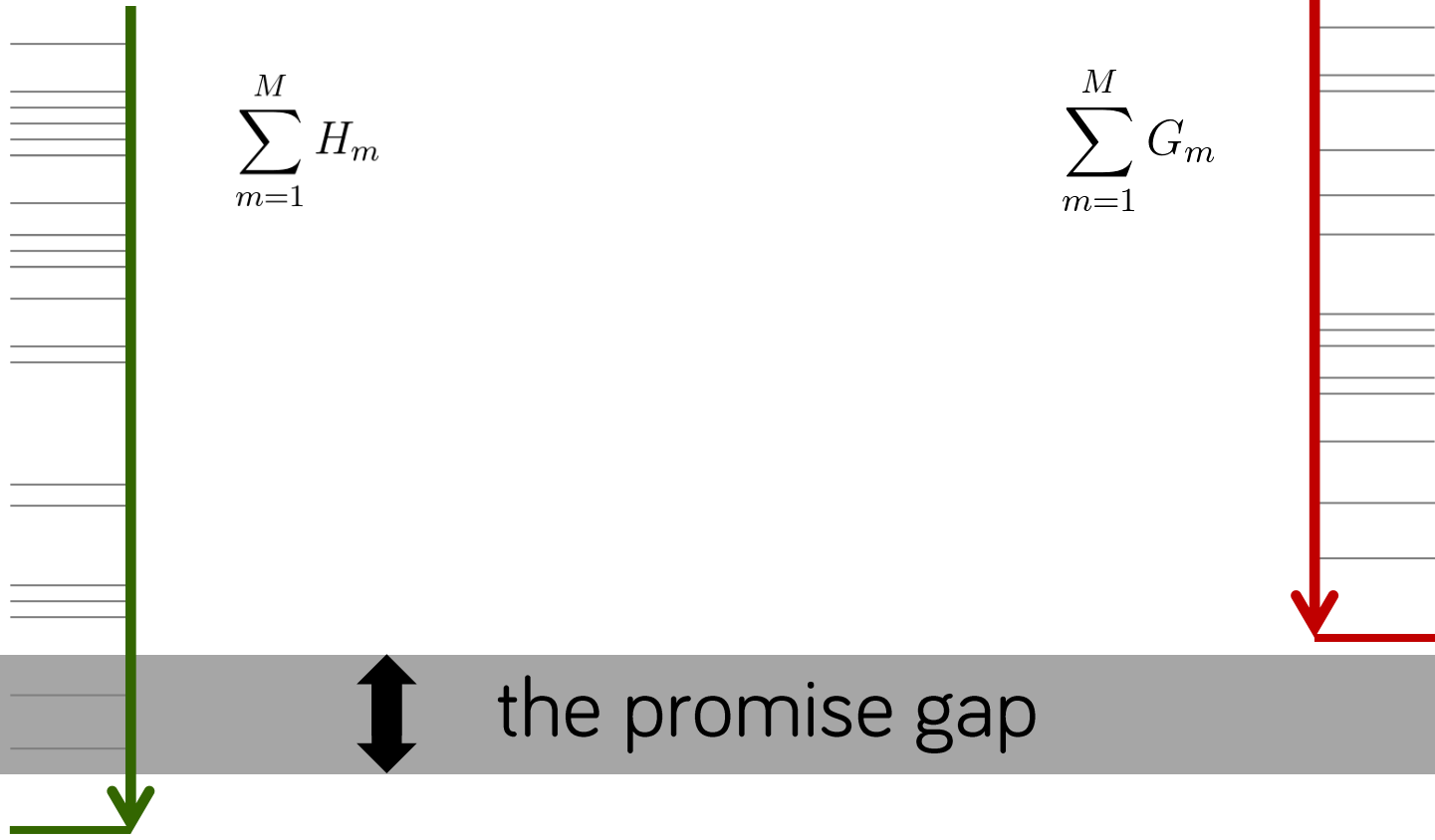
1 The QMA protocol

YES? Accept a good proof with $p > a$.
the promise
NO? Probability of accepting $p < b$.



- Is there an acceptable witness for this circuit?
- Does some local Hamiltonian have a low ground energy?

1 The promise gap for a problem



YES

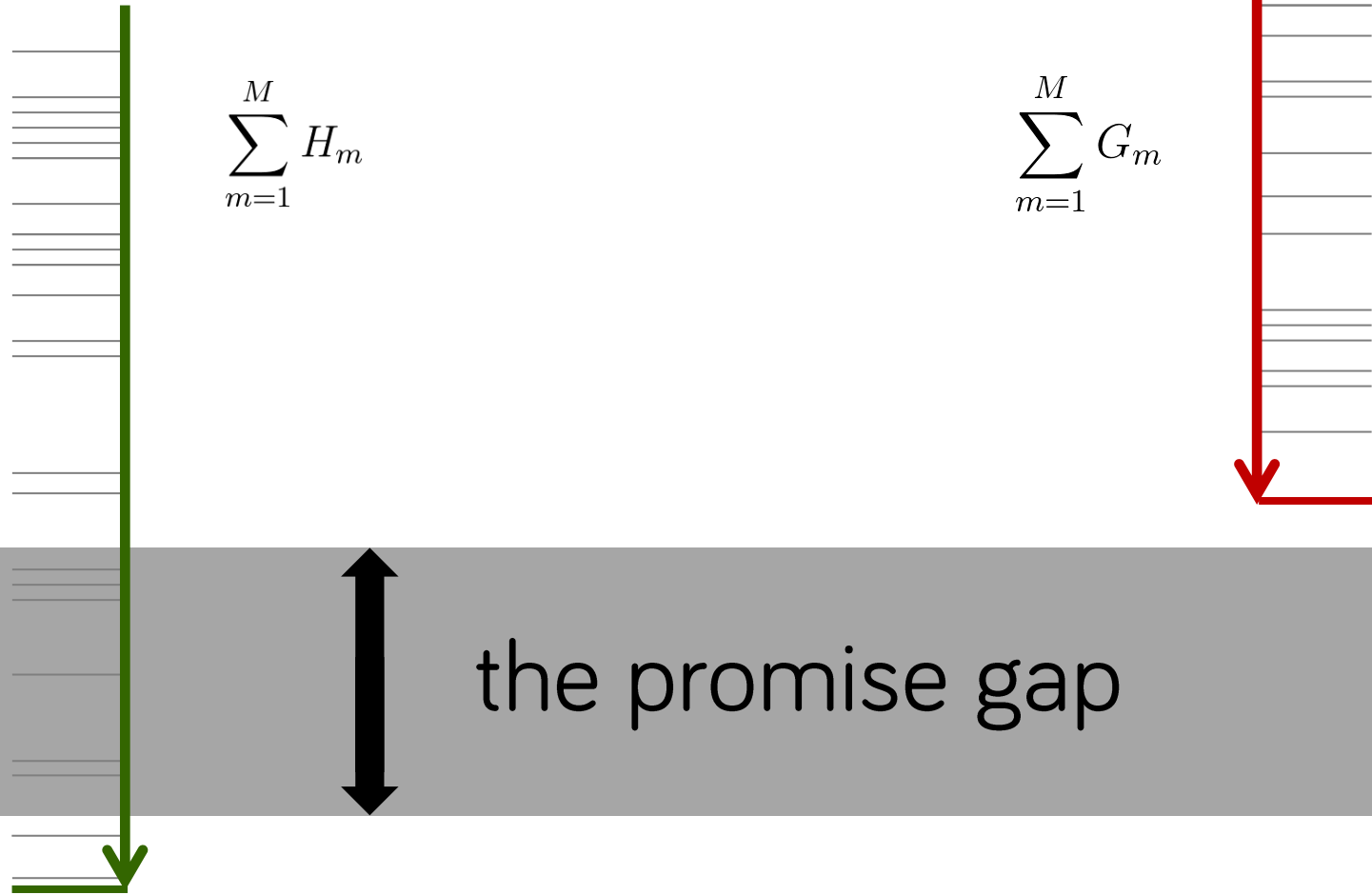
little frustration

lots of frustration

NO

1

The promise gap for a simpler problem?

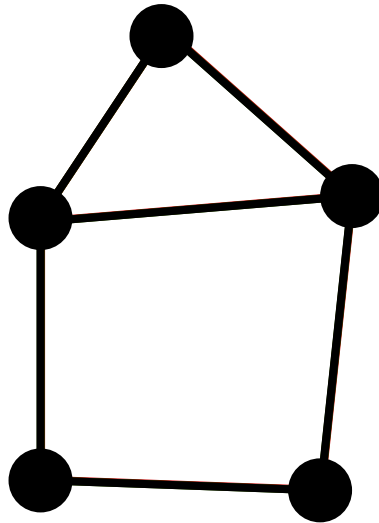
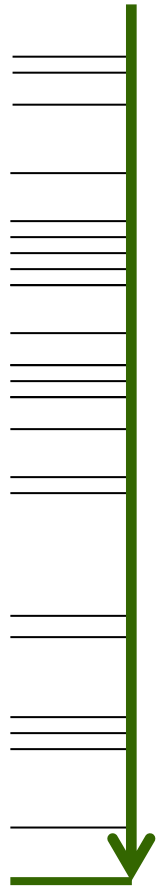


YES

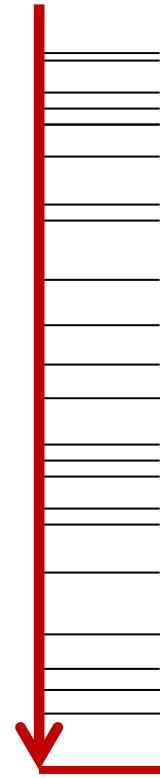
little
frustration

lots of
frustration

NO



$$\sum_{m=1}^M H_m$$

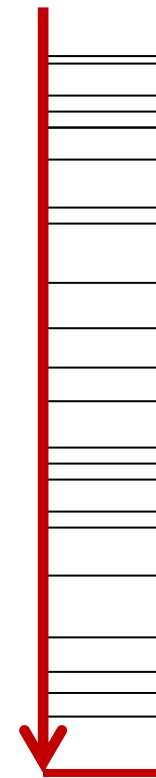
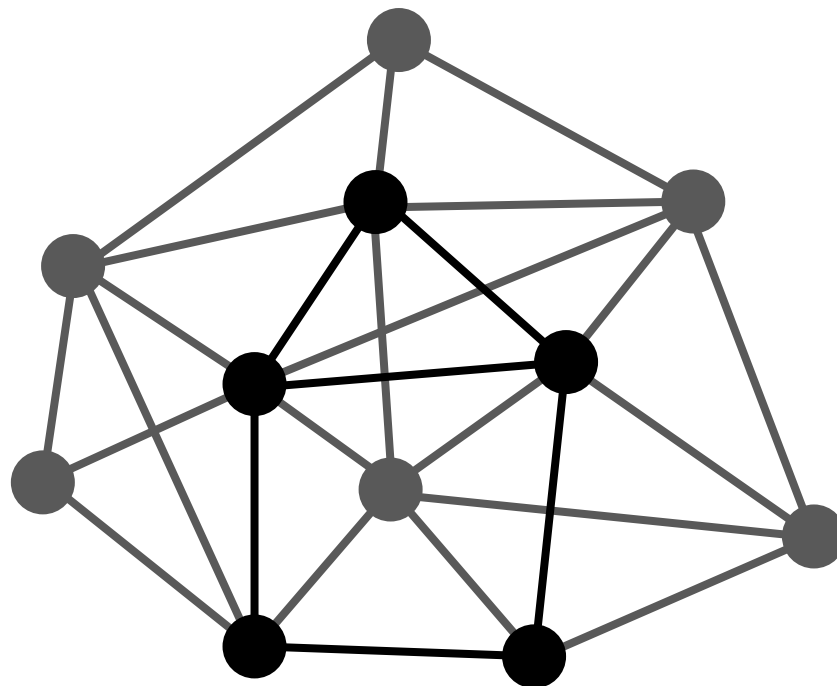
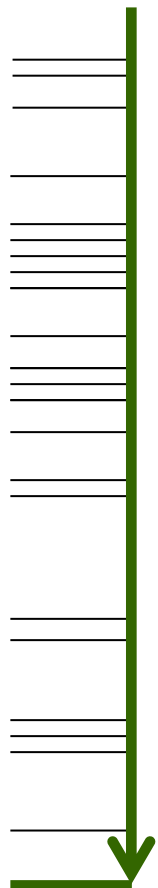


YES

little
frustration

lots of
frustration

NO



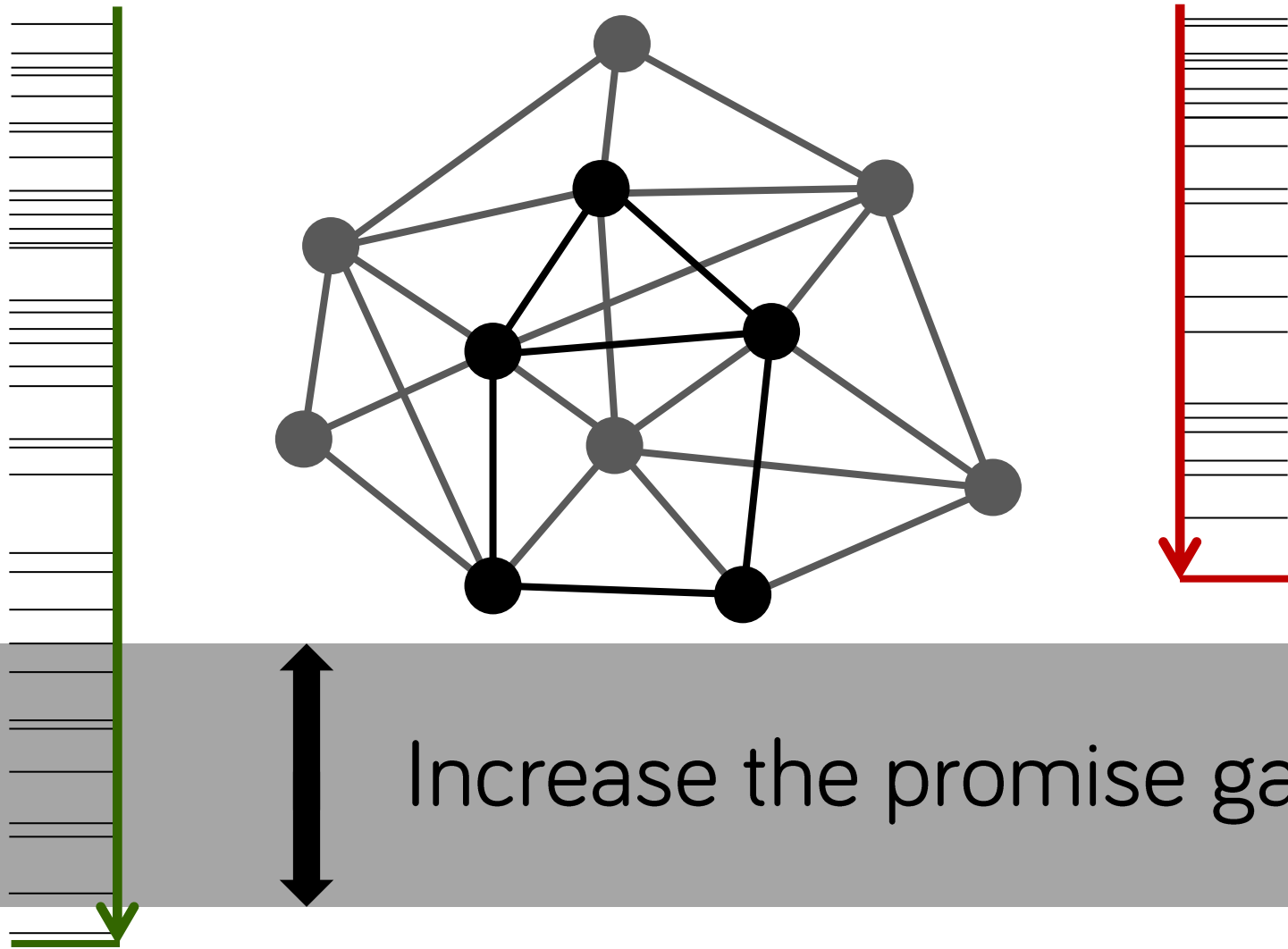
$$\sum_{m=1}^M H_m \rightarrow \sum_{m=1}^{M'} H'_m$$

YES

little
frustration

lots of
frustration

NO

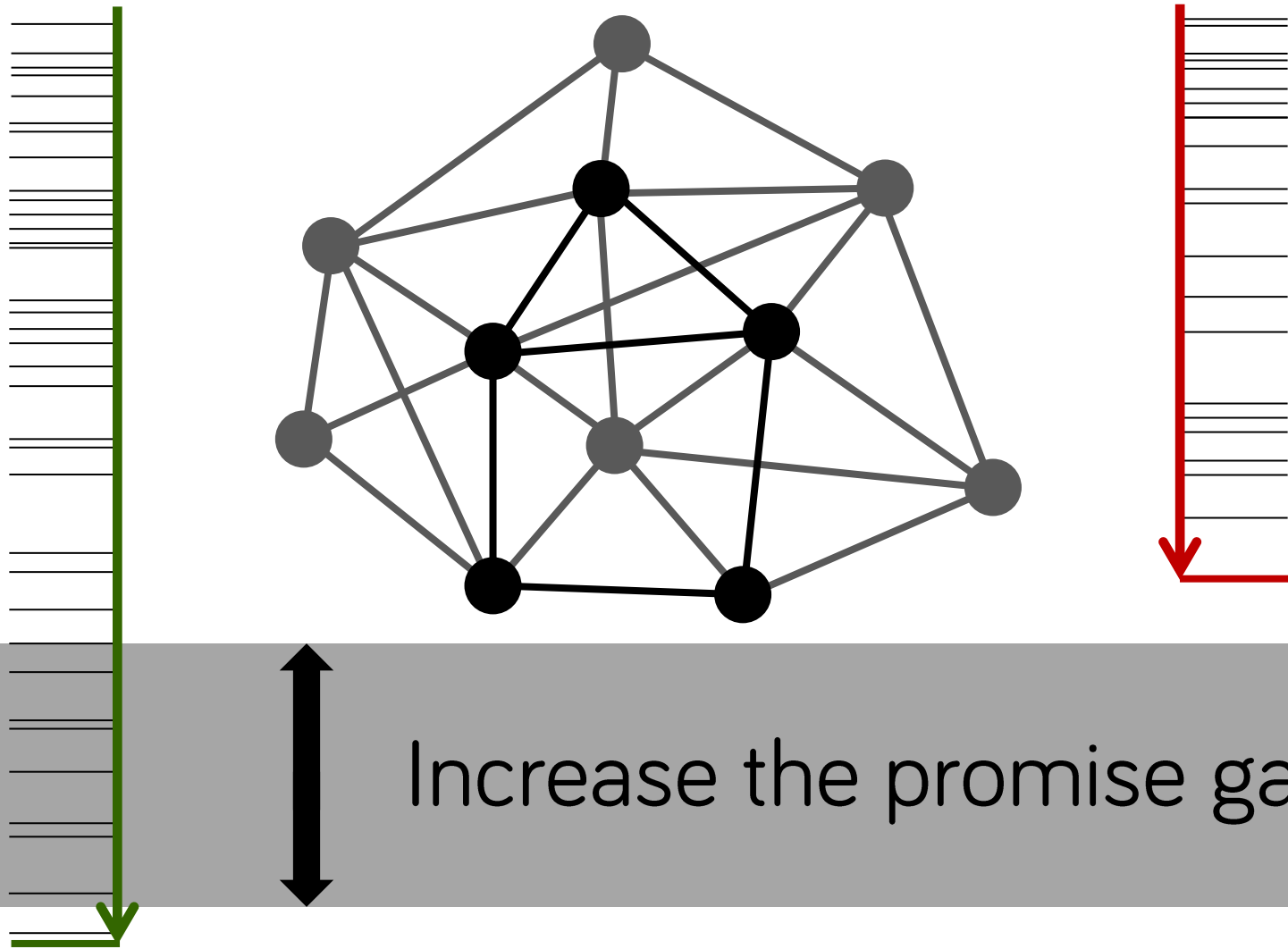


YES

little
frustration

lots of
frustration

NO



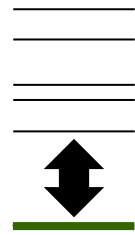
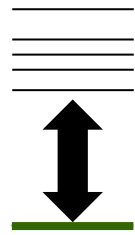
Increase the promise gap?

Using the usual circuit encoding
clock construction based ideas?

NO

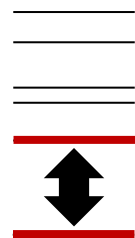
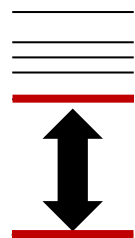
1

Small eigenvalue gaps ... small promise gaps



YES little frustration
a very low ground energy

$$H = A + B$$



NO without much frustration
a pretty low ground energy

■ geometric lemma

$$\lambda_0 \geq \sin^2 \frac{\vartheta}{2} \times \min(\Delta_A, \Delta_B)$$

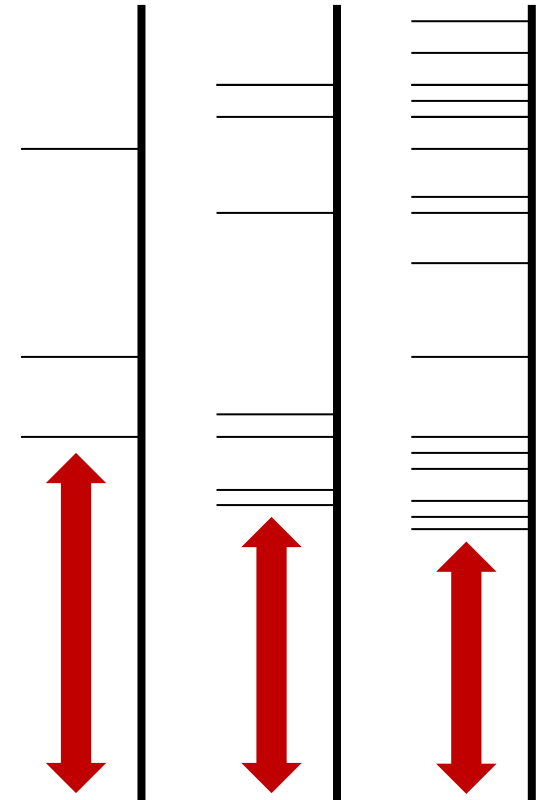
1 Small or large eigenvalue gaps?

$$N \rightarrow \infty$$

- Anything close to the ground state?

constant gap? $\Delta \geq \text{const.}$

1D: area law, an algorithm
2D: area law?



a Heisenberg XXX spin-1 chain (AKLT)

$$\sum_{j=1}^{N-1} X_j X_{j+1} + Y_j Y_{j+1} + Z_j Z_{j+1}$$

a biased walk in 1D

$$\sum_{j=1}^{N-1} (|j\rangle - B|j+1\rangle) (\langle j| - B\langle j+1|)$$

1 Small or large eigenvalue gaps?

$$N \rightarrow \infty$$

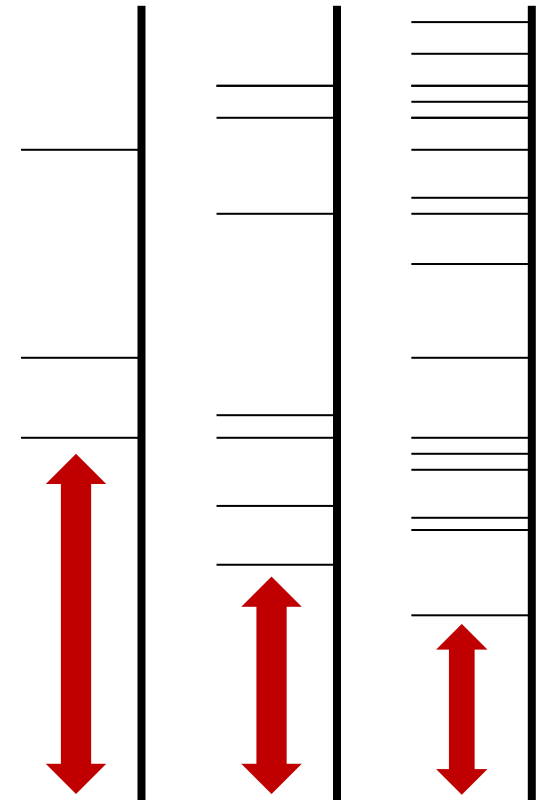
- Anything close to the ground state?

constant gap? $\Delta \geq \text{const.}$

1D: area law, an algorithm
2D: area law?

inverse-poly gap? $\Delta \propto N^{-c} \rightarrow 0$

clock constructions
NP, QCMA hard
qubits? 1D?



transverse-field Ising

$$\sum_{j=1}^{N-1} X_j - \sum_{j=1}^{N-1} Z_j Z_{j+1}$$

quantum walk on a line

$$\sum_{j=1}^{N-1} |j\rangle\langle j+1| + |j+1\rangle\langle j|$$

1 Small or large eigenvalue gaps?

$$N \rightarrow \infty$$

- Anything close to the ground state?

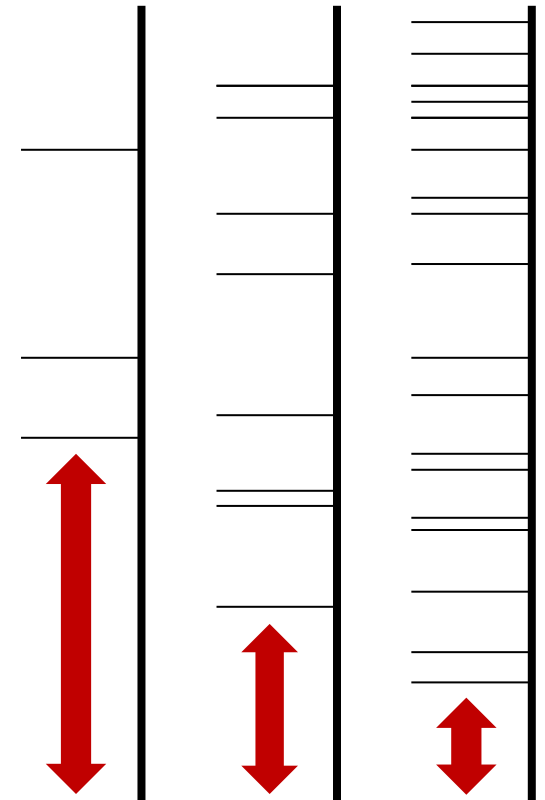
constant gap? $\Delta \geq \text{const.}$

1D: area law, an algorithm
2D: area law?

inverse-poly gap? $\Delta \propto N^{-c} \rightarrow 0$

clock constructions
NP, QCMA hard
qubits? 1D?

exp-small gap? $\Delta \propto 2^{-cN} \rightarrow 0$

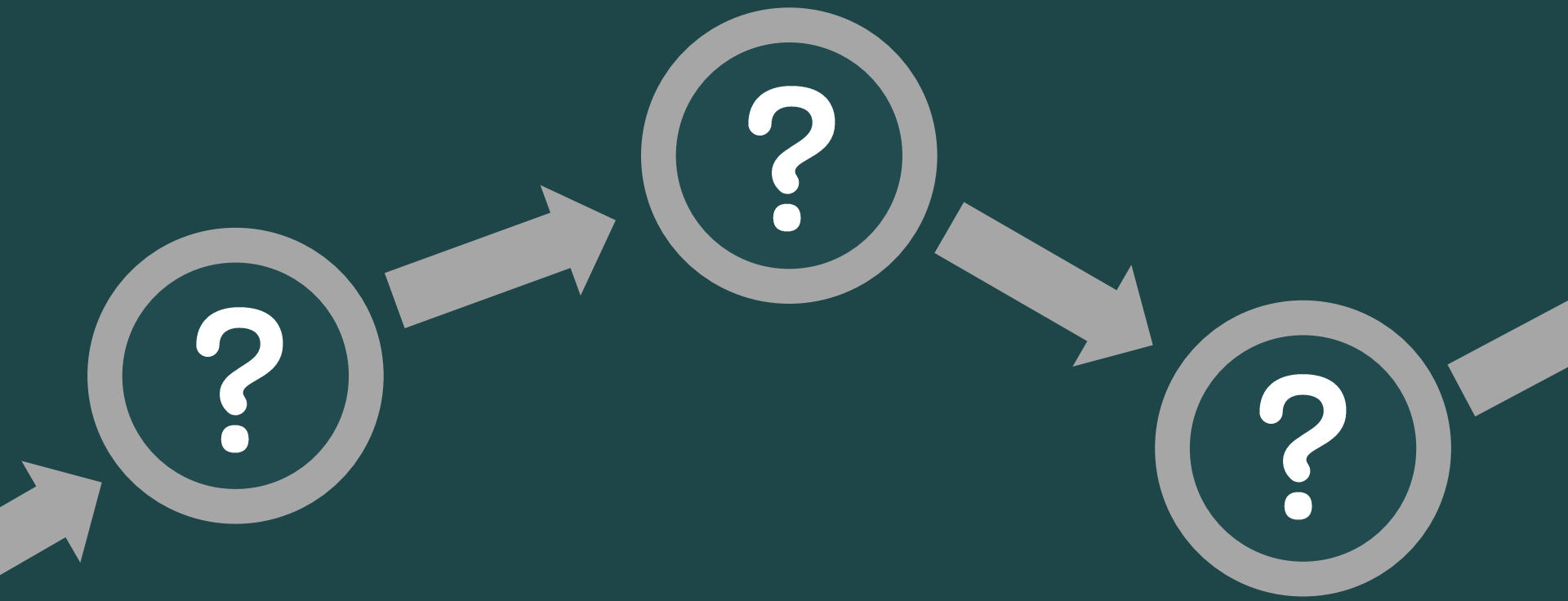


- Constant degree LH: at most constant gap.



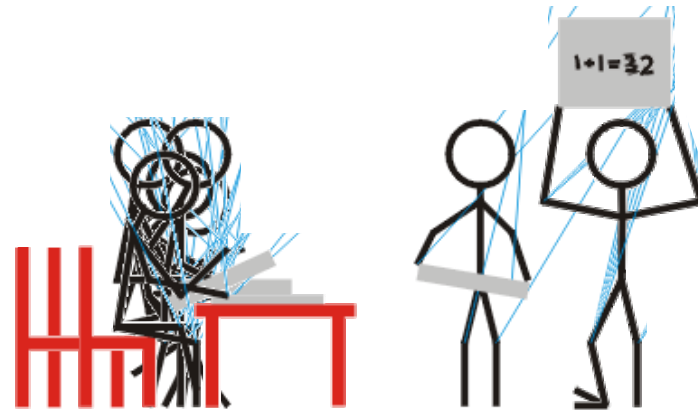
- Degeneracy: help or trouble?



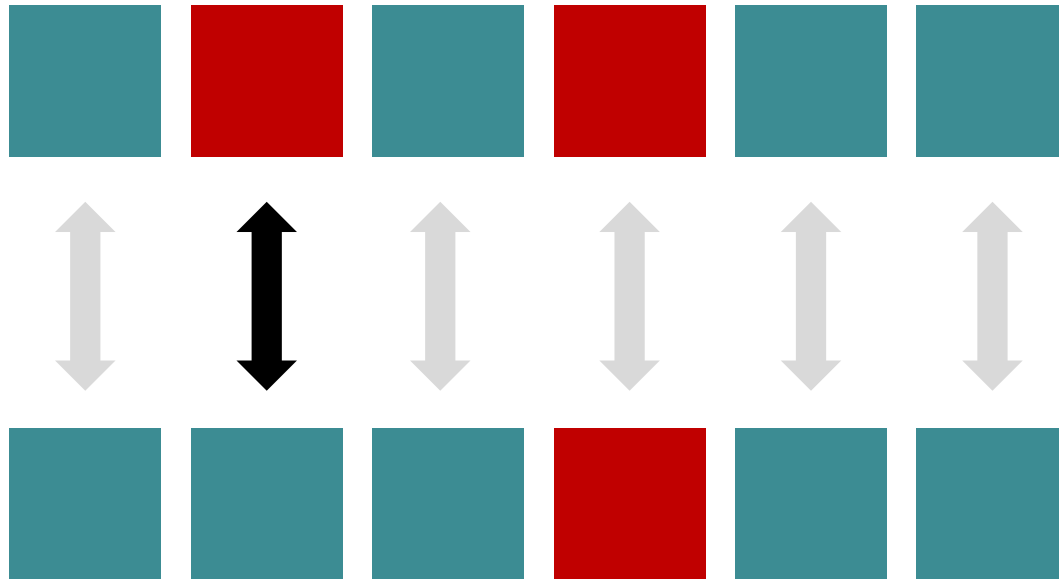


the history state
ground

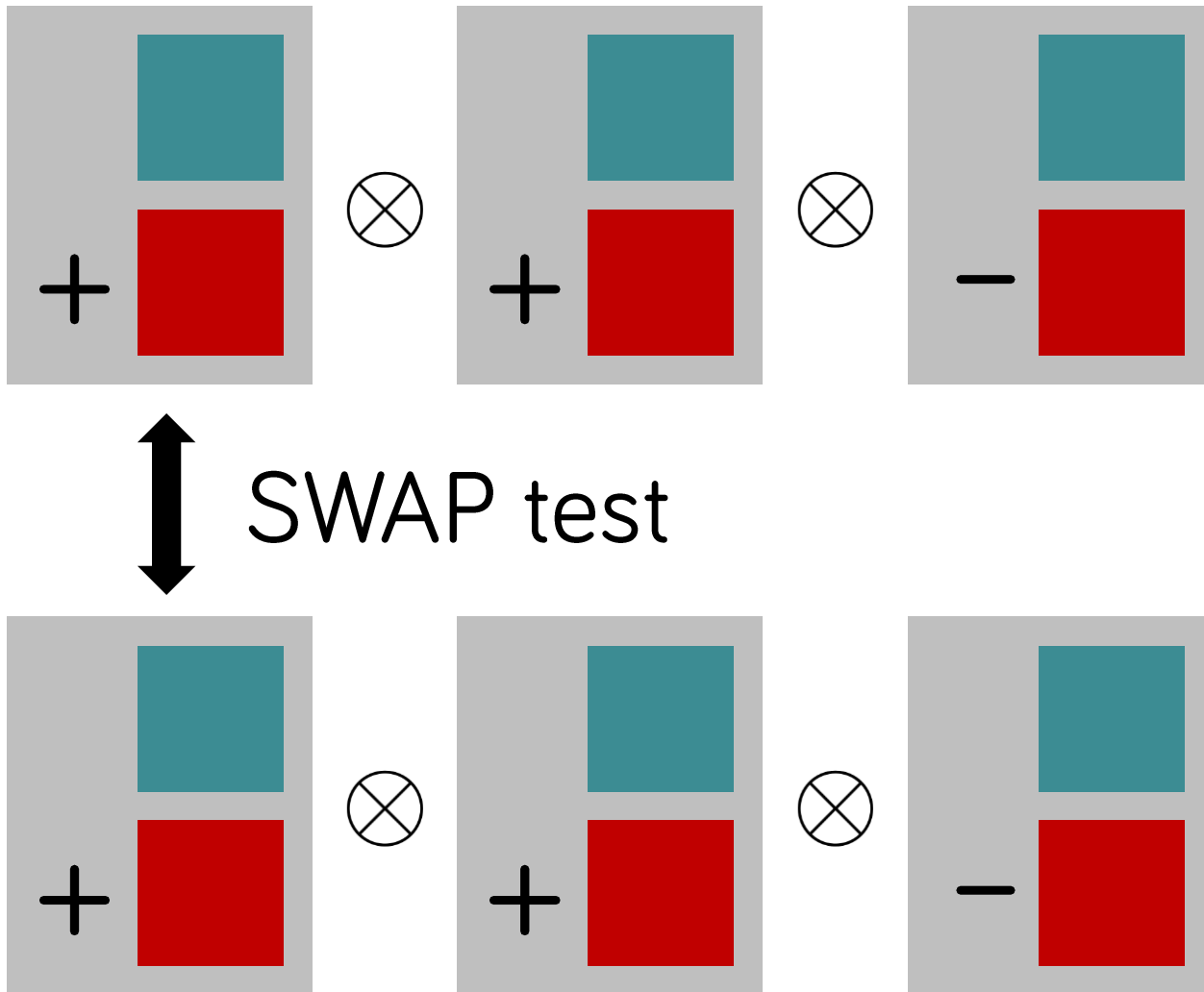
2 Snapshots of a computation



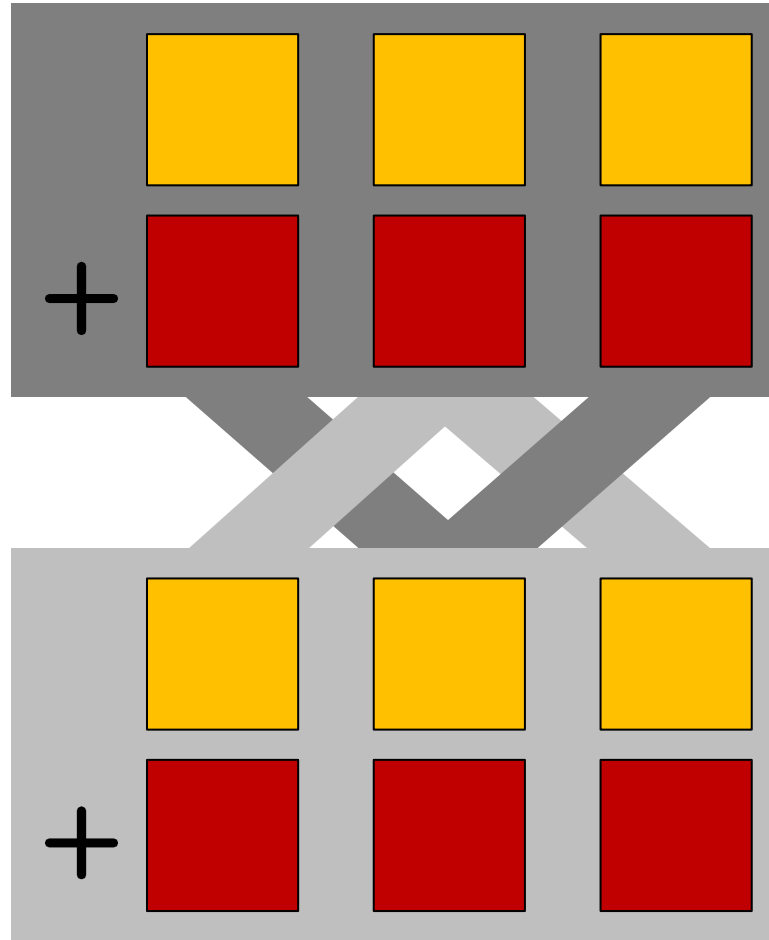
Locally comparing **strings**.



Locally comparing products.

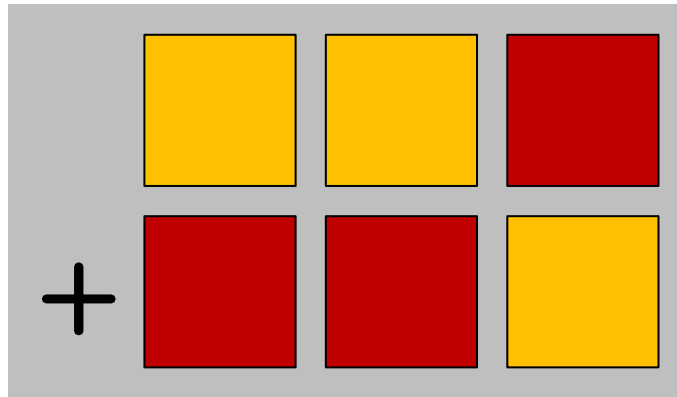


Locally comparing **entangled** states?



UGH!

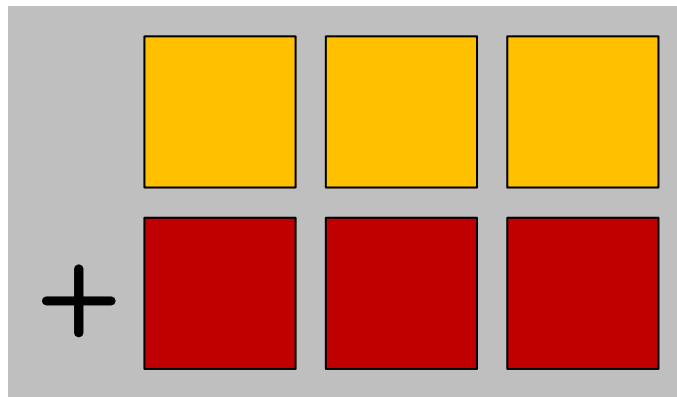
2 Labeling the data



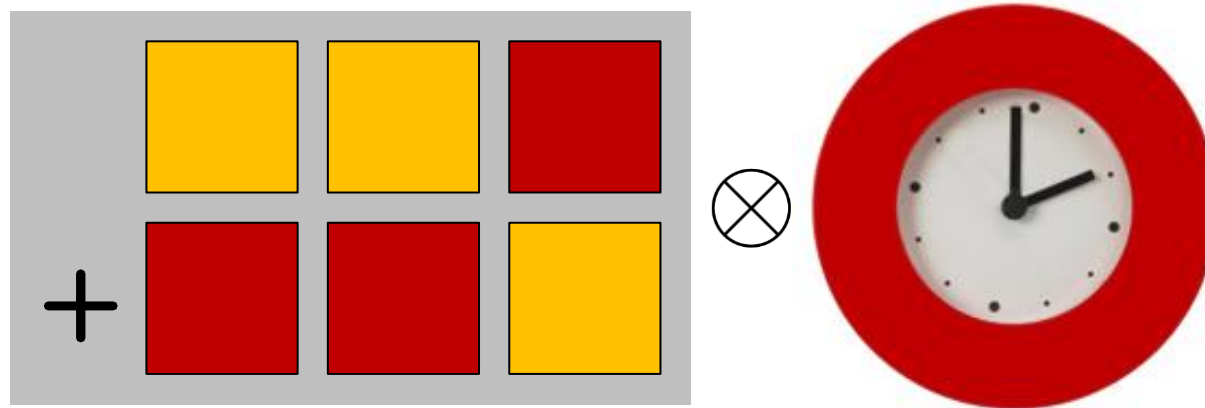
Hard to compare directly (locally).

U^\dagger   U

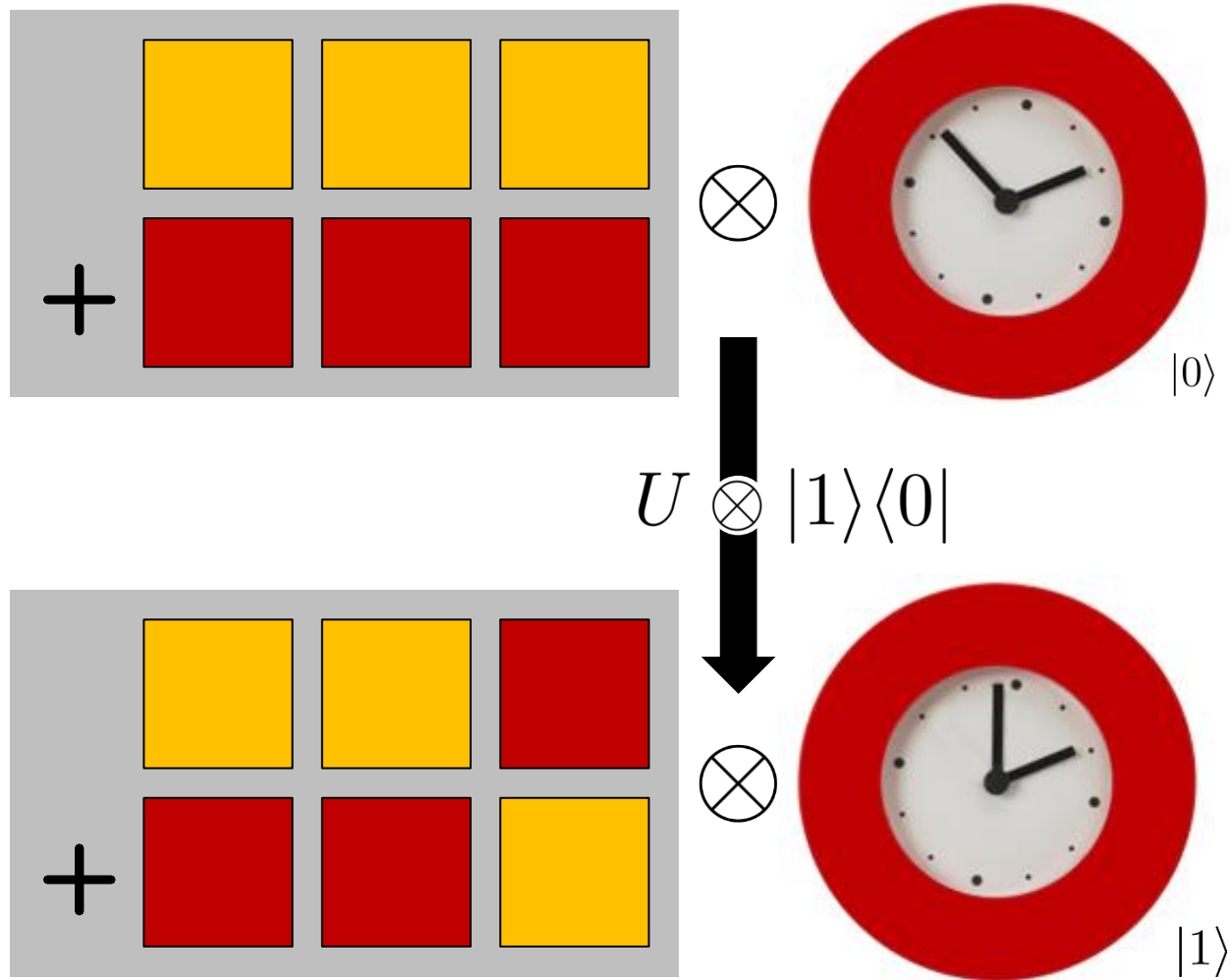
a clock



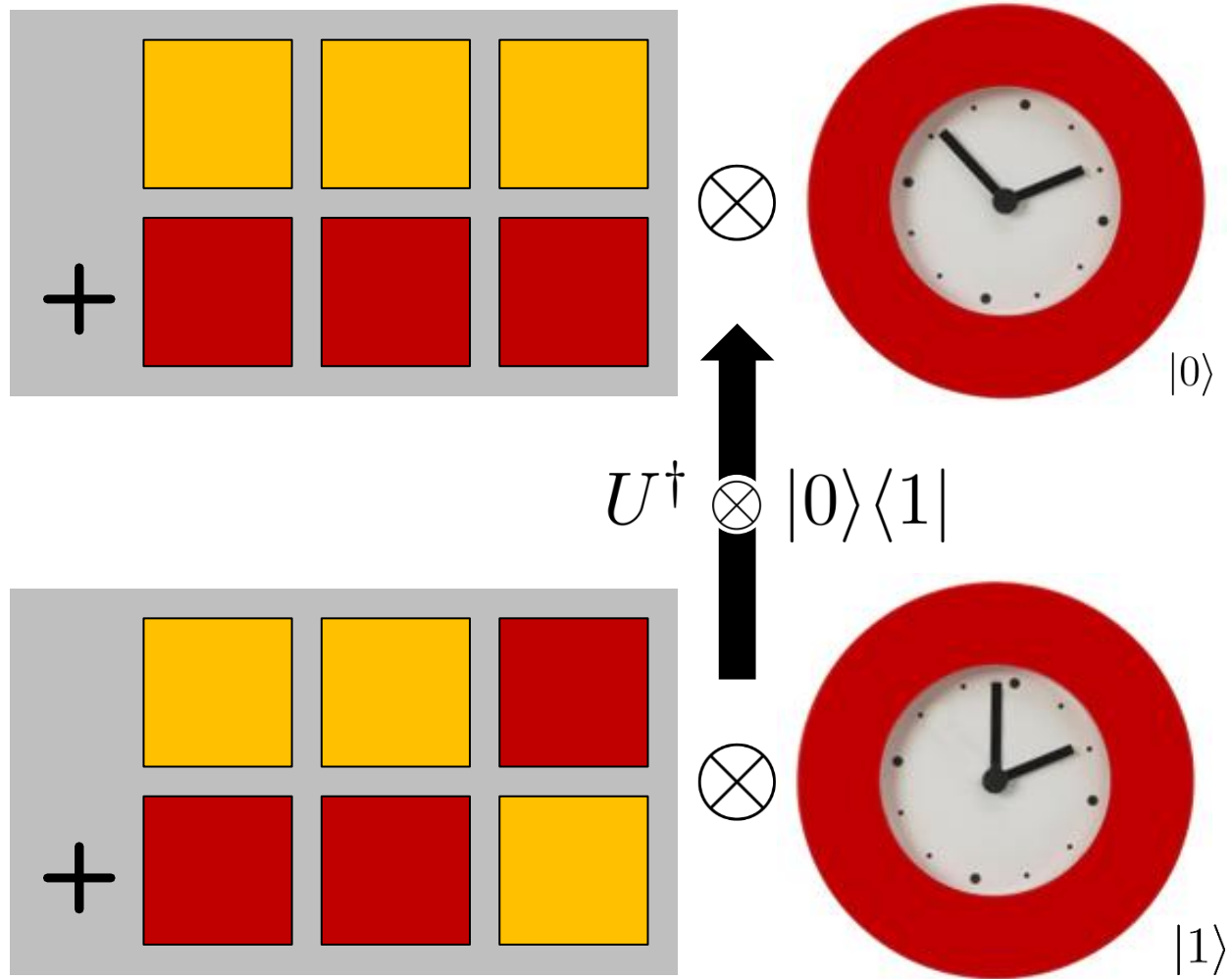
2 Labeling the data



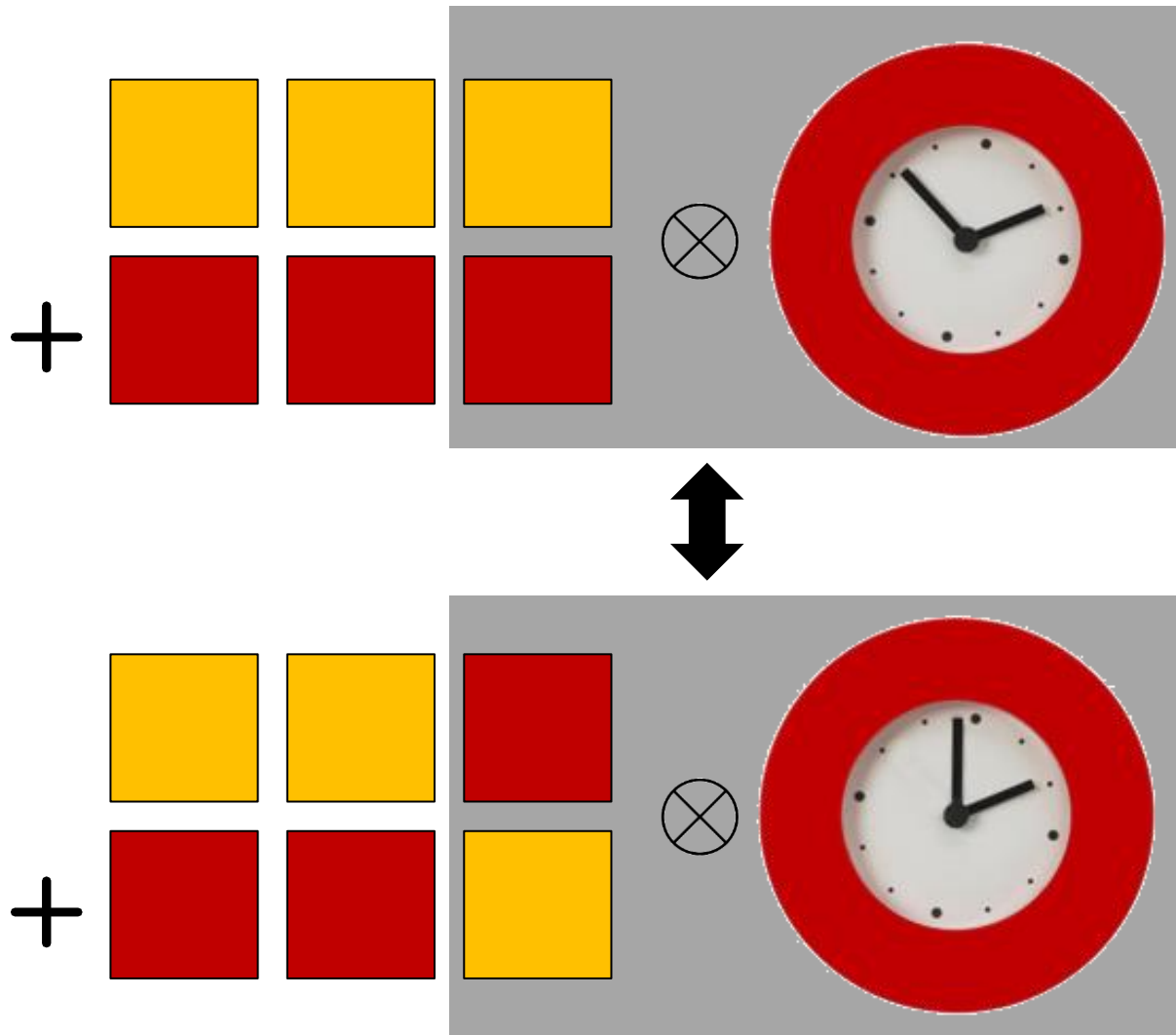
2 The data & the clock



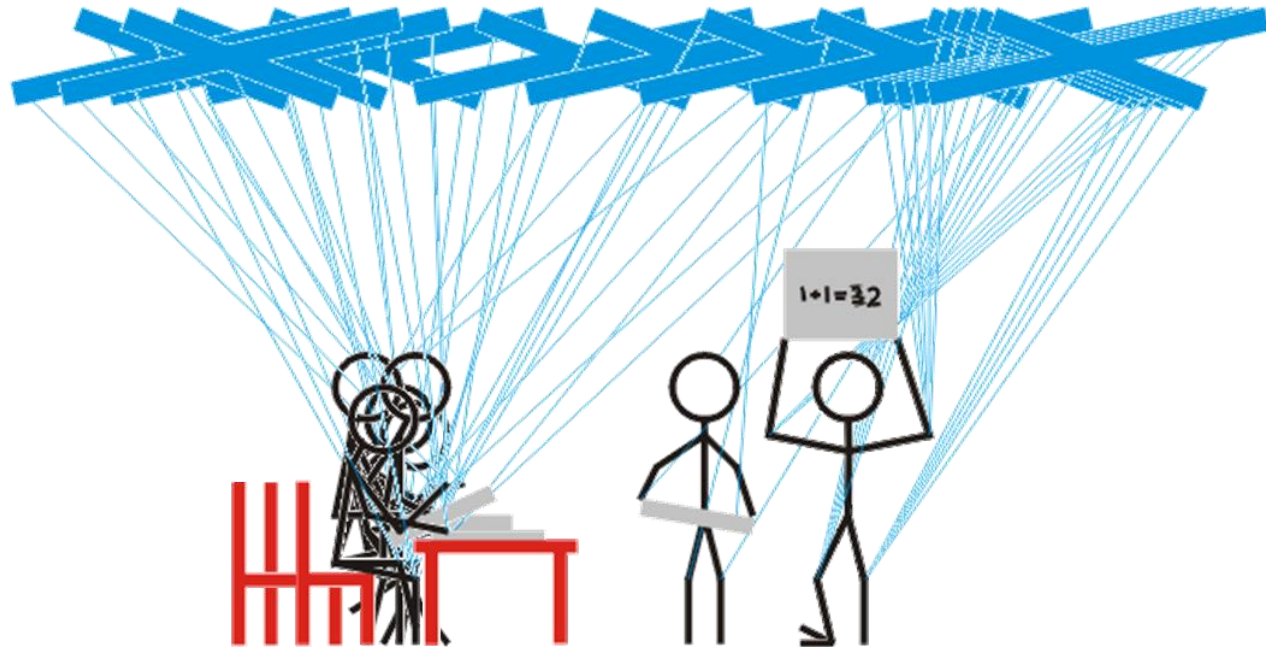
2 The data & the clock



2 The data & the clock: locally comparing related states



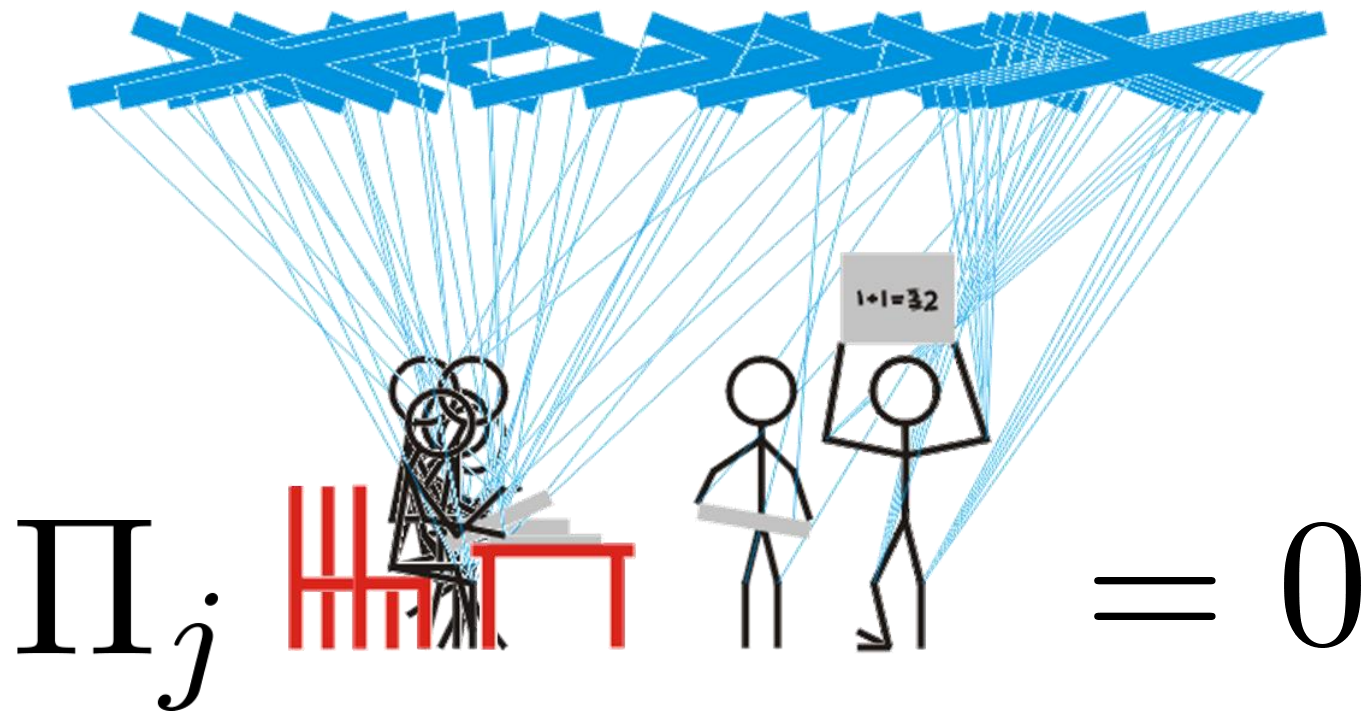
2 The history state



$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

$\underbrace{U_t \cdots U_1 |\varphi_0\rangle}_{|t\rangle}$

2 The history state: a ground state



$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T \underbrace{|\varphi_t\rangle \otimes |t\rangle}_{U_t \cdots U_1 |\varphi_0\rangle}$$

k-local *c-o-n-d-i-t-i-o-n-s*

clock encoding
state progression
initialization

$$|\dots 0\rangle \otimes |0\rangle$$

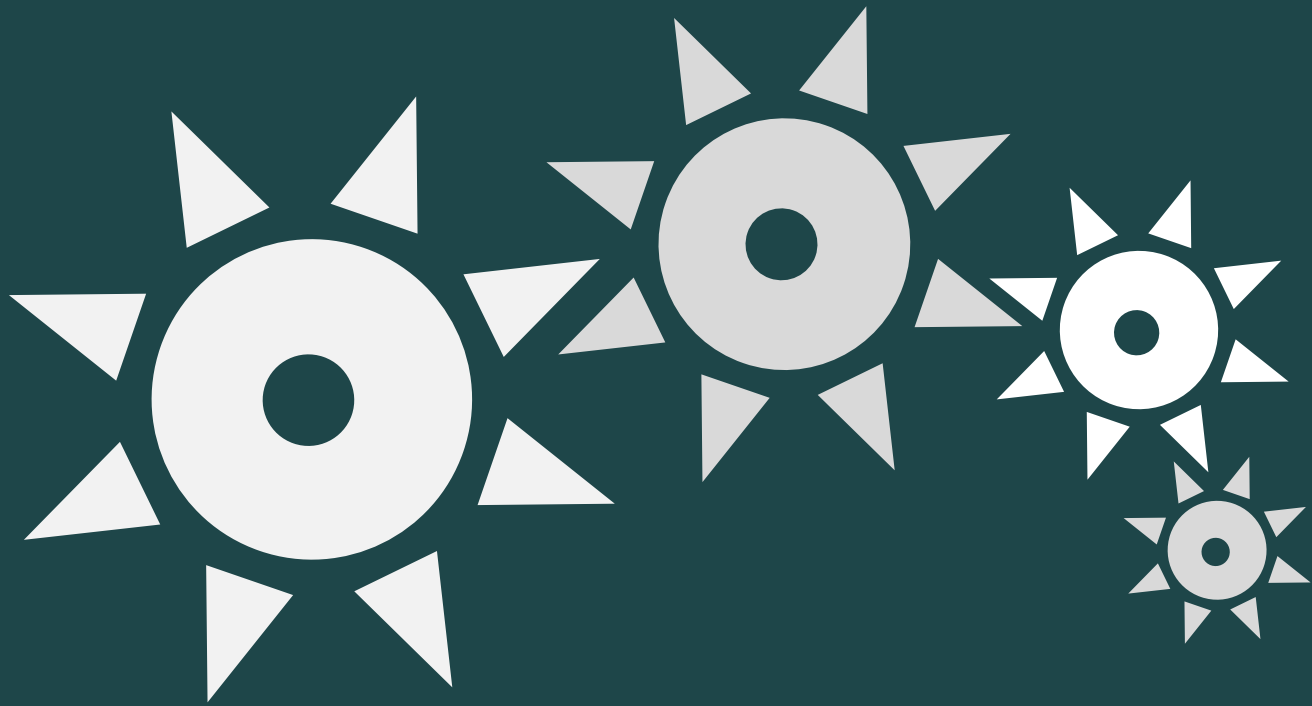
$$|\varphi_t\rangle \otimes |t\rangle$$

$$|\varphi_{t+1}\rangle \otimes |t+1\rangle$$

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

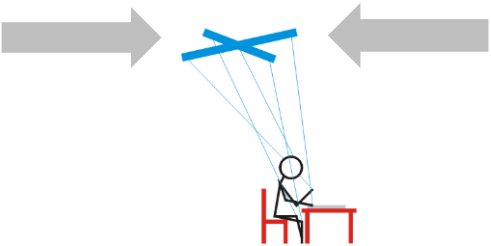

output $|\dots 1\rangle \otimes |T\rangle$

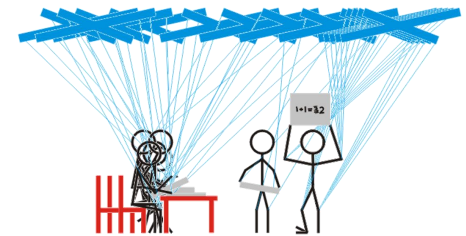
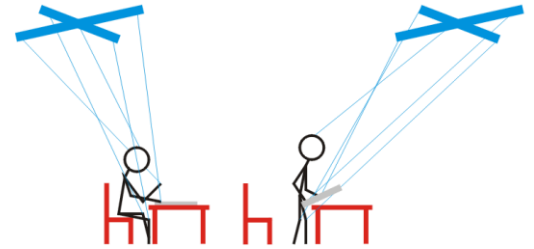




a clock workshop

3 Making a local clock

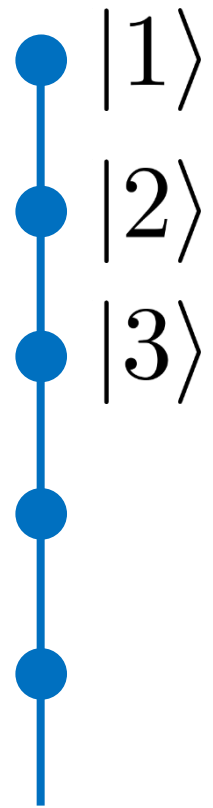
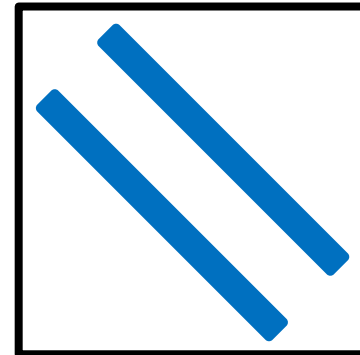
- dynamic: a system that ticks local ticks (transitions)
- static: a unique ground state the uniform tick superposition
- identifiable states 
- checking for **bad** states 



3 A quantum walk on a line is a clock


■ transitions $T_{st} = \overset{\bullet \text{---} \bullet}{|s\rangle\langle t|} + |t\rangle\langle s|$

Hamiltonian $H_w = \sum_{\langle s,t \rangle} T_{st}$




3 A line is a clock

transitions $T_{st} = |s\rangle\langle t| + |t\rangle\langle s|$

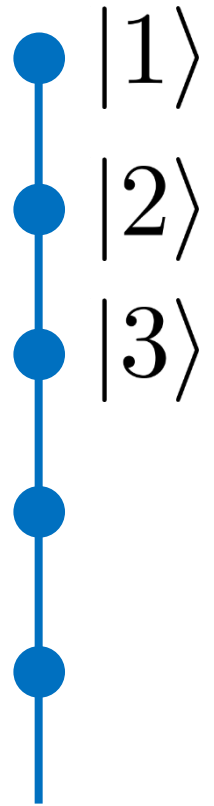
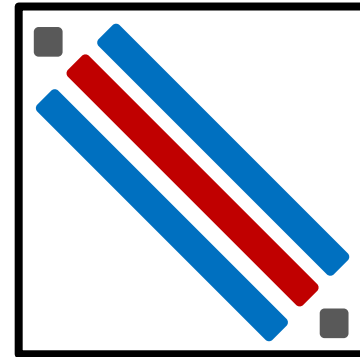


■ projections $P_{st} = \frac{1}{2} (|s\rangle - |t\rangle) (\langle s| - \langle t|)$




Hamiltonian $H = \sum_{\langle s,t \rangle} P_{st}$

the ground state $|1\rangle + |2\rangle + |3\rangle + \dots$




3 A line is a clock

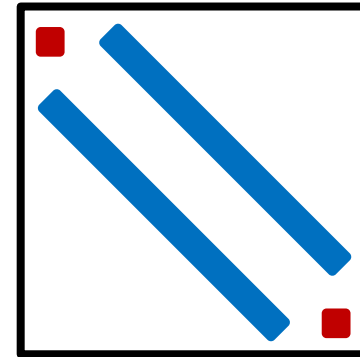
transitions $T_{st} = |s\rangle\langle t| + |t\rangle\langle s|$



■ projections $P_{st} = \frac{1}{2} (|s\rangle - |t\rangle) (\langle s| - \langle t|)$



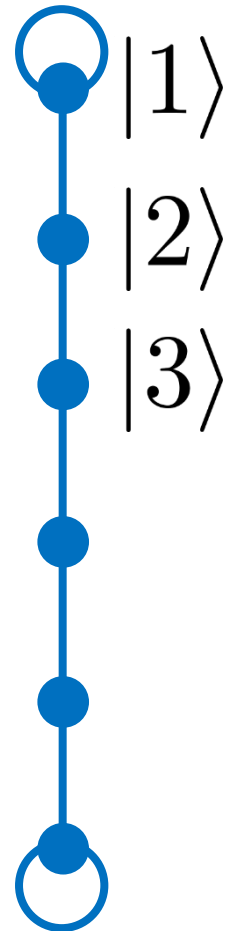
Hamiltonian $H = \sum_{\langle s,t \rangle} P_{st}$



the ground state $|1\rangle + |2\rangle + |3\rangle + \dots$

other eigenstates $|\varphi_p\rangle \propto \sum_{s=1}^N \cos(ps) |s\rangle$ $E_p = 2 \cos p$

the gap $\Delta = \Theta \left(\frac{1}{N^2} \right)$ $p = \frac{k\pi}{N}$



3 A pulse clock



- transitions

2-local

$$|10\rangle\langle 01| + |01\rangle\langle 10|$$

- identification

$$|1\rangle\langle 1|$$

- projections

get a superposition
for the ground state

$$|01 - 10\rangle\langle 01 - 10|$$

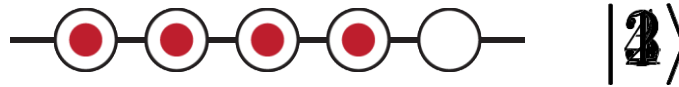
- invariant subspaces & tuning

a given number of excitations
tuning for a single excitation: prefer 1, hate 11

$$\begin{aligned} &1000 \\ + &0100 \\ + &0010 \\ + &0001 \end{aligned}$$

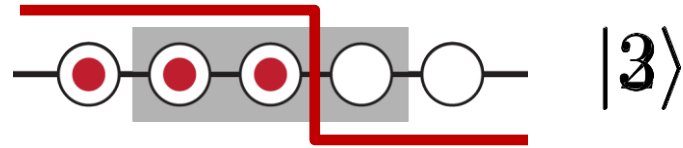


3 A domain wall (unary) clock



- clock checking $|01\rangle\langle 01|$
2-local
- identification $|10\rangle\langle 10|$

3 A domain wall (unary) clock



■ clock checking $|01\rangle\langle 01|$
2-local

■ identification $|10\rangle\langle 10|$

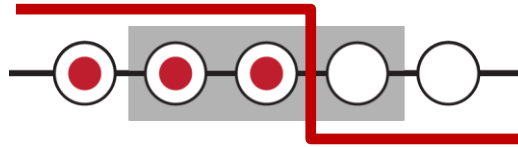
■ projections $|100 - 110\rangle\langle 100 - 110|$
3-local

■ a single domain wall: fix the ends
a unique ground state



$$10000 + 11000 + 11100 + 11110$$

3 The DW clock in Kitaev's 5-local Hamiltonian



- clock checking $|01\rangle\langle 01|$
2-local


- identification $|10\rangle\langle 10|$

- projections $|100 - 110\rangle\langle 100 - 110|$
3-local

- interacting with data $\frac{1}{2} (|t+1\rangle\langle t+1| + |t\rangle\langle t|)$
5-local
 $-\frac{1}{2} (U_{t+1} \otimes |t+1\rangle\langle t| + U_{t+1}^\dagger \otimes |t\rangle\langle t+1|)$

3 Kitaev's LH: the playground

- Invariant subspaces

with bad clock states 
labeled by the initial state


$$|\text{stuff}\rangle \otimes |\text{whatever}\rangle$$
$$\begin{aligned} & |\varphi_1\rangle \otimes |0\rangle_c \\ & U_1 |\varphi_1\rangle \otimes |1\rangle_c \\ & U_2 U_1 |\varphi_1\rangle \otimes |2\rangle_c \\ & U_3 U_2 U_1 |\varphi_1\rangle \otimes |3\rangle_c \\ & U_4 U_3 U_2 U_1 |\varphi_1\rangle \otimes |4\rangle_c \end{aligned}$$

$$\begin{aligned} & |\varphi'_1\rangle \otimes |0\rangle_c \\ & U_1 |\varphi'_1\rangle \otimes |1\rangle_c \\ & U_2 U_1 |\varphi'_1\rangle \otimes |2\rangle_c \\ & U_3 U_2 U_1 |\varphi'_1\rangle \otimes |3\rangle_c \\ & U_4 U_3 U_2 U_1 |\varphi'_1\rangle \otimes |4\rangle_c \end{aligned}$$

3 The history state: a line of states

- a projector Hamiltonian
kernel: the uniform superposition

$$\begin{aligned} &|t + 1\rangle\langle t + 1| + |t\rangle\langle t| \\ &- U_{t+1} \otimes |t + 1\rangle\langle t| \\ &- U_{t+1}^\dagger \otimes |t\rangle\langle t + 1| \end{aligned}$$

$$|\varphi_t\rangle \otimes |t\rangle$$

$$|\varphi_{t+1}\rangle \otimes |t + 1\rangle$$

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$



- endpoints: ancilla initialization/final acceptance



● YES

ground state

- NO

ground state

● NO



lower bound on the
ground state energy

good
clock
states

not
clock
states

history states

non-uniform
superpositions

history states



a polynomially small gap

$$\Delta = O(L^{-2})$$



history states

well

badly

initialized history states

well

initialized histories

accepted
states

well

initialized histories

accepted
states

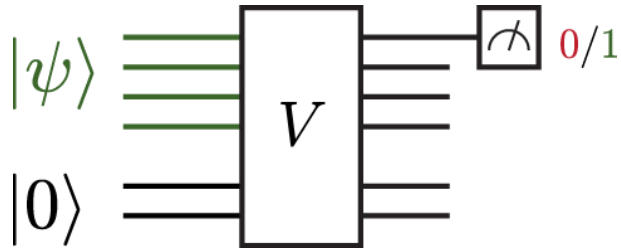
$$H_A + H_B$$

$$\lambda_0 \geq \sin^2 \frac{\vartheta}{2} \times \min(\Delta_A, \Delta_B)$$

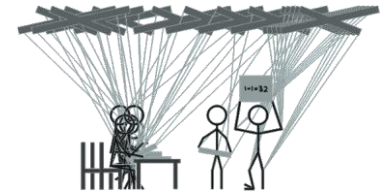
\uparrow L^{-2} \uparrow L^{-1}

3-LH and QMA verification

[N., Mozes 07]



$$H_{clock} + H_{init} + H_{prop} + H_{out}$$



NO V is unlikely to accept anything (ϵ)

lowest eigenvalue

$$\geq \frac{c(1 - \sqrt{\epsilon})}{L^2}$$



promise gap L^{-2}

(needs $\epsilon=L^{-1}$)


YES some proof is likely $(1-\epsilon)$ accepted

energy of the history

$$\leq \frac{\epsilon}{L + 1}$$

projections & gadgets

3 Lower locality for the price of bad transitions

■ the domain wall 

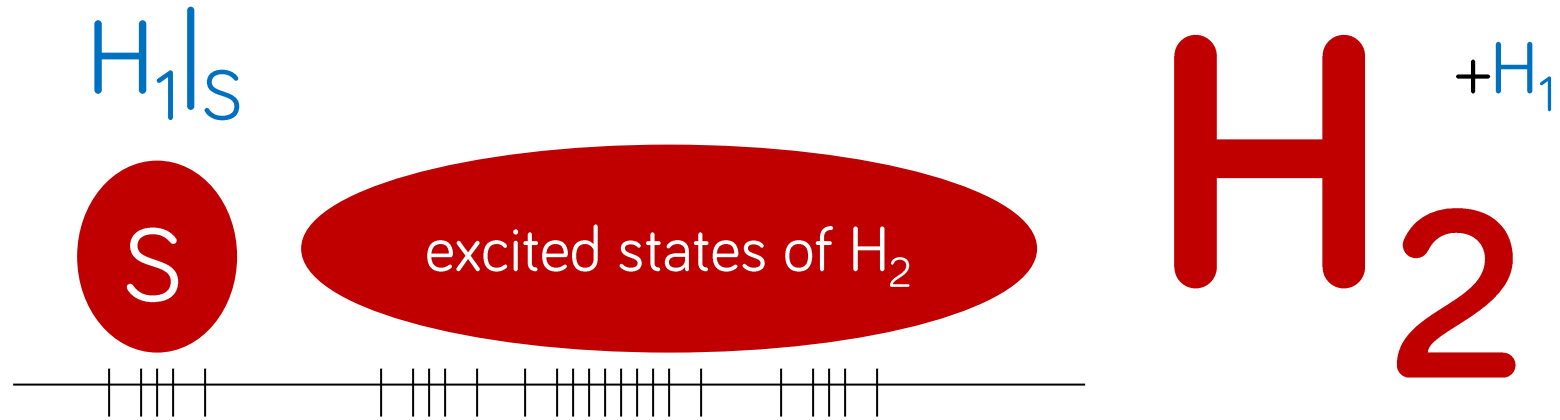
■ projections
now just 1-local $|1 - 0\rangle\langle 1 - 0|$

■ clock checks
2-local, STRONG $|01\rangle\langle 01|$

■ the ground state is close to

$$10000 + 11000 + 11100 + 11110$$

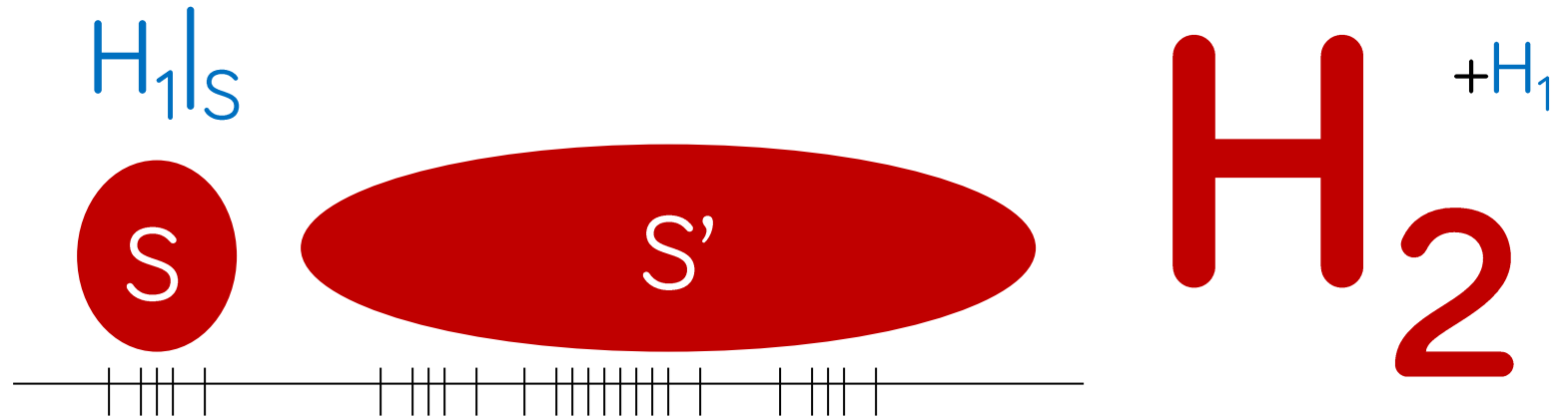
3 The projection lemma to estimate eigenvalues



Lemma 1 Let $H = H_1 + H_2$ be the sum of two Hamiltonians operating on some Hilbert space $\mathcal{H} = \mathcal{S} + \mathcal{S}^\perp$. The Hamiltonian H_2 is such that \mathcal{S} is a zero eigenspace and the eigenvectors in \mathcal{S}^\perp have eigenvalue at least $J > 2\|H_1\|$. Then,

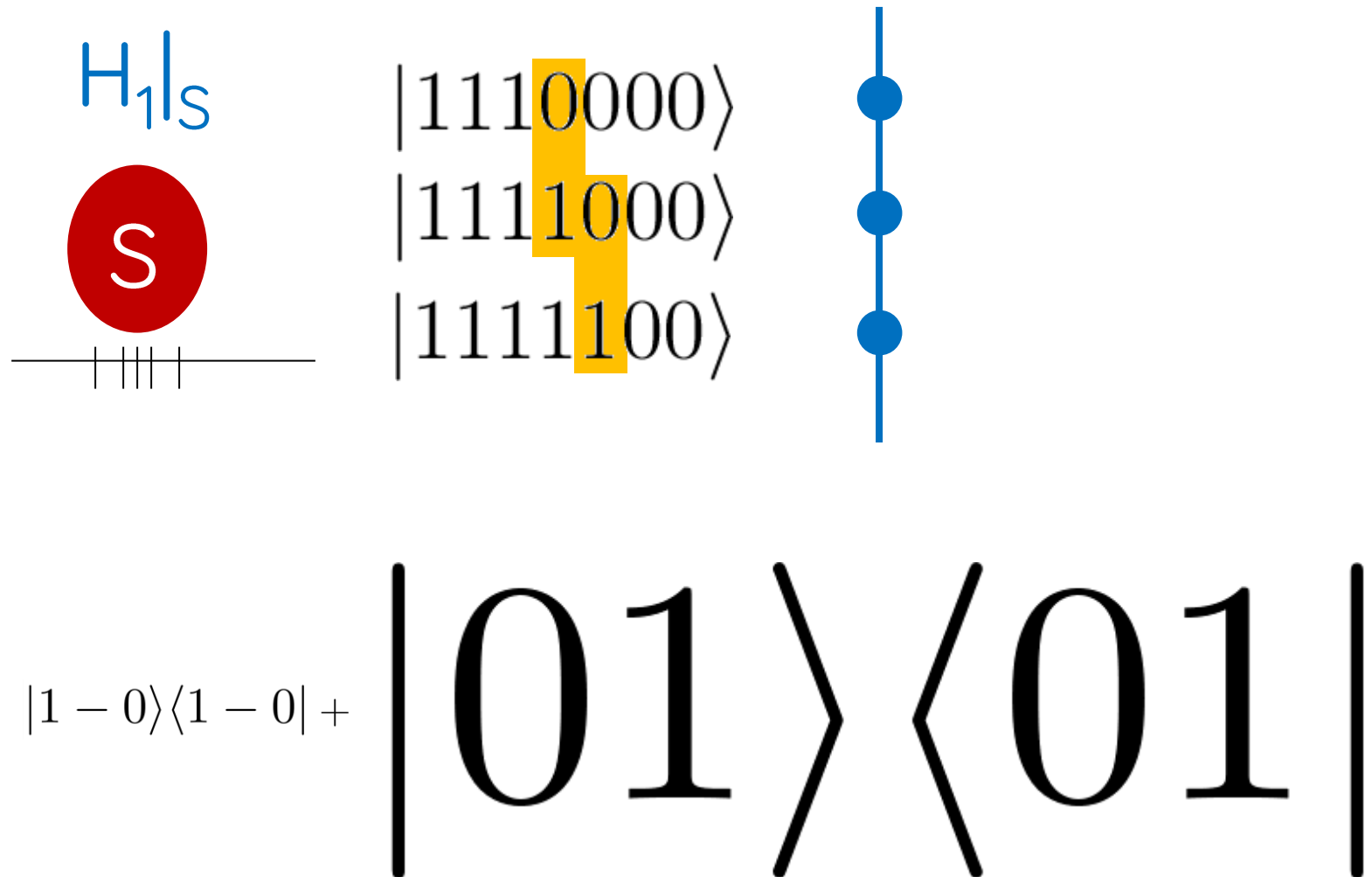
$$\lambda(H_1|_{\mathcal{S}}) - \frac{\|H_1\|^2}{J - 2\|H_1\|} \leq \lambda(H) \leq \lambda(H_1|_{\mathcal{S}})$$

3 The projection lemma to estimate eigenvalues



$$|1-0\rangle\langle 1-0| + |01\rangle\langle 01|$$

3 The projection lemma to estimate eigenvalues



3 Lower locality (3-LH) for the price of bad transitions

■ the domain wall 

■ non-projector
2-local terms $|10\rangle\langle 10|_{1,2} + |10\rangle\langle 10|_{2,3} - X_2$

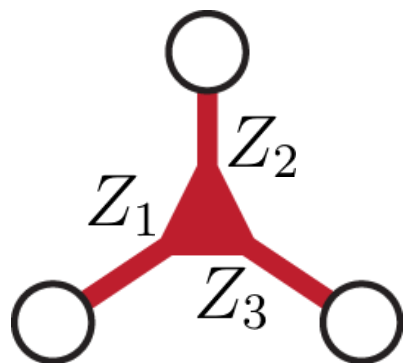
■ clock checks
2-local, BIG $|01\rangle\langle 01|$

■ from 5- to 3-local Hamiltonian [Kempe, Regev]

$$|10\rangle\langle 10|_{1,2} + |10\rangle\langle 10|_{2,3} - |1\rangle\langle 0|_2 \otimes U - |0\rangle\langle 1|_2 \otimes U^\dagger$$

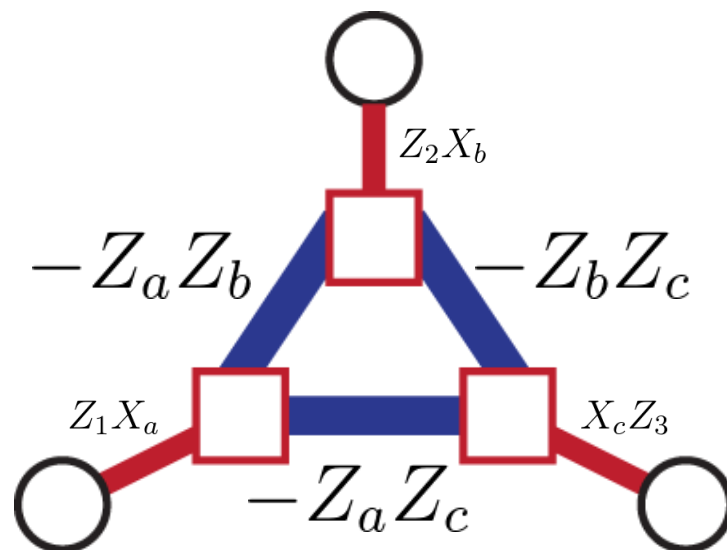
restricted to good clocks: qw on a line 

3 Further decreasing locality: a “3 from 2” gadget



- strongly coupled ancillas (a new energy scale)
- perturbation theory

$$G'(z) = (z\mathbb{I} - H')^{-1}$$

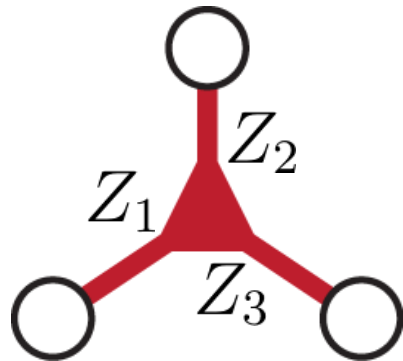


$$H' = H + V$$

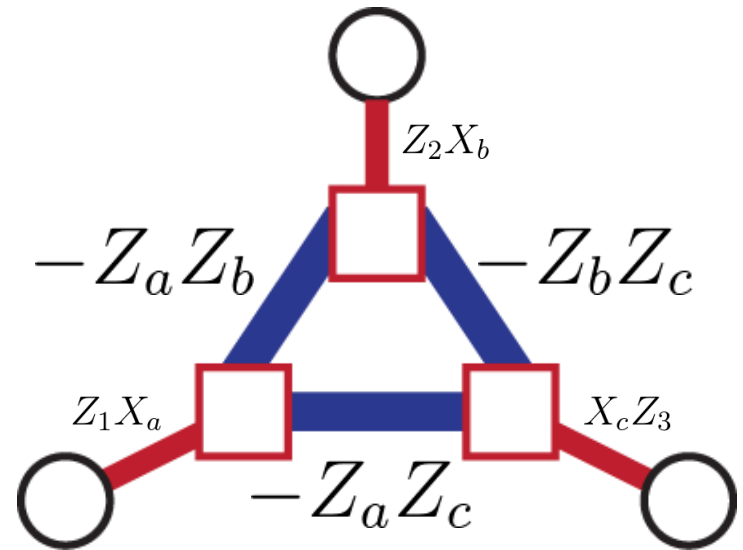
$\|H\| \gg \|V\|$

$$S = \text{span} \{ |000\rangle, |111\rangle \}$$

3 Further decreasing locality: a “3 from 2” gadget



- strongly coupled ancillas (a new energy scale)
- perturbation theory gives us an effective Hamiltonian



$$H' = H + V$$

$\|H\| \gg \|V\|$

$$S = \text{span} \{ |000\rangle, |111\rangle \}$$

$$V|_S$$

projection
lemma

$$V^2|_S$$

unwanted
(subtract)

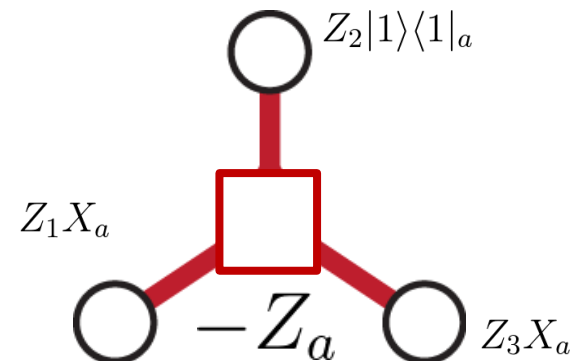
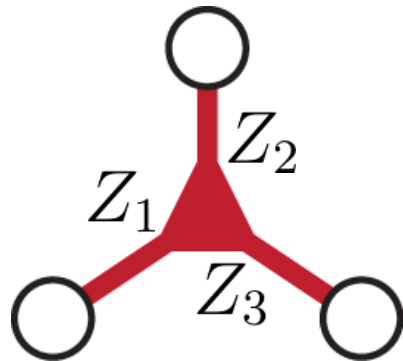
$$V^3|_S$$

the effective
3-local term

[Kempe, Kitaev, Regev '03]

3 STRONG local fields, OK interactions

[Cao et al., 1311.2555]



- strongly bound a single ancilla
no superstrong interactions
- perturbation theory gives us
an effective Hamiltonian

$$S = \{|0\rangle\}$$

$$H' = H + V$$

$$\|H\| \gg \|V\|$$

$$V|_S$$

projection
lemma

$$V^2|_S$$

unwanted
(subtract)

$$V^3|_S$$

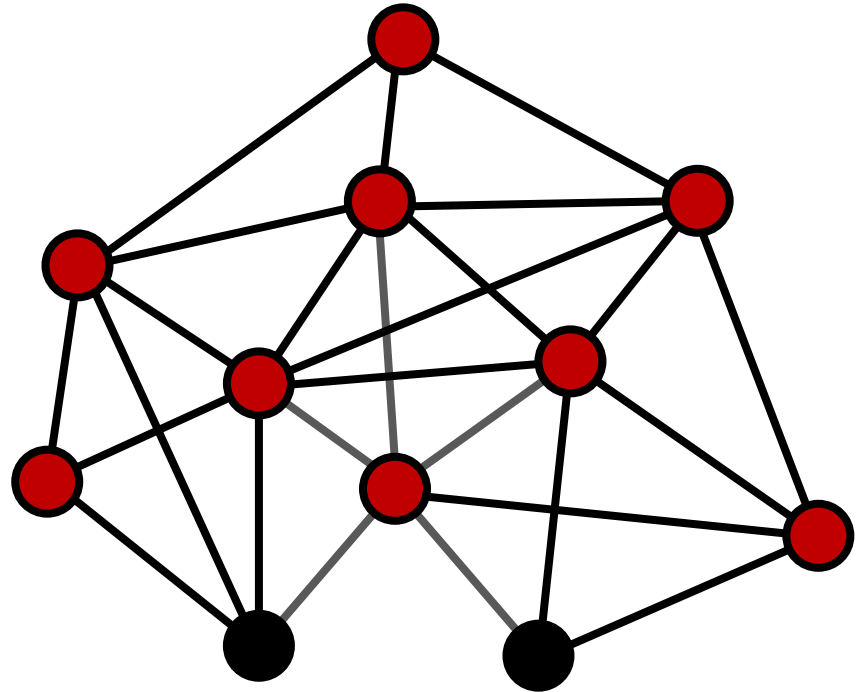
the effective
3-local term

special cases (Z-basis)
exact gadgets!

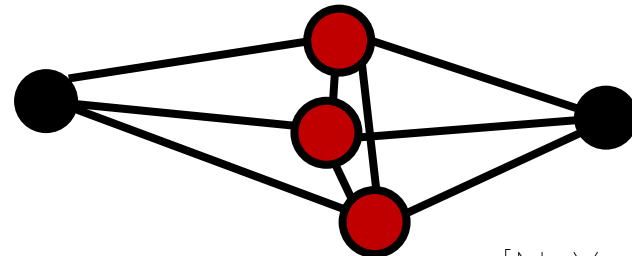
[Jacob Biamonte 0801.3800]

3 “Strengthening”, intermediary gadgets?

- classically easy: copy

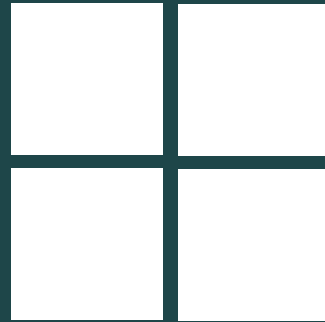
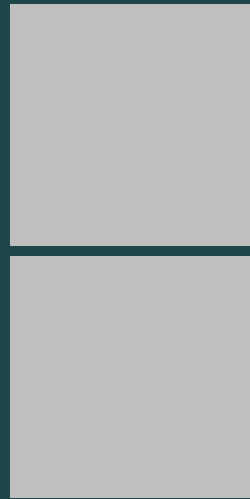


- quantumly?

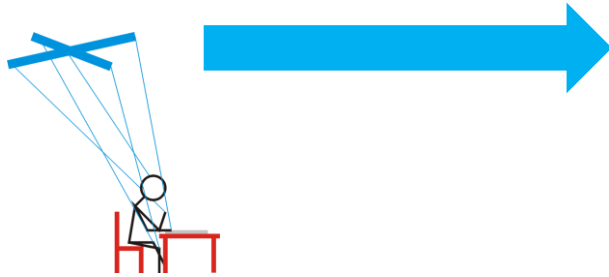


[N., Yudong Cao]

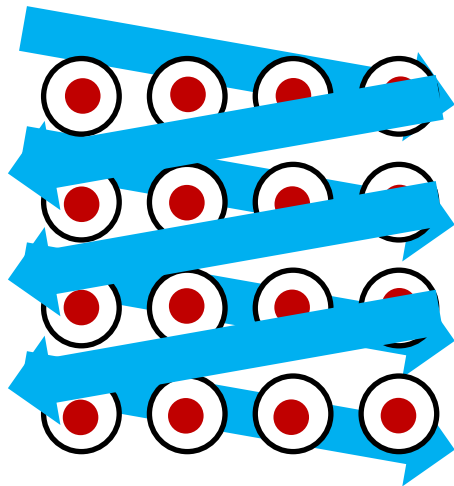
locality & dimensionality



clock/work registers

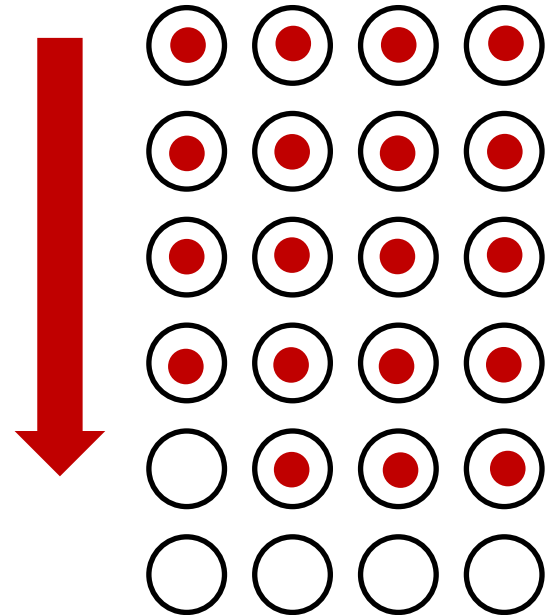


[KKR03, OT05]



constant degree
geometric locality

a geometric clock

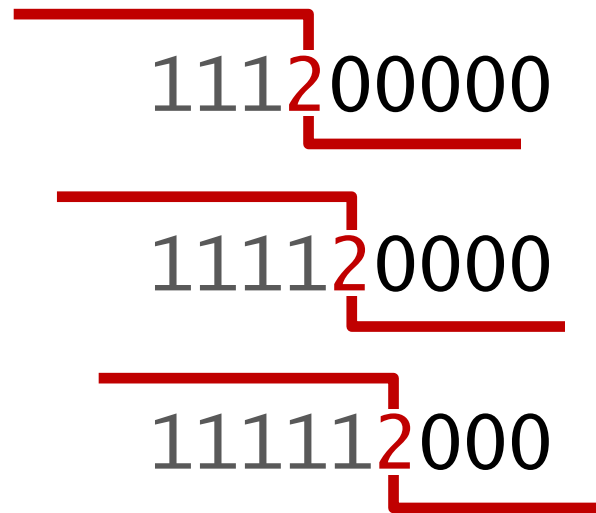


[Mize] [Janzing] [AvDKLLR] [BT13]



[AGIK07] moving data in 1D

3 Moving a special site: the qutrit surfer



- clock checking

$ 10\rangle\langle 10 $	$ 02\rangle\langle 02 $	$ 22\rangle\langle 22 $
$ 01\rangle\langle 01 $	$ 21\rangle\langle 21 $	

 2-local + ends

- identification

$ 2\rangle\langle 2 $

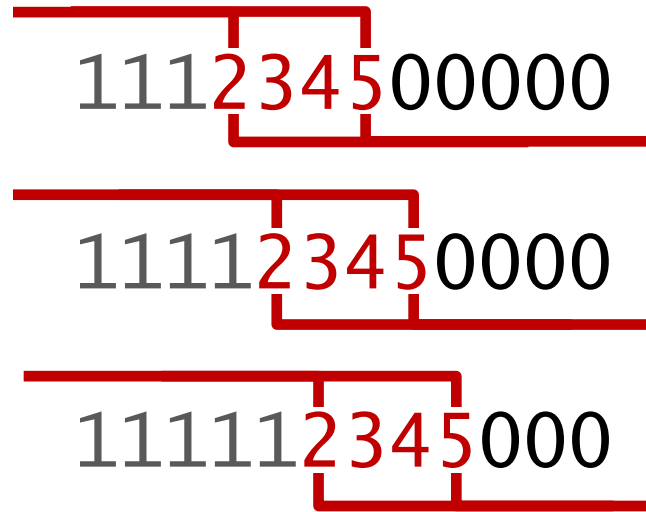
- projections

$ 20 - 12\rangle\langle 20 - 12 $

 2-local

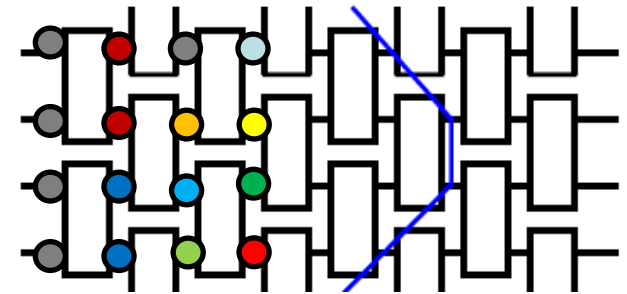
3 Constructing local, geometric clocks: moving the data

- telling “time”
by where
the data is



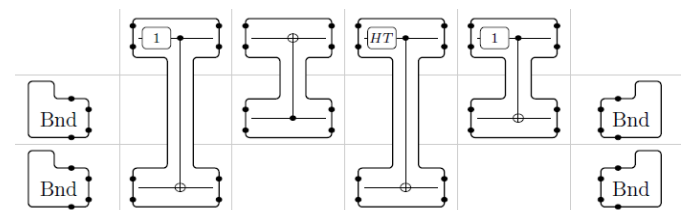
mutually
orthogonal
states

- carrying/moving data?
larger qudits (local dim.)
larger locality



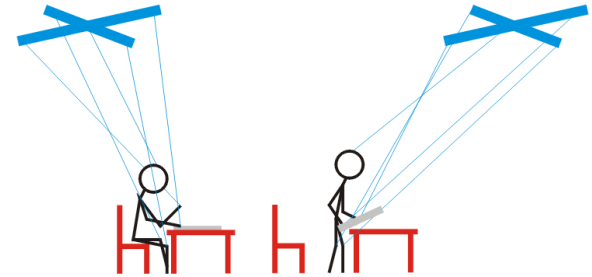
internal states ... dual-rail

[Childs Gosset Webb 13]

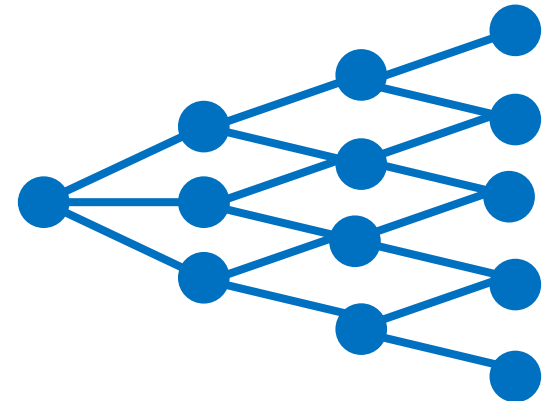


2 Making a good local clock

- identifiable states
 - domain-wall structure
 - local transitions
 - easily checkable **bad** states



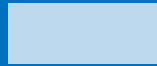
- different geometry?
 - simpler terms?
 - locality/qudits?
 - beyond linear?
 - larger (promise) gaps?



2 Hamiltonian Quantum Cellular Automata in 1D

- moving the program instead of the data

[N., Wocjan 07]



data

program particles diffuse above data
special states stand in their way

2 Hamiltonian Quantum Cellular Automata in 1D

- moving the program instead of the data

[N., Wocjan 07]



data

program particles diffuse above data
special states stand in their way

2 Hamiltonian Quantum Cellular Automata in 1D

- moving the program instead of the data

[N., Wocjan 07]



data

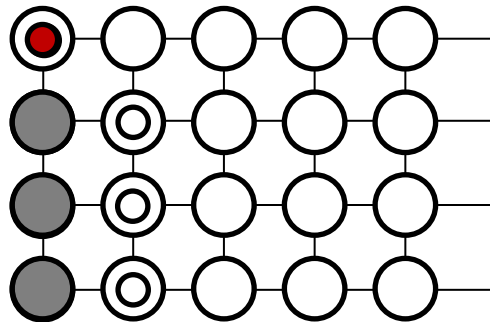
program particles diffuse above data
special states stand in their way

- BQP in 1D with a trans. invariant, time independent LH computational basis programmable
- a nonlinear clock, polynomial expected runtime

3 A local, sequential geometric clock in 2D

- 2D “sequential” evaluation

[Aharonov van Dam Kempe Landau Lloyd Regev 04]



Forbidden	Guarantees that
$\bigcirc \oplus, \bigcirc \otimes, \bigcirc \otimes$	\bigcirc is to the right of all other qubits
$\bigcirc \otimes, \oplus \otimes, \otimes \otimes$	\otimes is to the left of all other qubits
$\bigcirc \otimes, \otimes \bigcirc$	\bigcirc and \otimes are not horizontally adjacent
$\oplus \oplus, \oplus \otimes, \otimes \oplus, \otimes \otimes$	only one of \oplus, \otimes per row
$\bigcirc, \oplus, \otimes$ \otimes, \oplus, \oplus	only \otimes above \oplus
\oplus, \oplus, \oplus $\bigcirc, \otimes, \otimes$	only \oplus below \oplus
\bigcirc, \otimes \otimes, \bigcirc	\bigcirc and \otimes are not vertically adjacent
\otimes, \otimes \bigcirc, \oplus	no \bigcirc below \otimes and no \oplus below \otimes

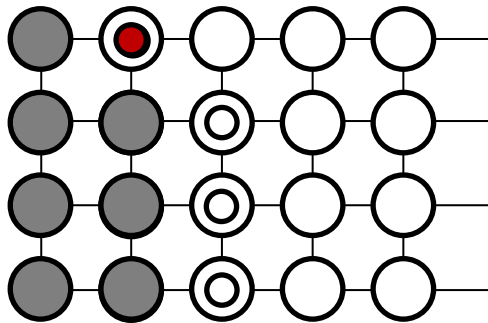
universality of adiabatic QC

2-local interactions, $d=6$ qudits

3 A local, sequential geometric clock in 2D

- 2D “sequential” evaluation

[Aharonov van Dam Kempe Landau Lloyd Regev 04]



Forbidden	Guarantees that
$\circ \oplus, \circ \otimes, \circ \otimes$	\circ is to the right of all other qubits
$\circ \otimes, \oplus \otimes, \otimes \otimes$	\otimes is to the left of all other qubits
$\circ \otimes, \otimes \circ$	\circ and \otimes are not horizontally adjacent
$\oplus \oplus, \oplus \otimes, \otimes \oplus, \otimes \otimes$	only one of \oplus, \otimes per row
\circ, \oplus, \otimes \otimes, \oplus, \oplus	only \otimes above \oplus
\oplus, \oplus, \oplus \circ, \otimes, \otimes	only \oplus below \oplus
\circ, \otimes \otimes, \circ	\circ and \otimes are not vertically adjacent
\otimes, \otimes \circ, \oplus	no \circ below \otimes and no \oplus below \otimes

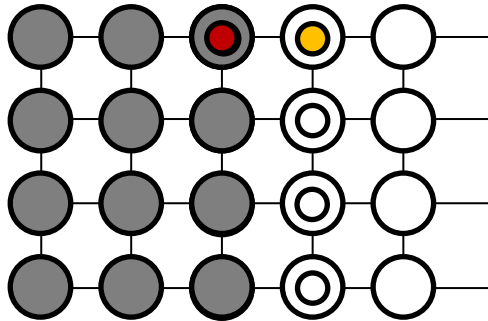
universality of adiabatic QC

2-local interactions, $d=6$ qudits

3 A local, sequential geometric clock in 2D

- 2D “sequential” evaluation

[Aharonov van Dam Kempe Landau Lloyd Regev 04]



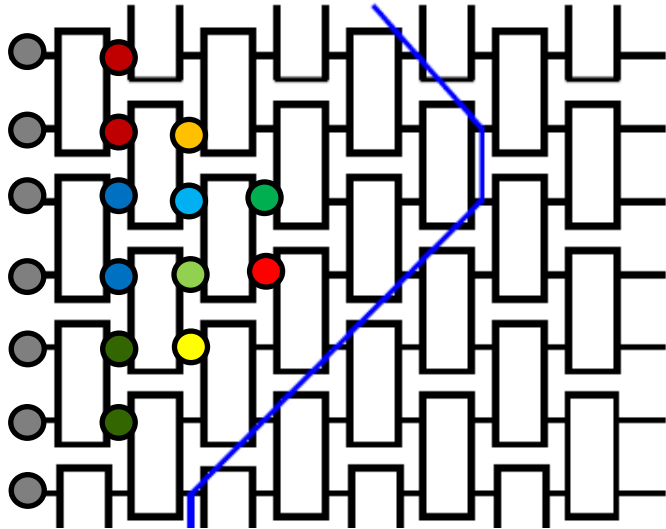
Forbidden	Guarantees that
$\circ\oplus, \circ\otimes, \circ\otimes$	\circ is to the right of all other qubits
$\circ\otimes, \oplus\otimes, \otimes\otimes$	\otimes is to the left of all other qubits
$\circ\otimes, \otimes\circ$	\circ and \otimes are not horizontally adjacent
$\oplus\oplus, \oplus\otimes, \otimes\oplus, \otimes\otimes$	only one of \oplus, \otimes per row
\circ, \oplus, \otimes \oplus, \oplus, \oplus	only \oplus above \oplus
\oplus, \oplus, \oplus \circ, \oplus, \otimes	only \oplus below \oplus
\circ, \otimes \otimes, \circ	\circ and \otimes are not vertically adjacent
\oplus, \otimes \circ, \oplus	no \circ below \oplus and no \oplus below \otimes

universality of adiabatic QC

2-local interactions, $d=6$ qudits

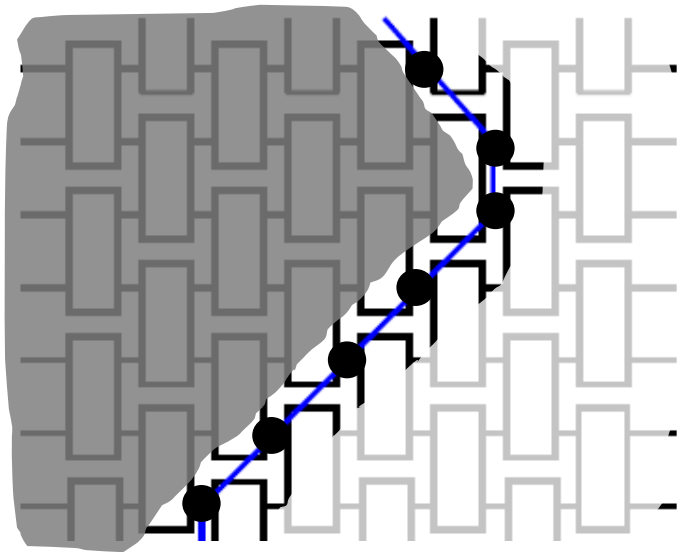
3 Another geometric clock in 2D: a string on a torus

- 2D “parallelized” evaluation [Mizel Lidar 06] [Janzing 07]
[Breuckmann Terhal 13]



3 Another geometric clock in 2D: a string on a torus

- 2D “parallelized” evaluation [Mizel Lidar 06] [Janzing 07]
[Breuckmann Terhal 13]




QMA-complete
4-local operations
 b, b^+ fermions, spin
or $d=4$ (spin $3/2$)
particle # tuning
(motivation: AQC)

- promise gap: proven $N^{-3}D^{-3}$, conjectured $N^{-2}D^{-2}=L^{-2}$

3 Constructing local, geometric clocks in 1D

- moving the data with 2-local interactions


$$(|s\rangle - |t\rangle)(\langle s| - \langle t|)$$

$$\begin{aligned} & |XY\rangle \langle XY|_{j,j+1} + |ZW\rangle \langle ZW|_{k,k+1} \\ & - |PQ\rangle \langle NO|_{i,i+1} - |NO\rangle \langle PQ|_{i,i+1} \end{aligned}$$

- higher local dimension: qudits

carry the data

mark transitions

detect bad states





we've been here

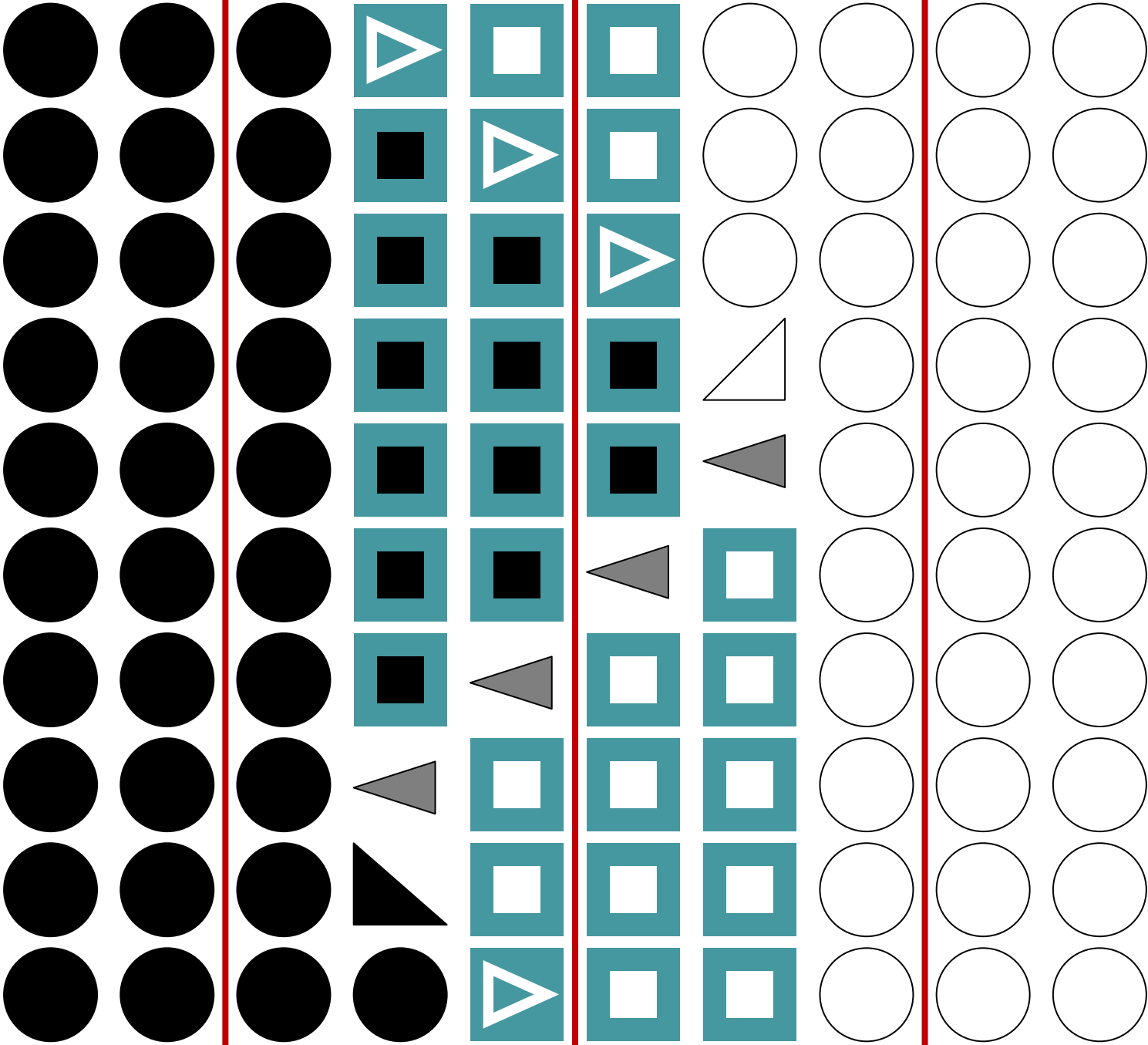
the data

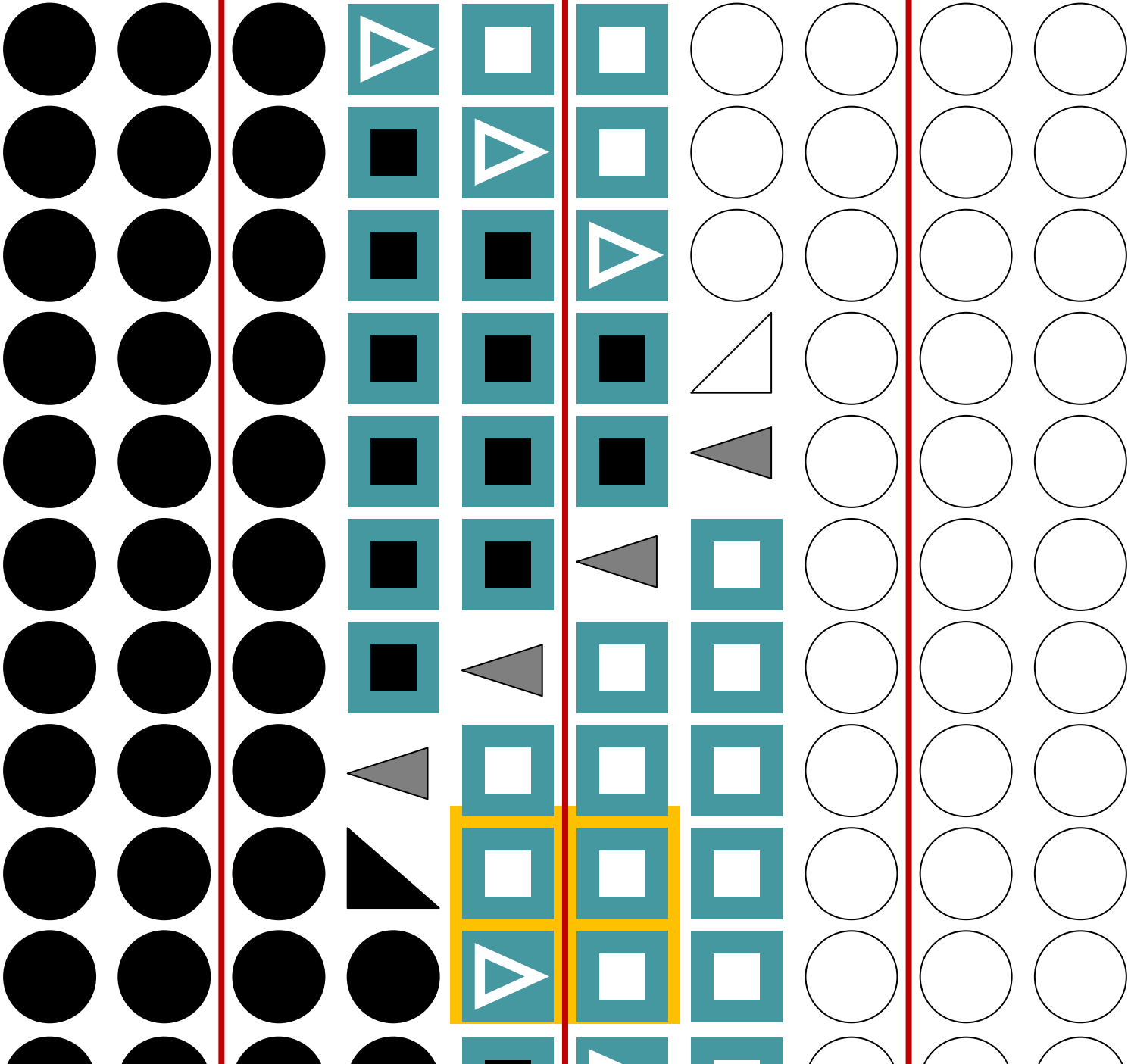
undiscovered territory

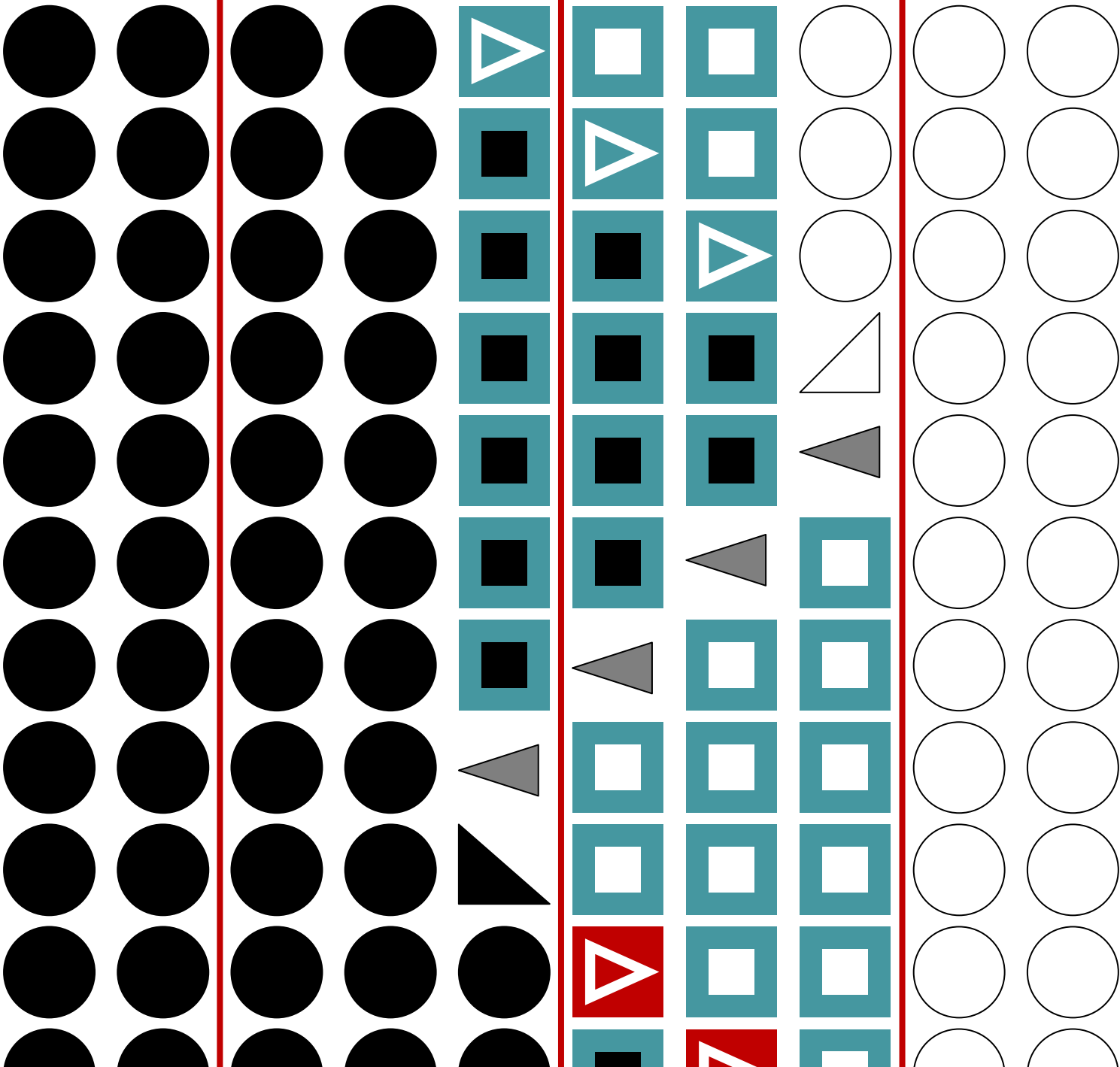
moves

the **power** of **quantum**
systems on a line

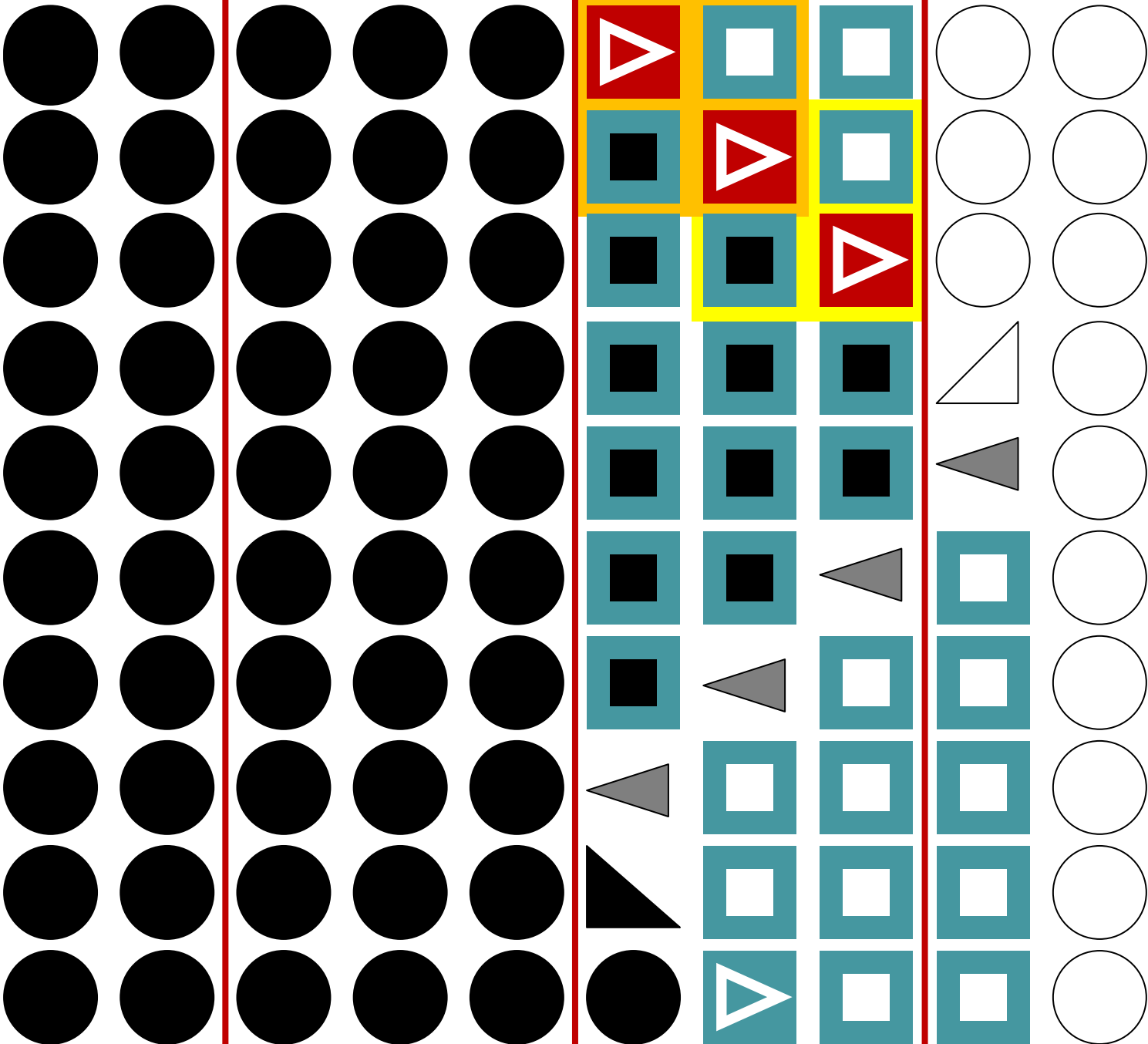
[Aharonov, Gottesman, Irani, Kempe]







U_{ab}
 U_{bc}



1 LH in 1D (2-local) with qudits

- unique state progression

every legal state goes
to exactly 2 states

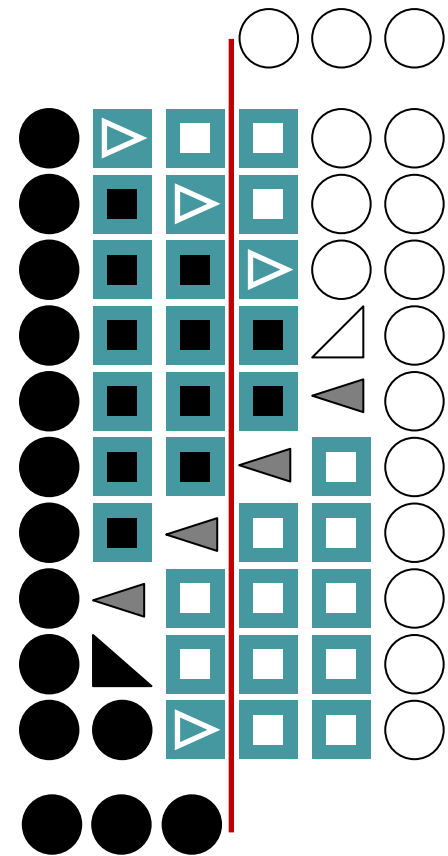
- clairvoyance

allowed but illegal states
evolve to forbidden ones

- the promise gap: L^{-3}

- an entangled ground state

special case: NP-hard [Schuch]



QMA-complete

d=13

[AGIK '06]

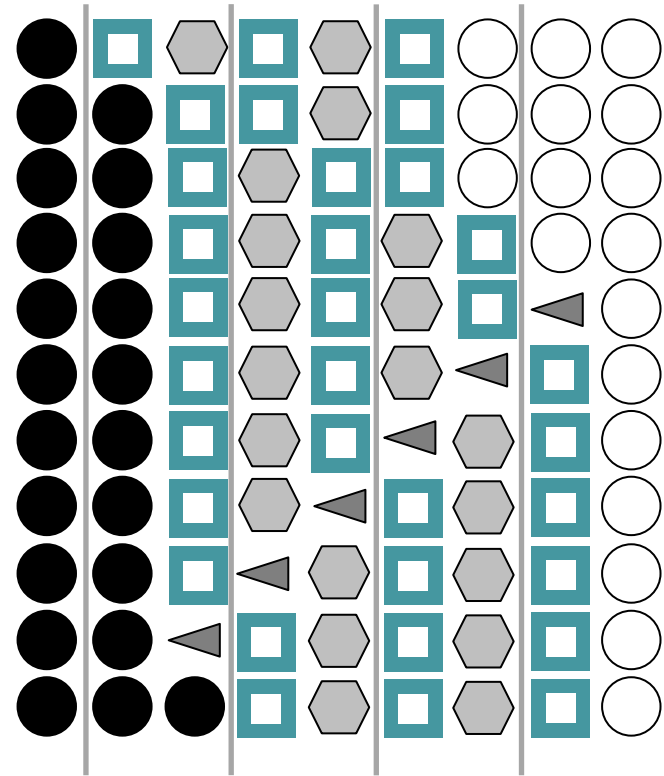
1 LH in 1D: more space = smaller qudits

- unique state progression

every legal state goes
to exactly 2 states

d=11

[N. 08]

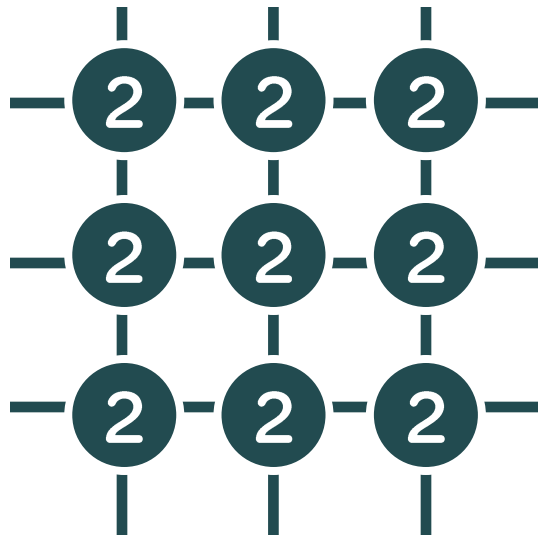


- bad but detectable transitions

d=8

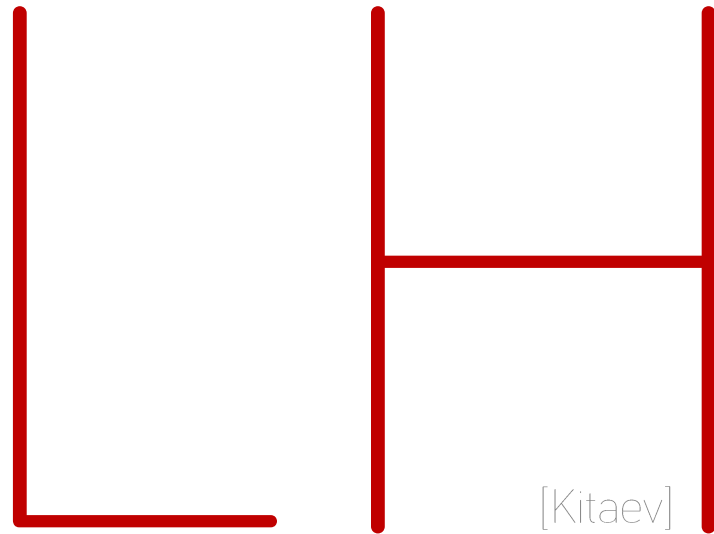
[Hallgren, N., Narayanaswami '13]

1 2-local Hamiltonian is QMA-complete

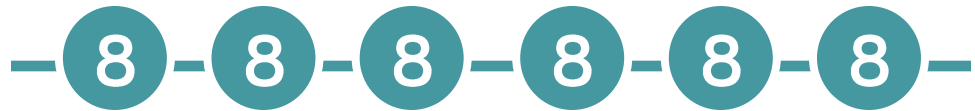


[Oliveira, Terhal '05]

a global minimum



[Kitaev]

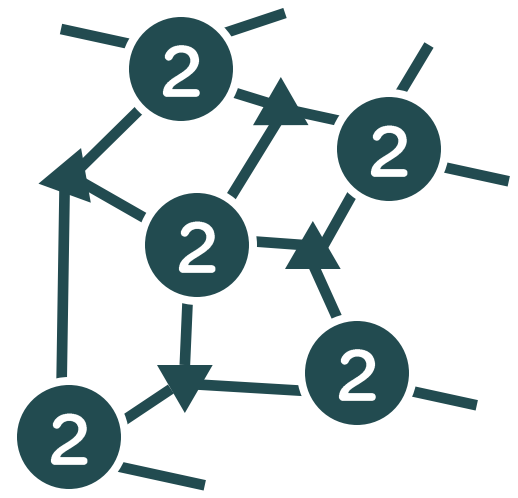


[Hallgren, N, Narayanaswami '13]

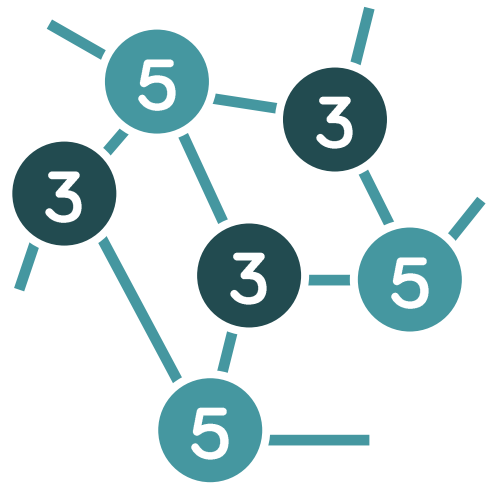
1 QMA₁-complete problems



[N. '08]



[Gosset, N. '13]



[Eldar, Regev '08]

unfrustrated

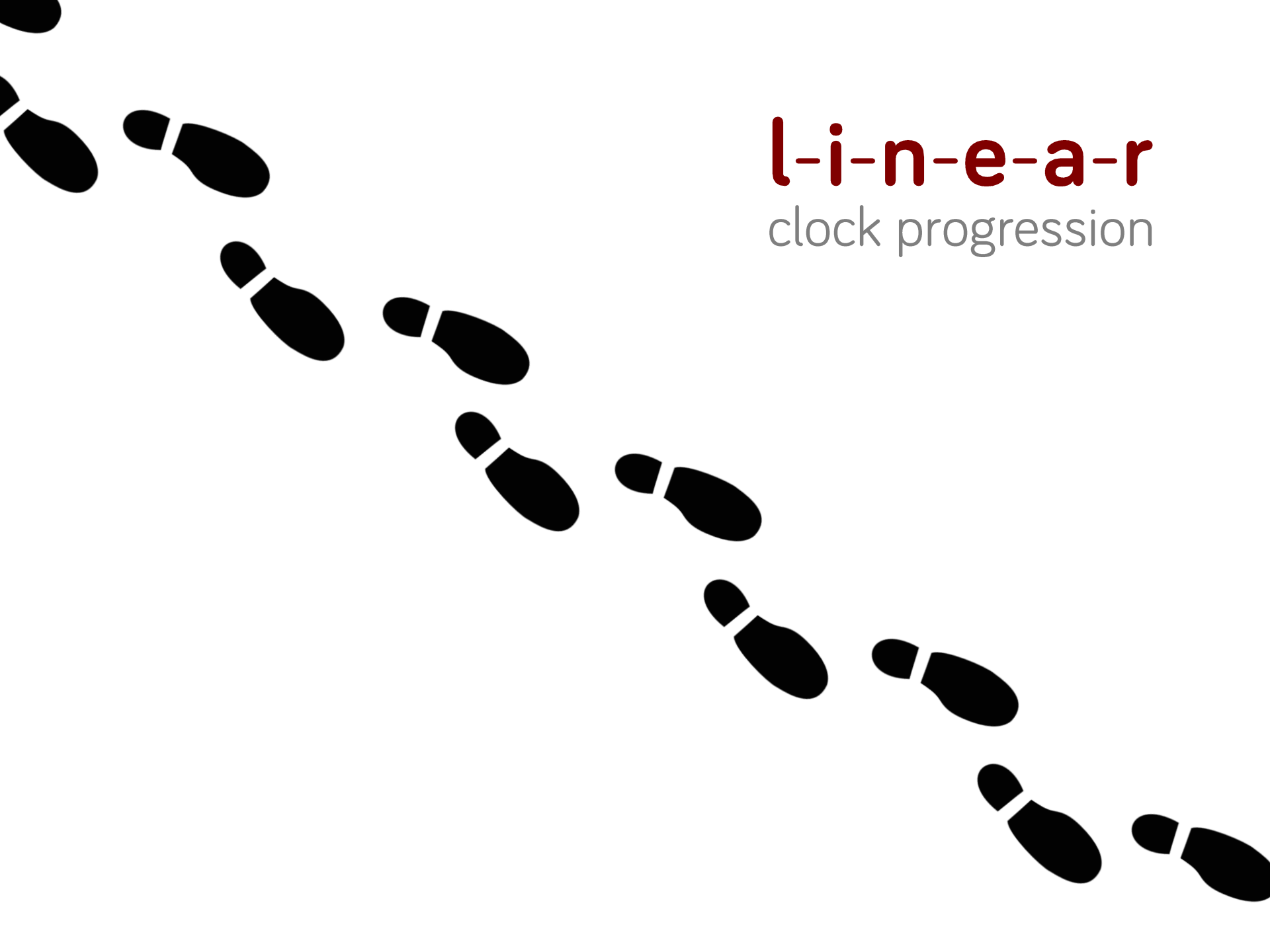
qSAT

[Bravyi]



clock

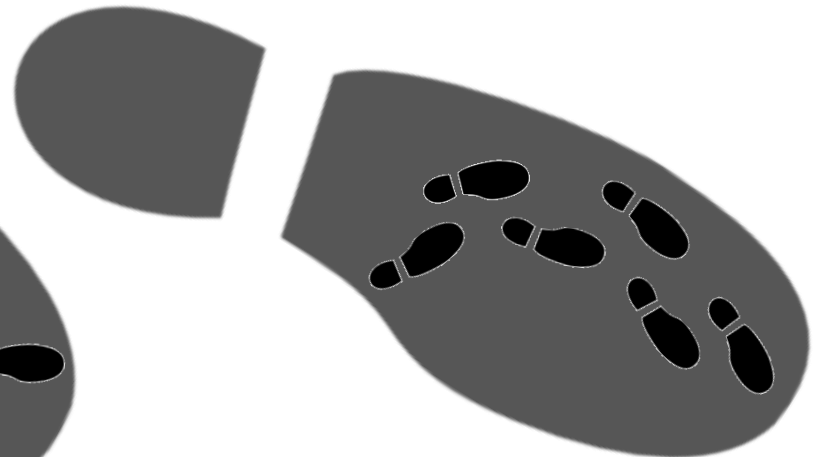
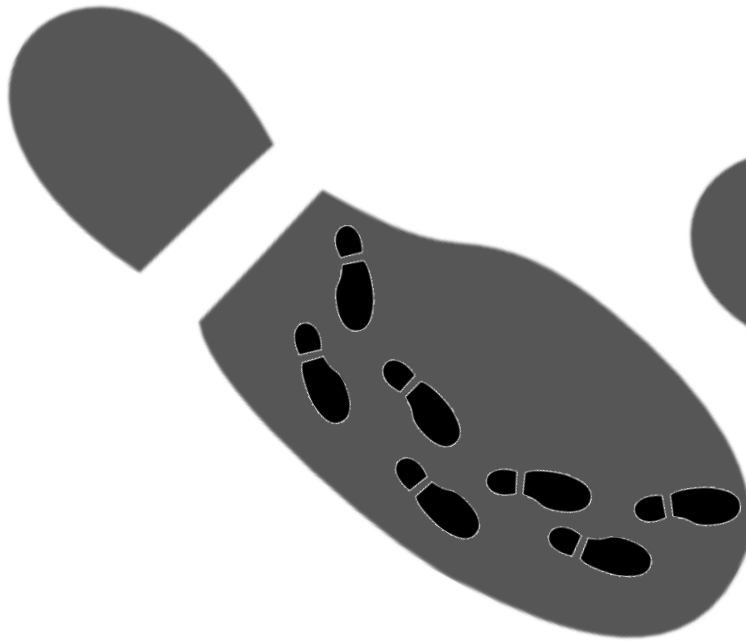
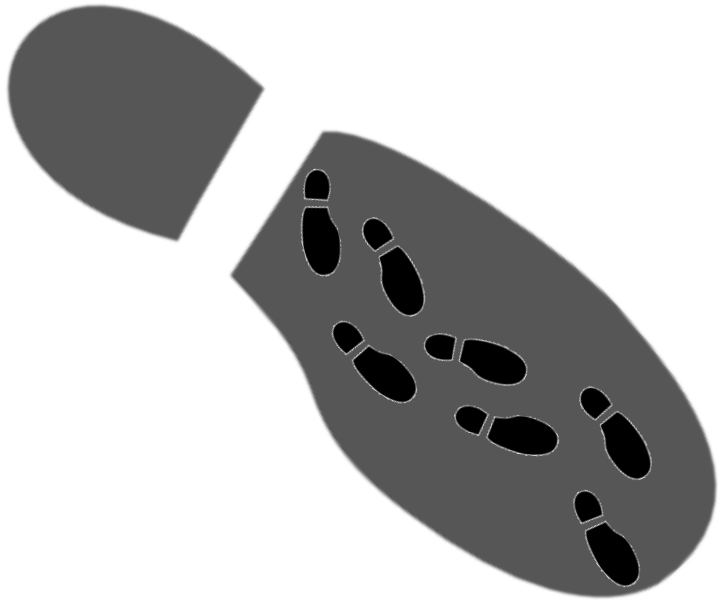
constructing 2



l-i-n-e-a-r
clock progression

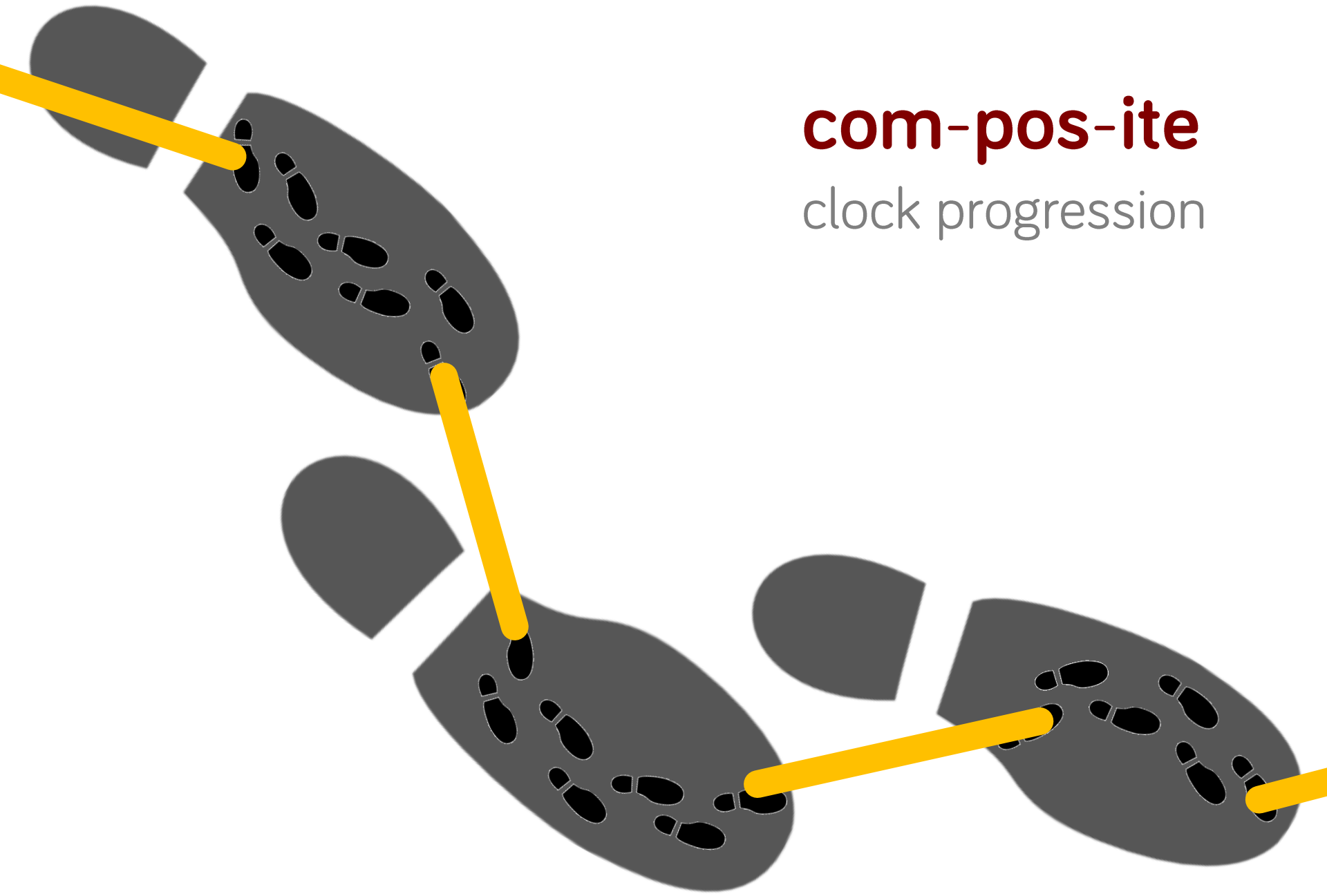
com-pos-ite

clock progression



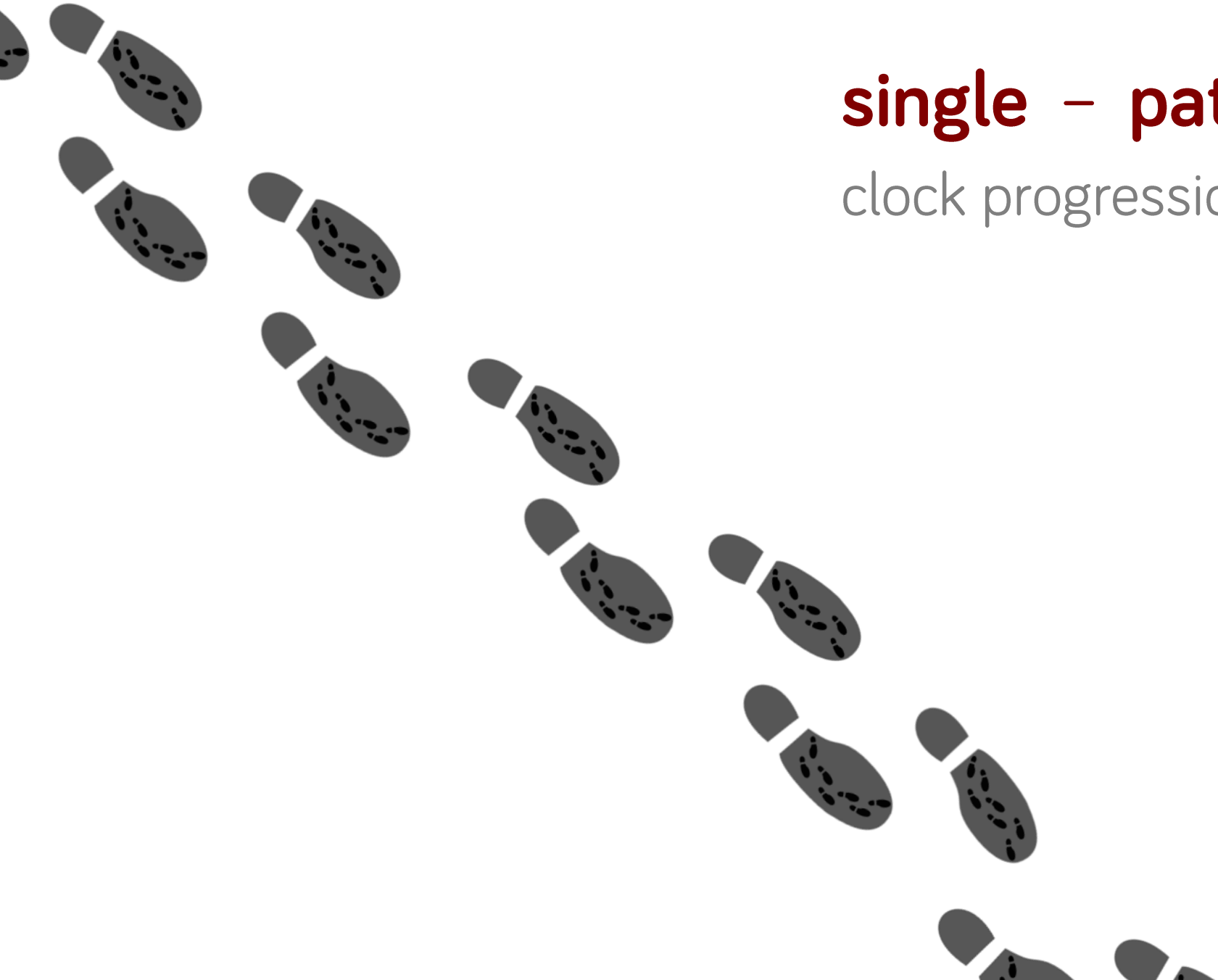
com-pos-ite

clock progression



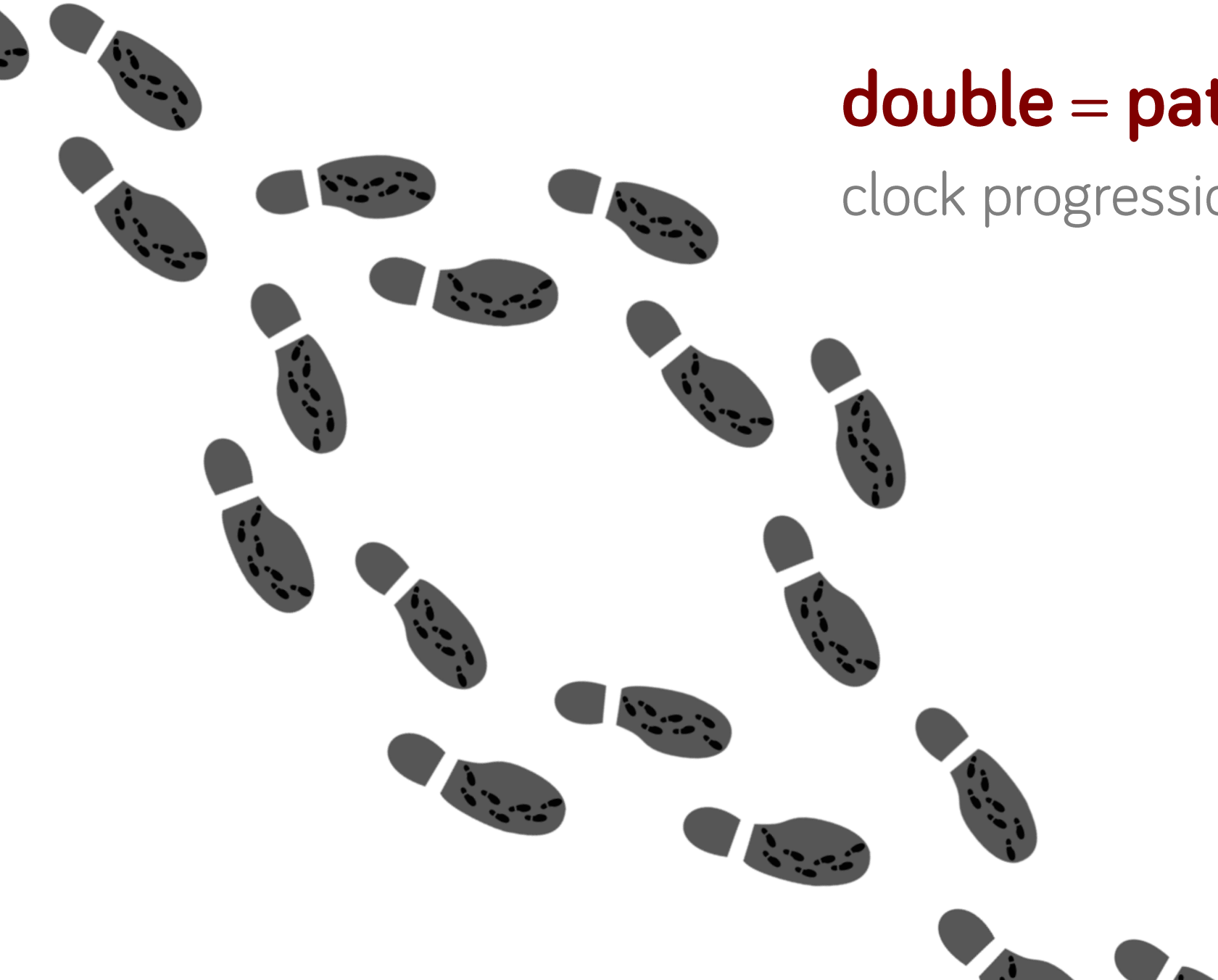
single – path

clock progression



double = path


clock progression



double = path

clock progression



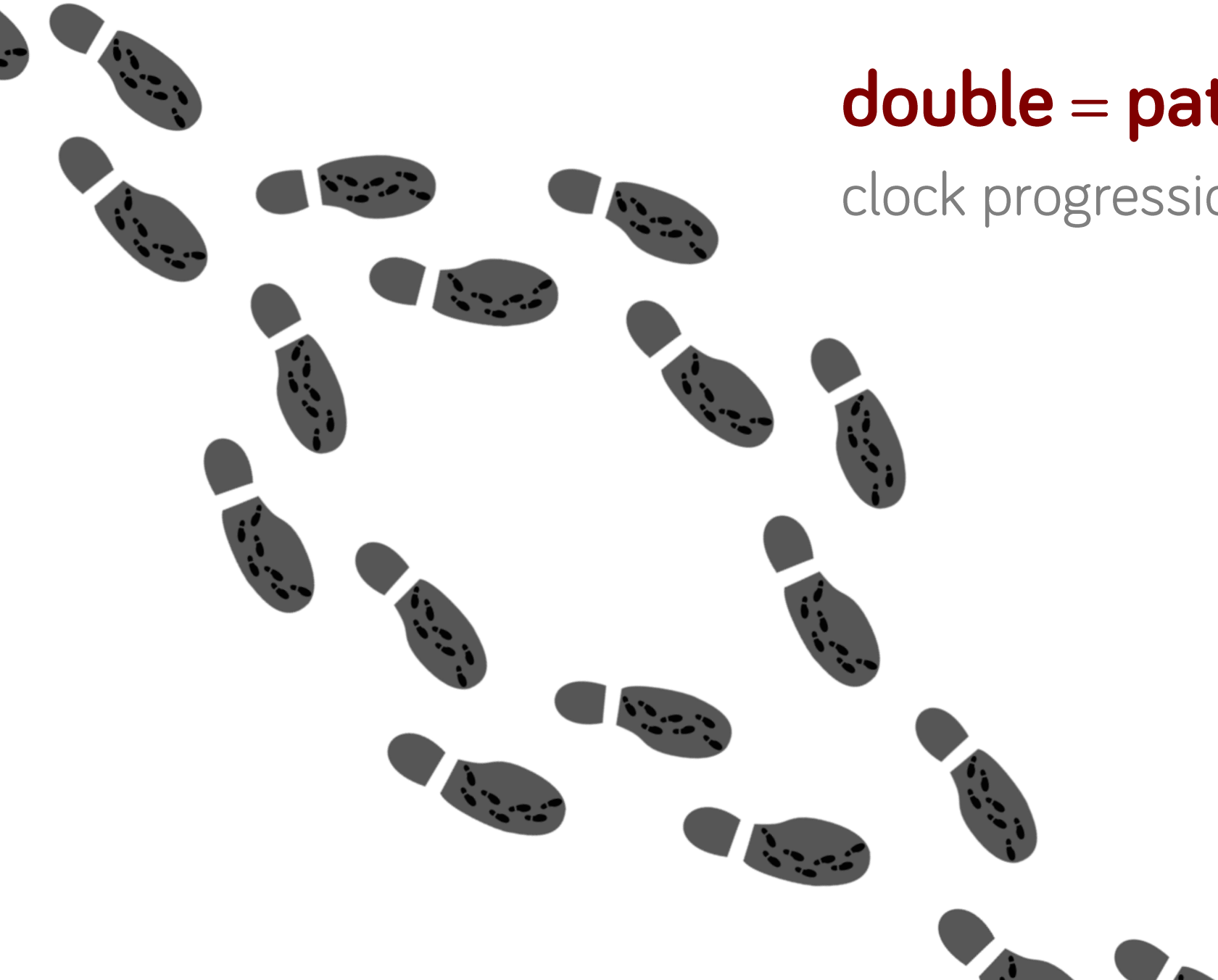


double = path

clock progression

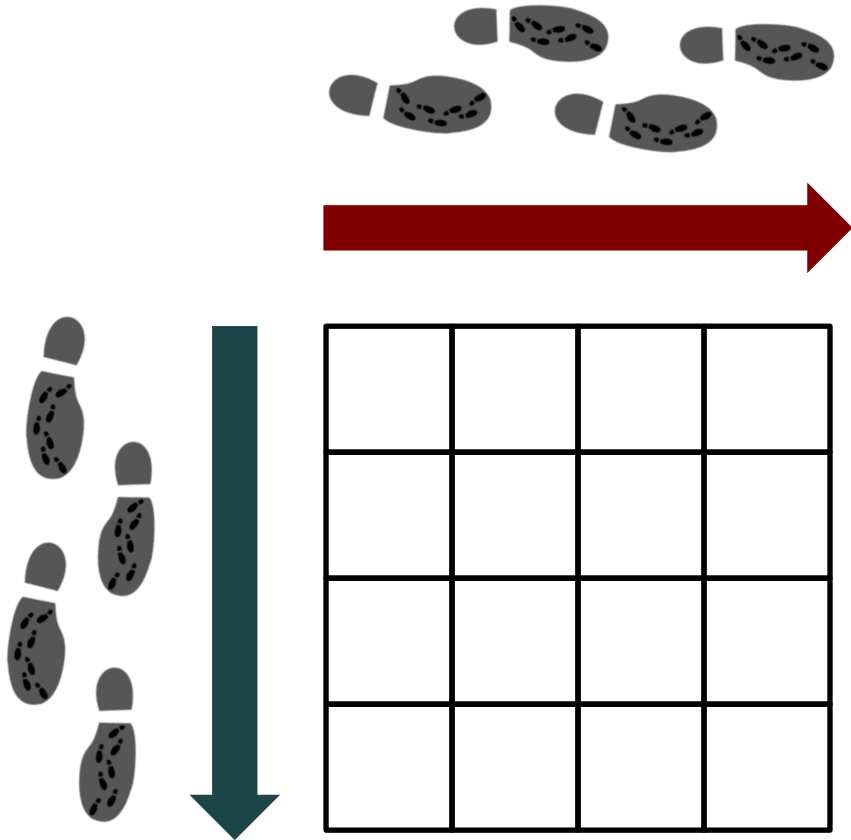
double = path

clock progression



2 clocks: 2D

clock progression



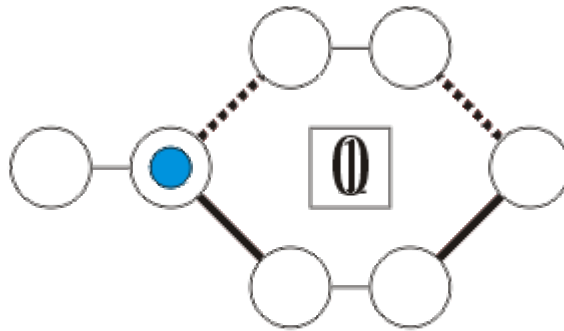
5 Applying 2-qubit gates 3-locally

- the railroad switch



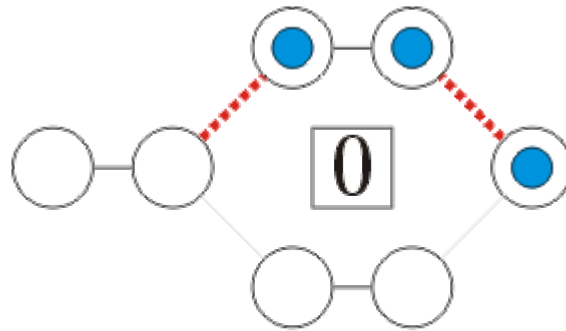
5 Applying 2-qubit gates 3-locally

- the railroad switch



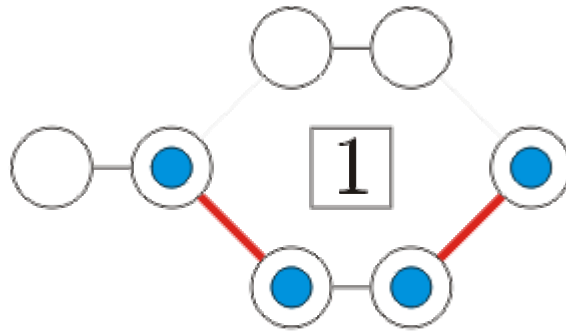
5 Applying 2-qubit gates 3-locally

- the railroad switch



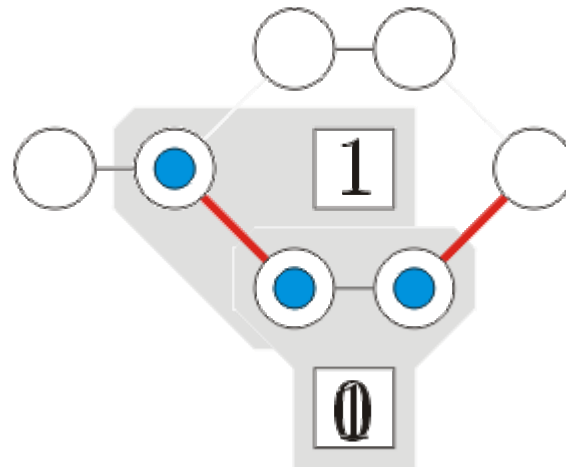
5 Applying 2-qubit gates 3-locally

- the railroad switch

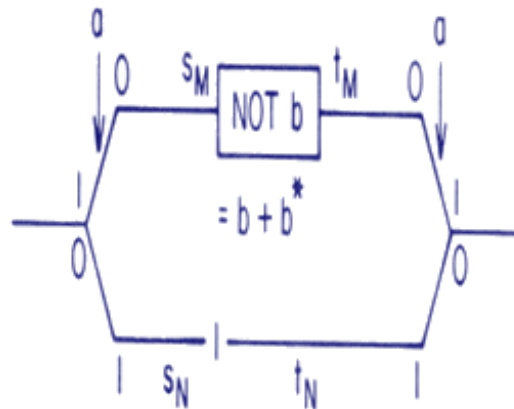


5 Applying 2-qubit gates 3-locally

- the railroad switch

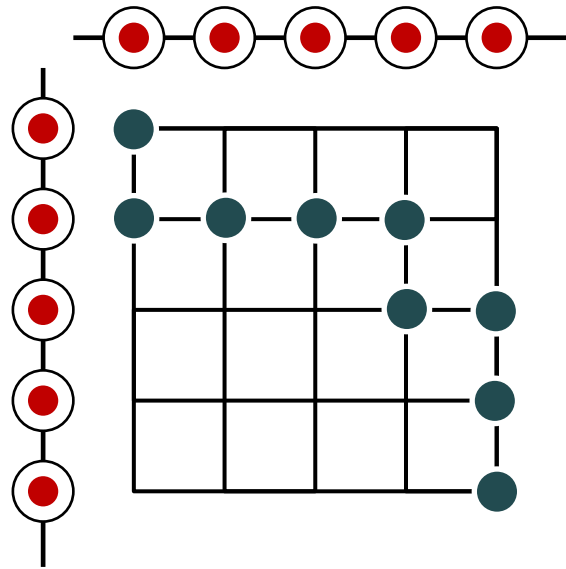


CNOT: 3-local
needs initialization



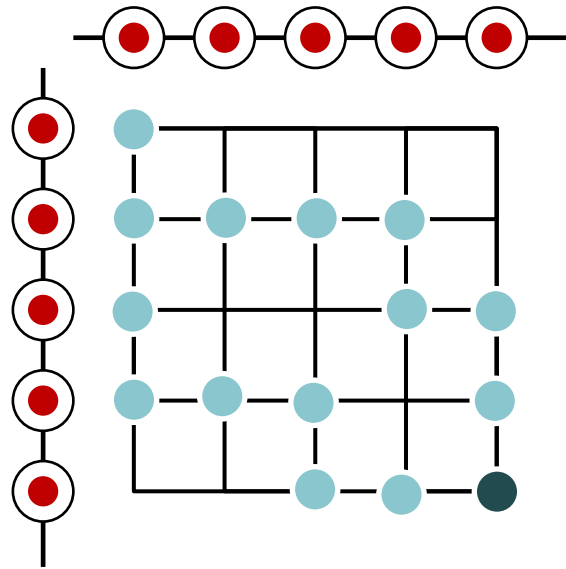
5 2D clocks (with two registers)

- two clocks



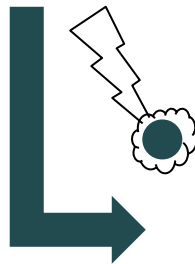
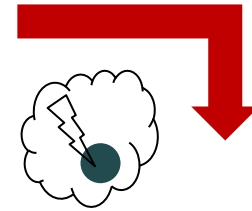
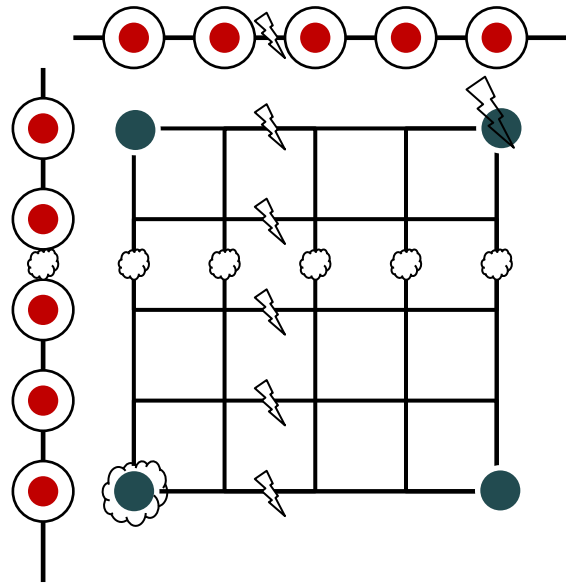
5 2D clocks (with two registers)

- two clocks



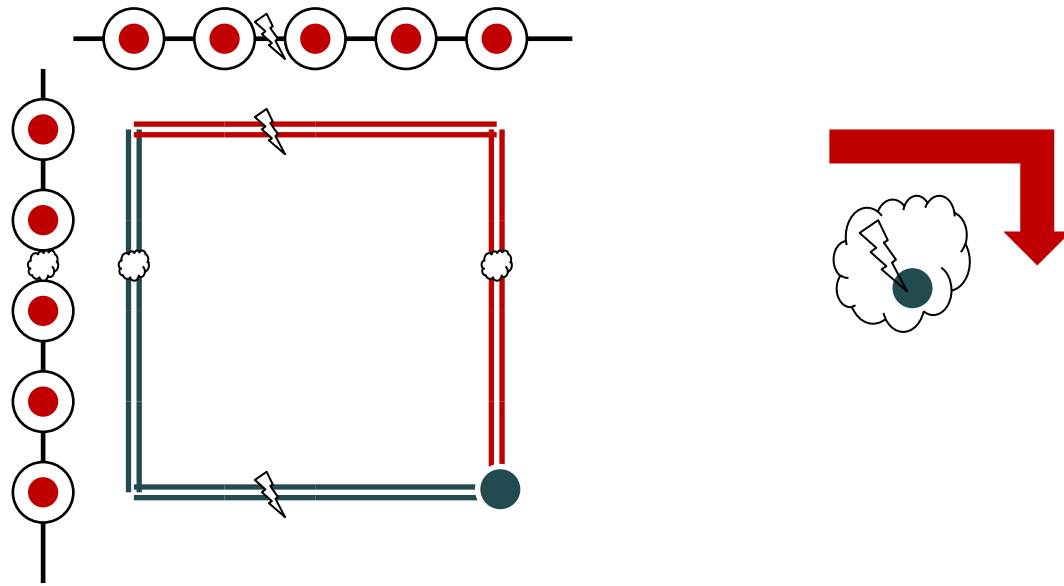
5 2D clocks (with two registers)

- add non-commuting (data) operations

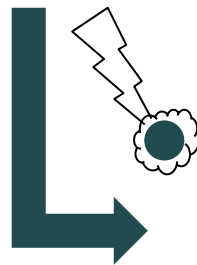


5 2D clocks (with two registers)

- like a railroad switch... with a single active site ensured

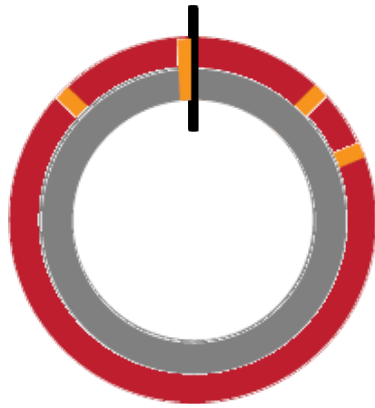


q3-SAT
QMA₁-comp.
[Gosset, N. '13]

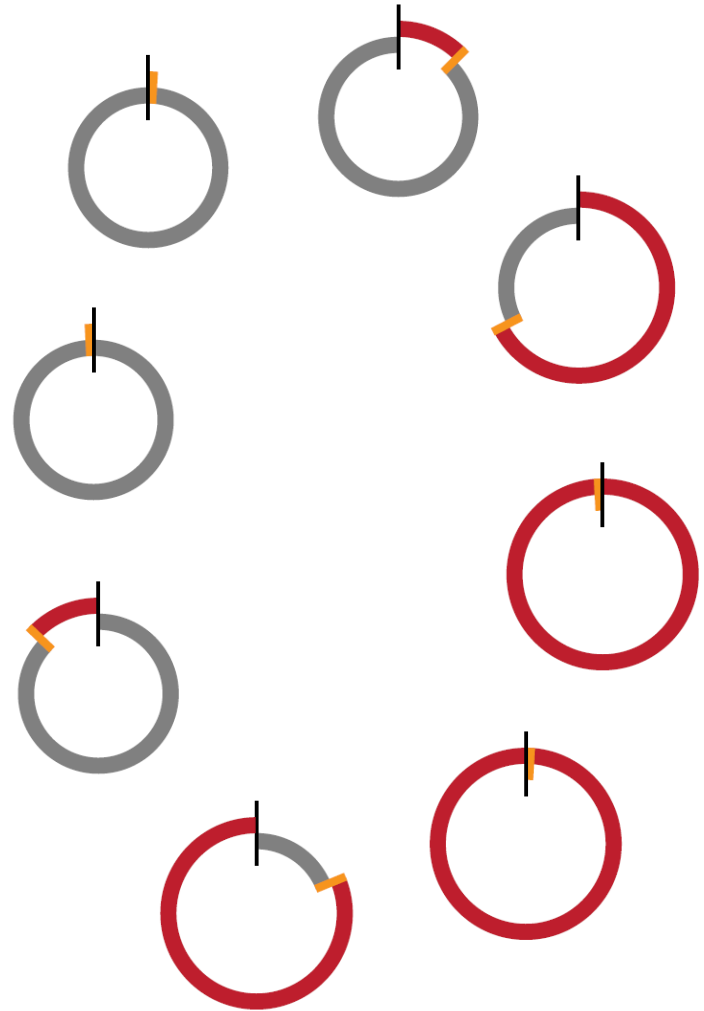


control: 0, 1

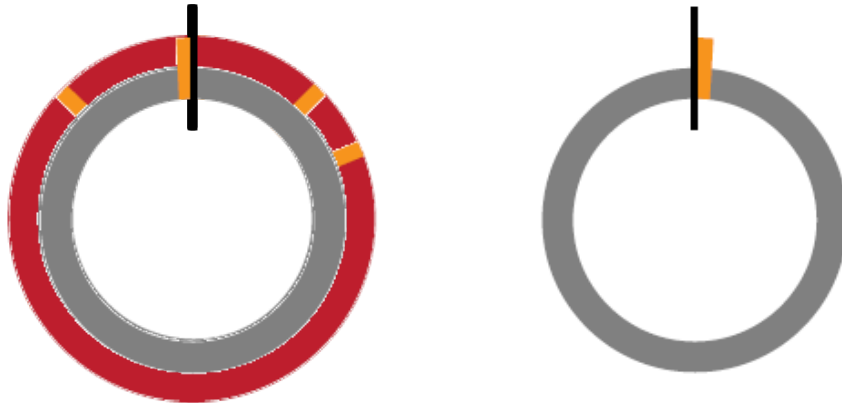
3 Can a clock be shorter than unary?



- a qutrit surfer on a cycle: $2C$ states



3 A smaller clock using two coupled cogs



3 A smaller clock using two coupled cogs



3 A smaller clock using two coupled cogs



- 2 cogs of length C give us $(2C)^2$ clock states
- transitions: 4-local, gates: 6-local (can be improved)
- the promise gap for a circuit with L gates: still L^{-2}



qPCP
& clocks

3 Questions about the qPCP conjecture

[Martin Schwarz]

- equivalence of the two formulations?

[AAV13]

LH with a fractional promise gap



translating the *random*
small verification to a LH?

look at a few qubits of a proof

- locally checking the
(expected) very entangled states?

[D.Aharonov, L.Eldar]

3 Questions about the qPCP conjecture

- clock constructions have a $1/\text{poly}$ promise gap
consistent, effective interaction strengthening?
error-correction/detection based quantum gadgets?

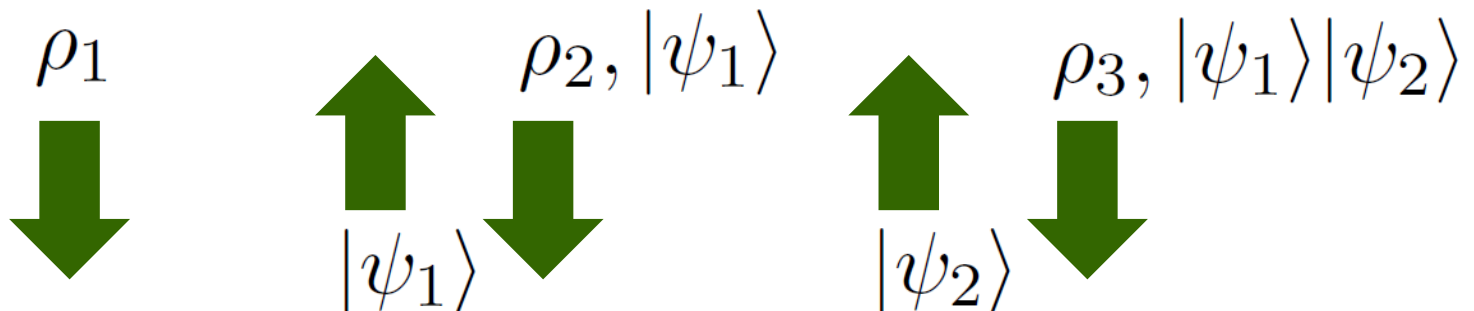
- direct Hamiltonian methods?

beyond the history state? [Itai Arad]

$$M := \mathbb{I} - H/m$$

$$\text{Tr}(M^\ell) \quad \begin{array}{l} \Gamma = 1/\text{poly}(n) \\ \ell = \Omega(mn/\Gamma) \end{array}$$

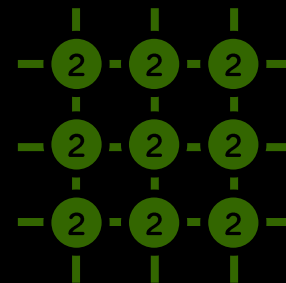
an interactive protocol to check the trace:



1

strong promises

and eigenvalue gaps



2

the history

of the history state



3

running the clock

precise/faulty, qubit/qudit, sequential/parallel



4

on the qPCP road

questions & warnings

